Neutron stars with Quark Core (Hybrid Stars)

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OUTLINE

- Introduction
- EOS of Nuclear matter and Neutron Star
- EOS of Quark matter
- Hadron-quark phase transition
- Structure of Hybrid stars
- Summary

Stellar evolution



Possible internal structures and compositions of four different types of compact stars





A semi-quantitative phase diagram on $T-\mu$ plane

EOS

Equation of states(EOS) of nuclear matter at high density

* Hadron matter

- * Nuclear many body theories: Brueckner-Bethe-Goldstone(BBG),....
- * Relativistic mean field theory
- * Quark matter
 - * MIT bag model
 - * Dyson-Schwinger equations (DSE)
 - * Nambu Jona-Lasino (NJL) model (color superconductivity)

* ...

Eos of Hadron Matter

We have calculated the equation of state of hadronic matter, we use statistical Thomas-Fermi approach by employing Myers-Swiatecki density dependent potential.

Relativistic Star

$$ds^{2} = e^{\nu(r)}c^{2}dt^{2} - e^{\lambda(r)}dr^{2} - r^{2}(d\theta^{2} + sin^{2}\theta d\phi^{2}) \longrightarrow \text{Metric}$$

$$T^{\mu}_{\nu} = (p + \epsilon)u^{\mu}u_{\nu} + p\delta^{\mu}_{\nu} \longrightarrow \text{Energy-Momentum Tensor (EOS)}$$

$$G_{\nu\mu} = 8\pi T_{\nu\mu} \longrightarrow \text{Einstein's Equation}$$

$$\frac{dp}{dr} = -\frac{[p(r) + \epsilon(r)][m(r) + 4\pi r^{3}p(r)]}{r(r - 2m(r))} \longrightarrow \text{TOV equation (General Relativistic Hydrostatic equilibrium)}$$

$$m(r) = 4\pi \int_{0}^{r} drr^{2}\epsilon(r) \longrightarrow \text{Stellar mass}$$

$$r \rightarrow 0 \quad m(r) \rightarrow 0 \quad P \rightarrow P_{c} \quad \epsilon \rightarrow \epsilon_{c}$$
can easily be solved numerically for given EoS $\epsilon(p)$
initial value: $p(r = 0) = p_{0}$
surface of the star: $p(r = R) = 0$
gravitational mass of the star: $M \equiv m(R)$

EOS of Quark Matter (MIT Bag Model)

$$\begin{split} P_q &= -B + P_q^{kin} + P_q^{int} \\ \varepsilon_q &= B + \varepsilon_q^{kin} + \varepsilon_q^{int} \end{split}$$

$$\Omega = \Omega_u + \Omega_d + \Omega_s + \Omega_e + B \qquad \qquad \Omega_e$$

$$\begin{split} \Omega_q = & - & \frac{3m_q^4}{8\pi^2} \big[\frac{\eta_q x_q}{3} (2x_q^2 - 3) + \ln(x_q + \eta_q) \big] \\ & + & \frac{3m_q^4 \alpha_s}{4\pi^3} \Big\{ 2 \big[\eta_q x_q - \ln(x_q + \eta_q) \big]^2 - \frac{4}{3} x_q^4 + 2\ln(\eta_q) \\ & + & 4 \ln(\frac{\sigma_{\text{ren}}}{m_q \eta_q}) \big[\eta_q x_q - \ln(x_q + \eta_q) \big] \Big\} \end{split}$$

 $\frac{\mu_{\theta}^{*}}{12\pi^{2}}$

 $\begin{array}{l} m_q \ , \ \mu_q \ : \ q \ \text{quark mass and chemical potential.} \\ x_q = \sqrt{\mu_q^2 - m_q^2}/m_q \\ \eta_q = \sqrt{1 + x_q^2} = \mu_q/m_q \\ \alpha_s \ : \ \text{QCD fine structure constant} \end{array}$

$$ho_q = -rac{\partial \Omega_q}{\partial \mu_q}$$

$$\epsilon_Q = \sum_q (\Omega_q + \mu_q \rho_q) + B$$
$$P_Q = -\sum_q \Omega_q - B$$

$$\begin{split} \mu_s &= \mu_d \equiv \mu \,, \quad \mu_u = \mu - \mu_e \,, \\ \frac{2}{3} n_u - \frac{1}{3} n_d - \frac{1}{3} n_s - n_e &= 0 \end{split}$$



$$B(\rho) = B_{\infty} + (B_0 - B_{\infty}) \exp\left[-\beta \left(\frac{\rho}{\rho_0}\right)^2\right]$$

with $B_{\infty} = 60 \text{ MeV fm}^{-3}$, $B_0 = 400 \text{ MeV fm}^{-3}$, and $\beta = 0.17$.

Hadron-quark phase transition

Gibbs condition at zero temperature between hadron phase and quark phase

 $p_{HP}(\mu_n, \mu_e) = p_{QP}(\mu_n, \mu_e).$

the total electrical charge is

$$0 = \frac{Q}{V} = \chi q_{QP} + (1 - \chi)q_{HP}$$

and the total energy density is

$$\epsilon = \frac{E}{V} = \chi \epsilon_{QP} + (1 - \chi) \epsilon_{HP}$$





Results of hadron matter







Results of quark matter





Conclusions





References:

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• Weber et al. ArXiv: 0705.2708

•H.R. Moshfegh et al. J.Phys.G38:085102, 2011.

Thank you for your attention.