MOND In Galactic Scales

H.Haghi

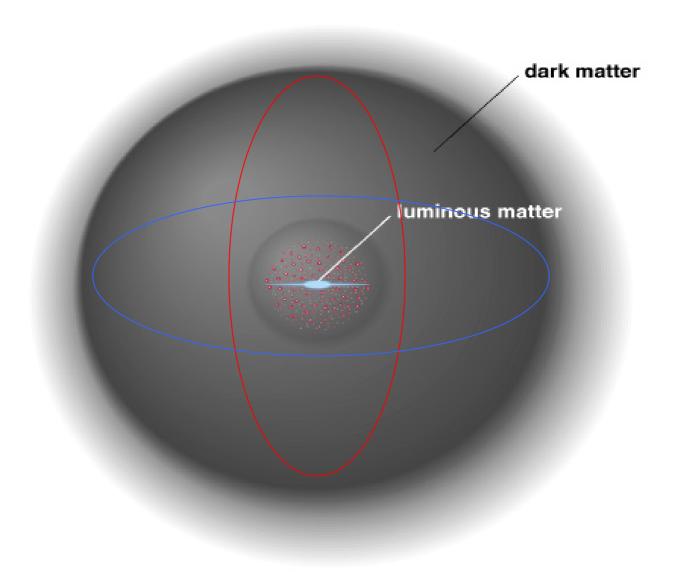
under supervising

S.Rahvar

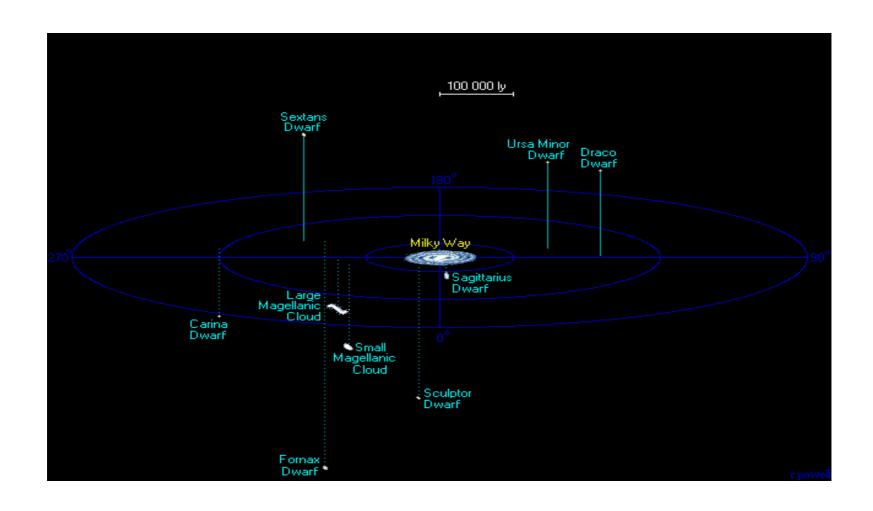
Sharif University of Technology - IRAN

How is the shape of halo??

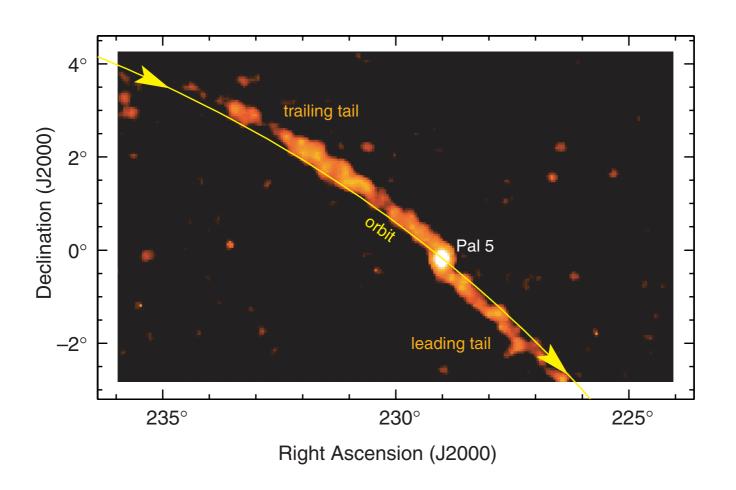
→ Galactic models



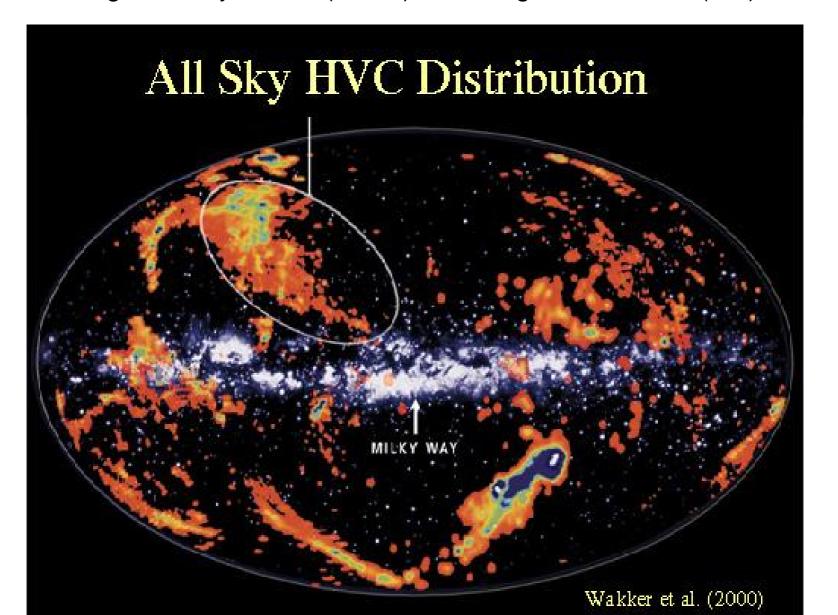
The MW's satellite provide information about the density of • Galactic halo at large radii.



Discover tidal streams from globular clusters and dwarf galaxies in the Galactic halo (Michael Odenkirchen, MPIA)



High velocity clouds (HVCs) and Magellanic Stream(MS)

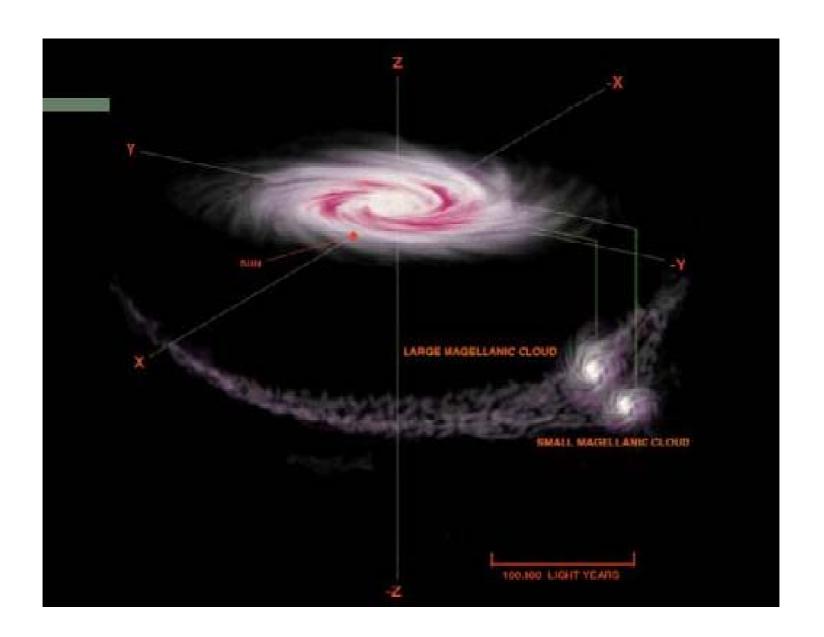


Magellanic systems

- A gaseous bridge connects the two galaxies (LMC, SMC) •

- The Magellanic Stream, is large gas stream extended from the MCs into the bridge between the Magellanic Clouds, and goes towards south galactic pole.

The Magellanic Clouds orbit each other and made a close • approach to the Milky Way 200-800 million years ago



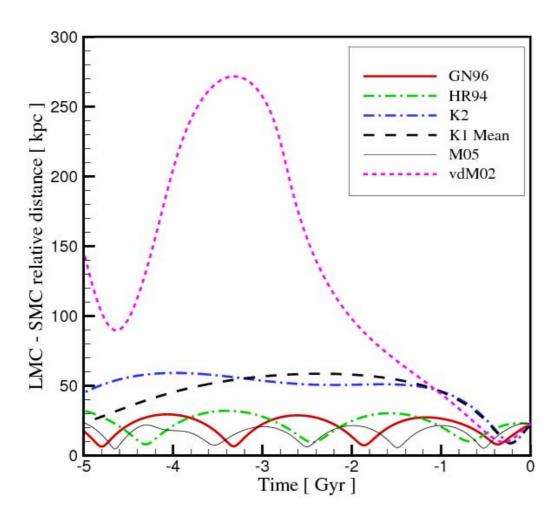
Dynamics of Magellanic Clouds •

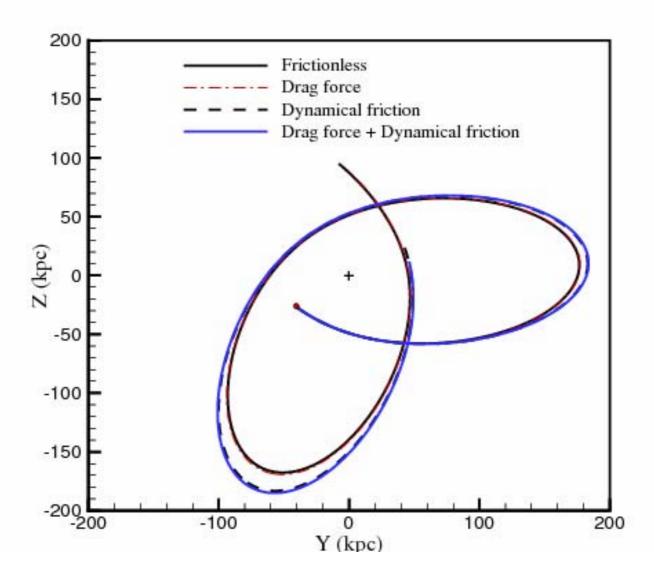
initial condition????

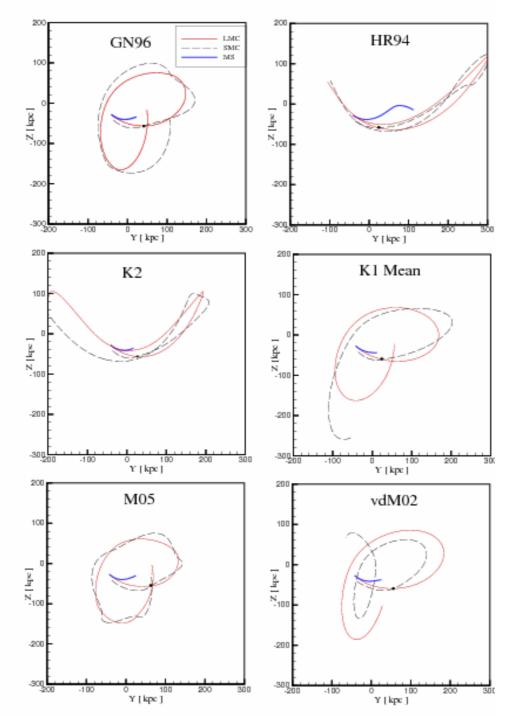
Dynamics of Magellanic Stream •

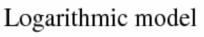
Various models for the origin and dynamics of MS •

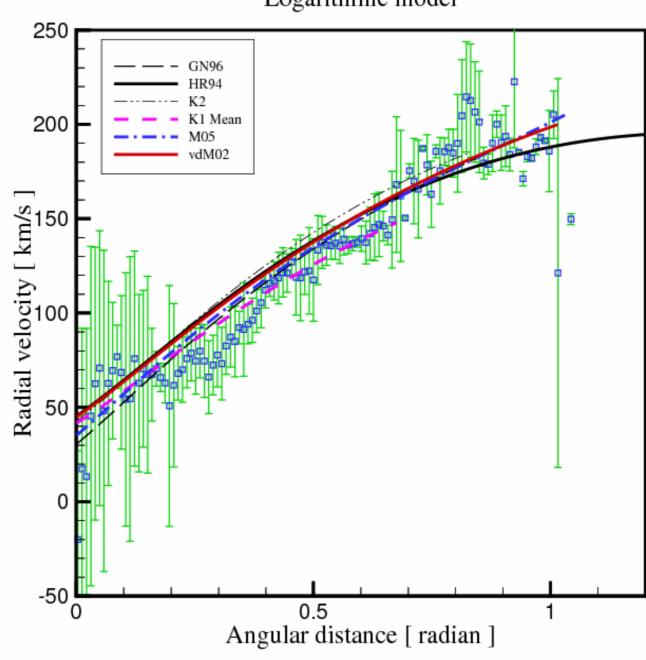
Magellanic Stream: A Possible Tool for Studying Dark Halo • Model (Submitted to A&A 2006)





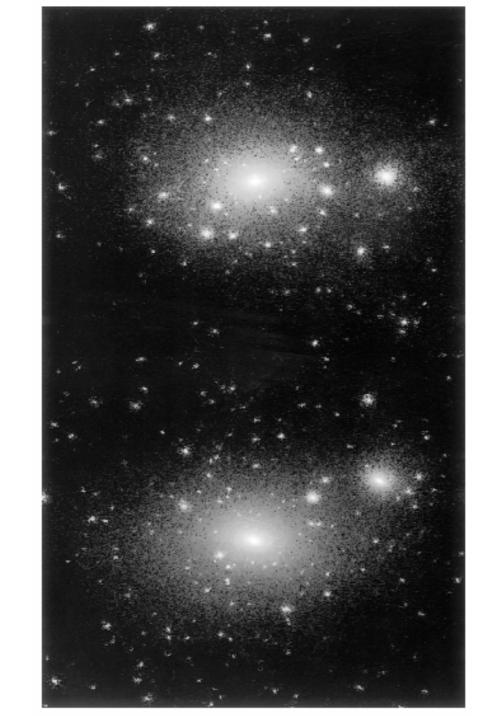


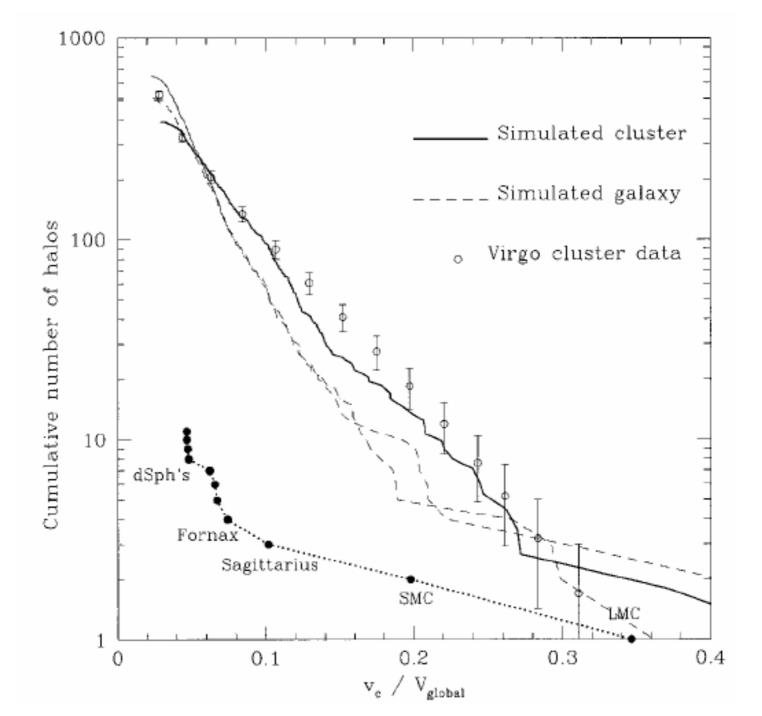


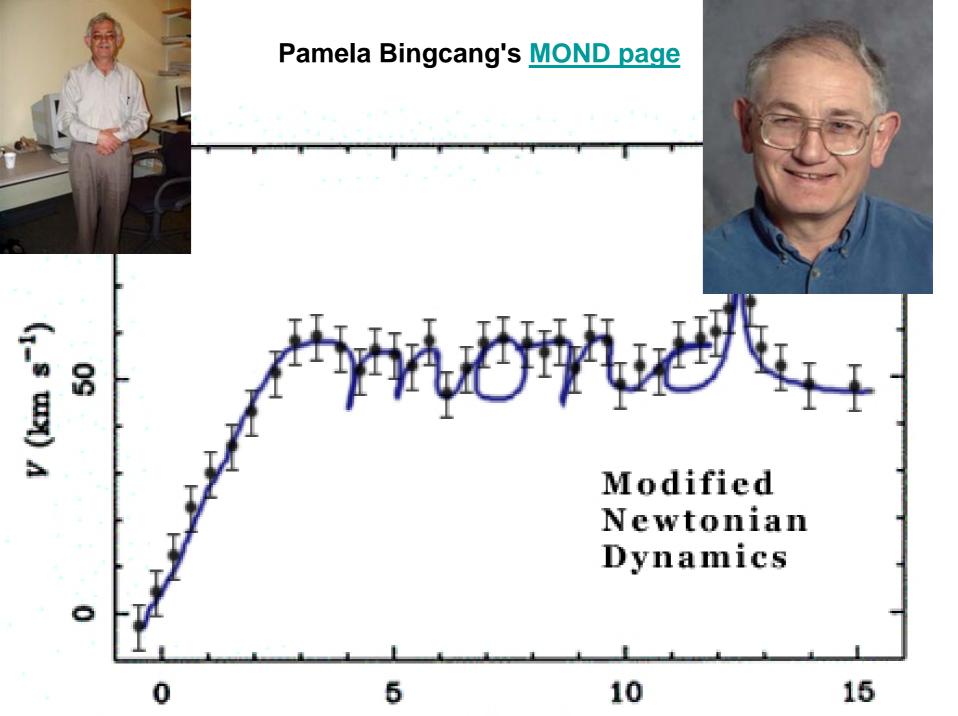


Problems with dark matter theory

- -- Prediction of more substructure than what is seen (Moore et al. 1999, Klypin et al 2000)
- -- Prediction of overly dense cores in the center of Galaxies and clusters (Moore et al. 1999, Klypin et al 2000).
- -- Fine tuning in dark matter halo shape parameters. In other word More parameters acquired to fit RC.
- -- Tully Fisher relation!! •
- -- and . . (Selwood and Kosowsky 2001)







MONDian potential of disk

$$\underline{\nabla} \cdot \left[\mu(|\underline{\nabla}(\Phi)|/a_0) \underline{\nabla}\Phi \right] = 4\pi G \rho \qquad \nabla \cdot \left[\frac{\|\nabla\Phi\|}{a_0} \nabla \Phi - \nabla \Phi_N \right] = 0$$

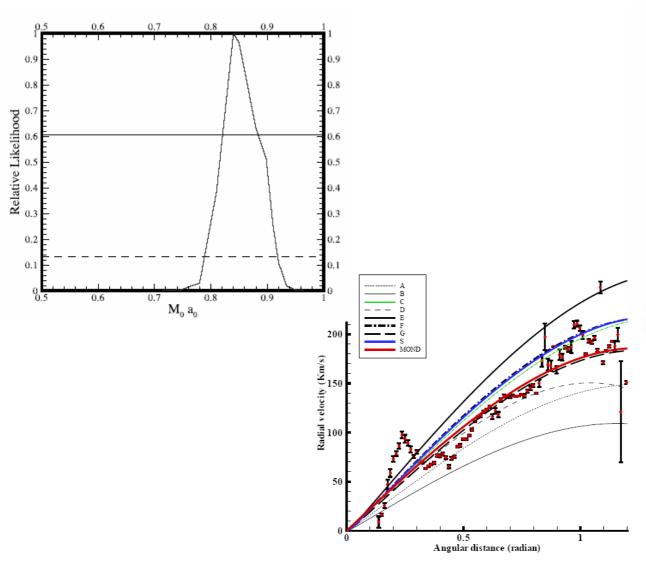
$$\frac{\|\nabla\Phi\|}{a_0} \nabla \Phi - \nabla \Phi_N = \nabla \times \mathbf{h}$$

$$\mu(|\underline{g}|/a_0)\underline{g} = \underline{g_N} \qquad \qquad \mu(x) = x(1+x^2)^{-1/2}$$

$$\underline{g} = \underline{g_N} \frac{\left(1 + \sqrt{1 + 4a_0^2/|\underline{g_N}|^2}\right)^{1/2}}{\sqrt{2}}$$

$$\Phi_N(R,z) = \frac{-GM}{\sqrt{R^2 + (a+|z|)^2}}$$

$$\Phi \simeq \frac{(MGa_0)^{1/2}}{2} \ln(R^2 + (|z| + a)^2)$$



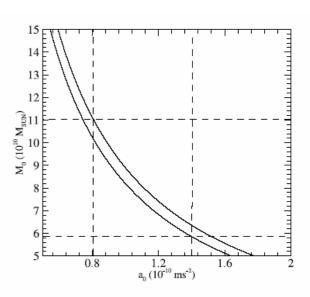


Fig. 4.—Constraints on $M_0a_0=8.4^{+0.5}_{-0.2}$ from the maximum likelihood analysis, with 1 σ confidence interval (M_0 in terms of $10^{10}~M_{\odot}$ and MONDian acceleration scale a_0 in terms of 10^{-10} m s⁻²). Using the range $a_0=0.8-1.4$ implies $M_0=6-11$ for the mass of the Kuzmin disk.

The Magellanic Stream in Modified Newtonian Dynamics Haghi, Rahvar & Hasani (2006 ApJ...652..354H)

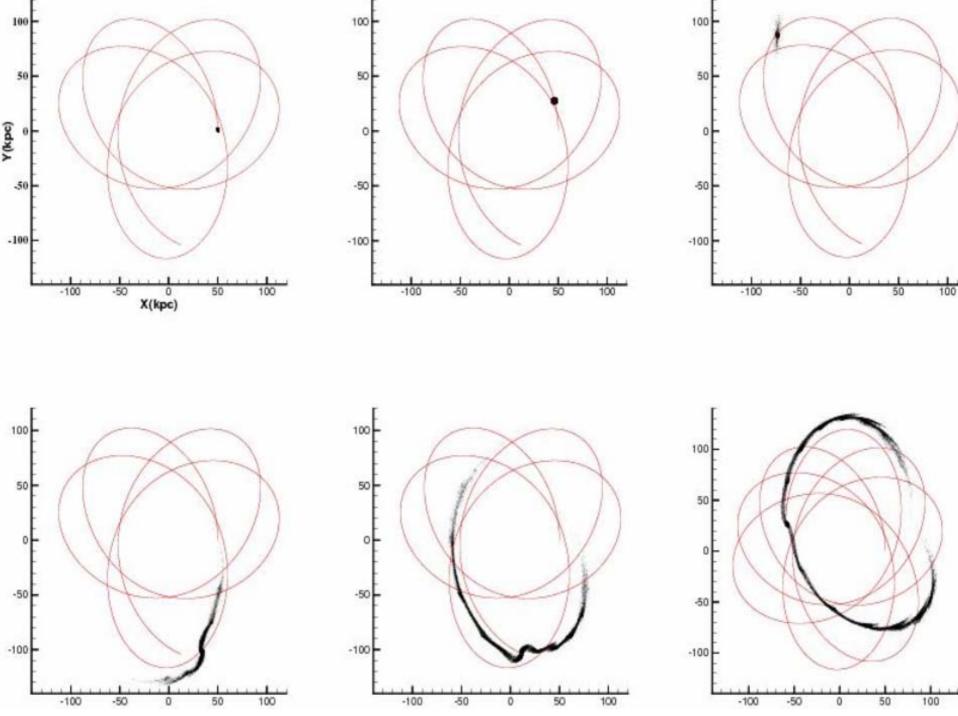
Globular Clusters

A Test of Newton's Law of Gravity in the Weak Acceleration Regime :

1- Globular Clusters as a Test for Gravity in the Weak Acceleration Regime

(R. Scarpa 2003, H.Baumgardt 2005, H.S.Zhao 2005)

2 -Testing MOND with Ultra-Compact Dwarf Galaxies (R. Scarpa 2005)



Orbital family of Kuzmin disk in MONDian regime (Haghi, Jalali, Zhao, Rahvar) In preparation

In dMOND regime the equation of motion is:

$$\ddot{z} = -\sqrt{GMa_0} \left(\frac{z \pm a}{R^2 + (z \pm a)^2}\right)$$

$$\ddot{R} = -\sqrt{GMa_0} \left(\frac{R}{R^2 + (z \pm a)^2}\right)$$

In TeVeS model $g = g_N + g_s$, So that for Bekenstein μ -function we have:

$$\mu(s)\mathbf{g}_s = \mathbf{g}_N, \mu(s) = \frac{g^s}{a_0}$$

$$\mathbf{g} = \mathbf{g}_N + \mathbf{g}_N \sqrt{\frac{a_0}{g_N}}$$

$$\ddot{z} = -\left(\frac{\sqrt{GMa_0}}{R^2 + (z \pm a)^2} + \frac{GM}{[R^2 + (z \pm a)^2]^{3/2}}\right)(z \pm a)$$

$$\ddot{R} = -\left(\frac{\sqrt{GMa_0}}{R^2 + (z \pm a)^2} + \frac{GM}{[R^2 + (z \pm a)^2]^{3/2}}\right)R$$

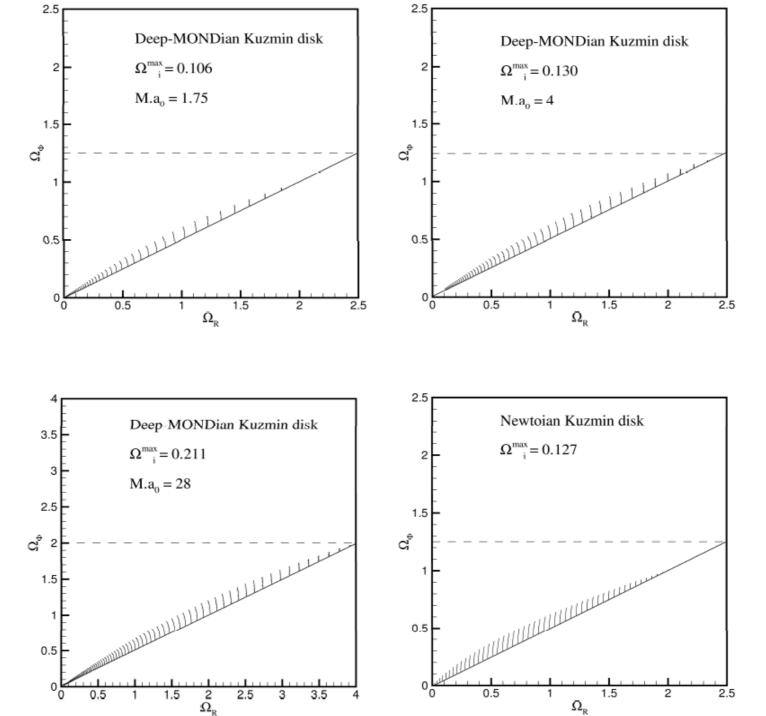
Orbital in Equatorial Plane

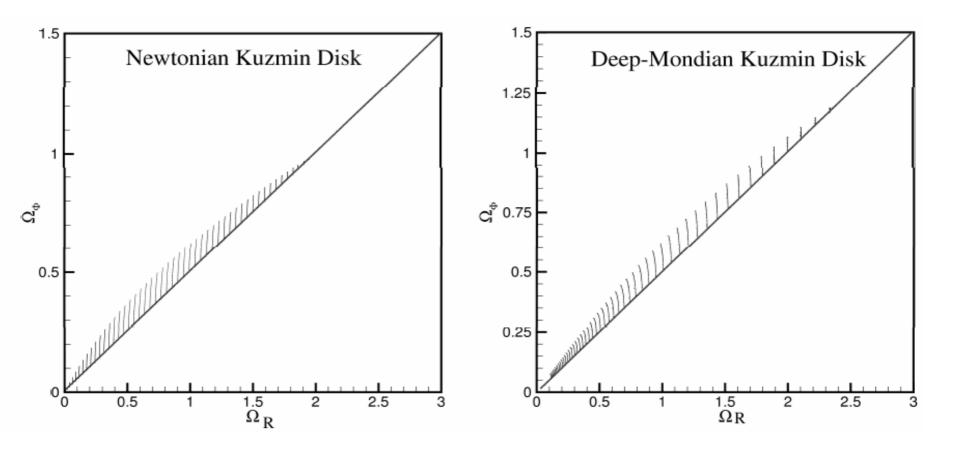
In Teves:

$$\ddot{R} = -\left(\frac{\sqrt{GMa_0}}{R^2 + a^2} + \frac{GM}{[R^2 + a^2]^{3/2}}\right)R$$

In CDM:

$$\ddot{R}^N = \left(\frac{-V_0^2}{[R^2 + R_c^2]} + \frac{GM}{[R^2 + a^2]^{3/2}}\right)R$$





Galaxy merging in MOND

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$$v_0^2 = 4G(M_* + M_{\rm DM}) \left(\frac{1}{d_0} - \frac{1}{d_{\rm rest}}\right),$$

while in the MOND cases

$$v_0^2 \simeq 0.8\sqrt{8GM_*a_0} \ln \frac{d_{\rm rest}}{d_0},$$

