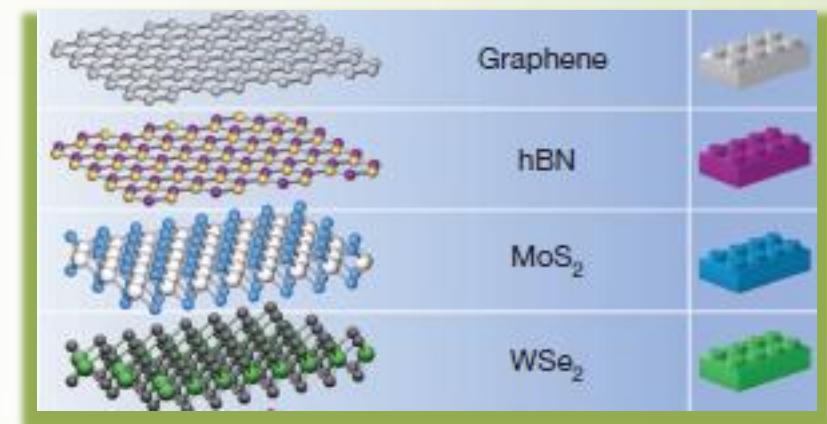
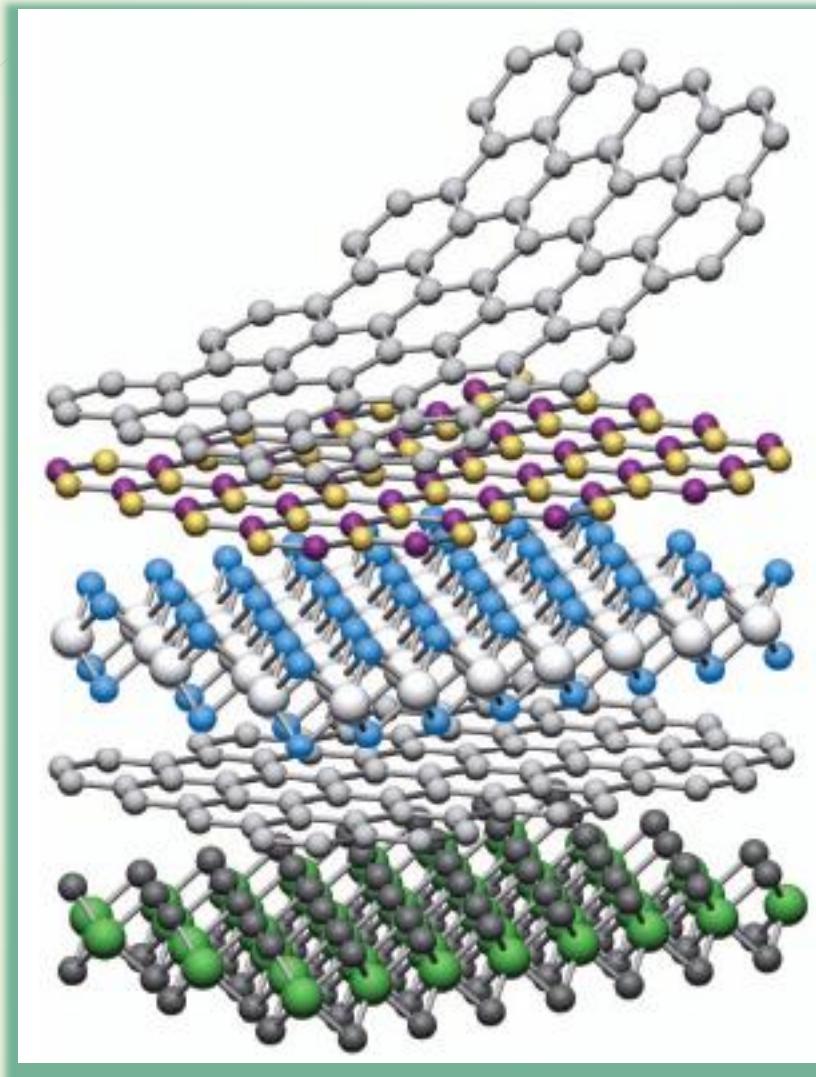


Introduction on van der Waals layered materials

Reza Asgari



van der Waals heterostructures



A. K. Geim and I . V. Grigoieva, Nature, **499**, 419 (2013)

Outline

- Introduction
- Graphene-like materials (BN, Silicene,...)
- TMDs, MoS₂
- Conclusion

Layered van der Waals solids

These crystal structures feature neutral, single-atom-thick or polyhedral-thick layers of atoms that are covalently or ionically connected with their neighbors within each layer, whereas the layers are held together via van der Waals bonding along the third axis.

Layered van der Waals solids

- Mechanical exfoliation of large crystals using “Scotch tape”
- Chemical exfoliation

Novoselov, *et al* Nature **438**, 197 (2005)

Epitaxy, requires ultrahigh vacuum conditions: Expensive Science **312**, 1192 (2006)
Various chemical methods. Nano Lett. **8**, 2442 (2008) , Nature Nanotech. **3**, 270 (2008); *ibid* **4**, 217(2009)
Chemical Vapour Deposition : Nature **457**, 706 (2009), Nano Lett. **9**, 30 (2009)

Prerequisite for having 2D

- 3D materials with melting temperature over 1000
- 3D parents must be chemically inert and exhibit no decomposed surface layer in air
- Insulating and semiconducting 2D crystals are more likely to be stable than metallic ones

Graphene

Many extraordinary properties, such as its 2.3% absorption in the white light spectrum, high Young's modulus and excellent thermal conductivity, have all been reported.

Using graphene in a wide range of areas, including high-speed electronic and optical devices, energy generation and storage, hybrid materials, chemical sensors, optoelectronics and even DNA sequencing.

Different techniques:

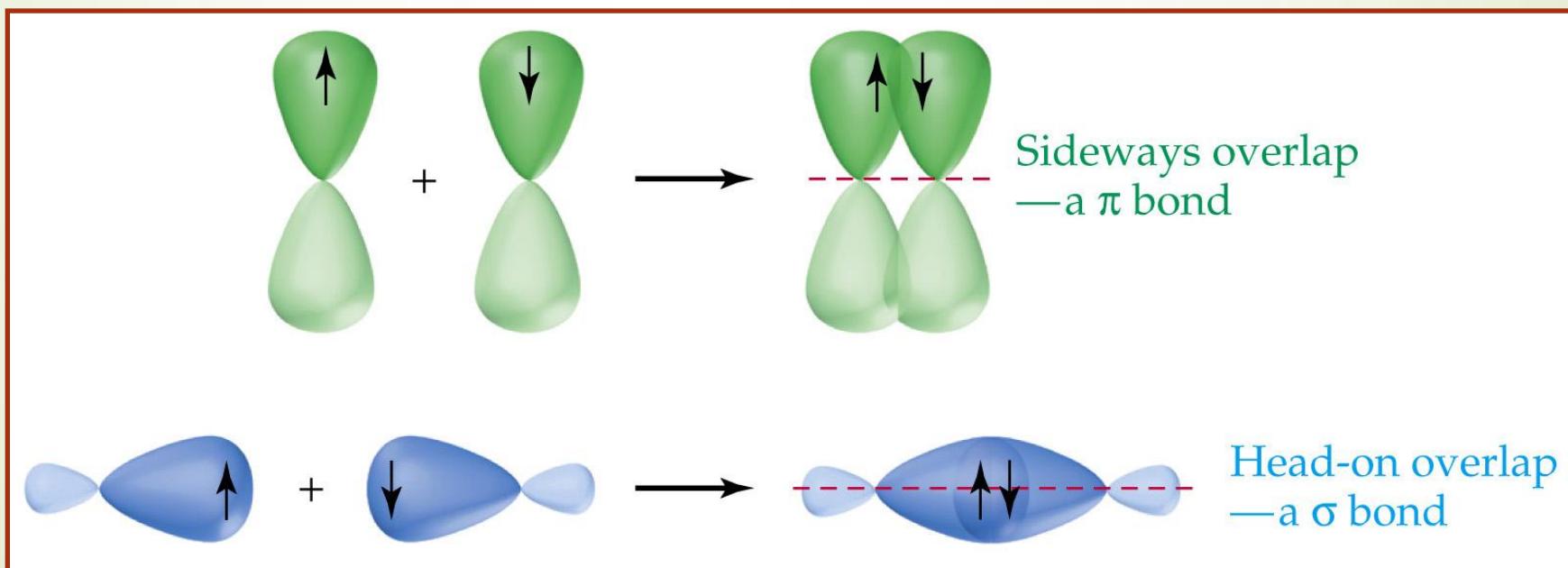
Mechanical exfoliation, liquid-phase exfoliation, reduction of graphene oxide, chemical vapor deposition (CVD), surface segregation, and molecular beam epitaxy (MBE).

The fine control of the number and structure of graphene sheets over an entire substrate remains a major challenge.

Atomic structure: Graphene

9

- There are four valence electrons.
- The 2s and 2p can form hybridized orbital
- Two side-to-side $2p_z$ orbitals form a p bond
- Two co-axial $2p$ or sp orbitals form the s bond



2D layered materials

metal chalcogenides, transition metal oxides, and other 2D compounds layered transition oxides such as MoO_3 and La_2CuO_4 ,²² insulator hexagonal boron nitride (h-BN), and topological insulators of Bi_2Te_3 , Sb_2Se_3 , Bi_2Se_3 and silicene and germanene

The common feature of these layered materials is that the bulk 3D crystals are stacked structures.

They involve van der Waals interactions between adjacent sheets with strong covalent bonding within each sheet.

2D layered materials: Library

Graphene family	Graphene	hBN 'white graphene'		BCN	Fluorographene	Graphene oxide
2D chalcogenides	MoS ₂ , WS ₂ , MoSe ₂ , WSe ₂		Semiconducting dichalcogenides: MoTe ₂ , WTe ₂ , ZrS ₂ , ZrSe ₂ and so on		Metallic dichalcogenides: NbSe ₂ , NbS ₂ , TaS ₂ , TiS ₂ , NiSe ₂ and so on	
				Layered semiconductors: GaSe, GaTe, InSe, Bi ₂ Se ₃ and so on		
2D oxides	Micas, BSCCO	MoO ₃ , WO ₃		Perovskite-type: LaNb ₂ O ₇ , (Ca,Sr) ₂ Nb ₃ O ₁₀ , Bi ₄ Ti ₃ O ₁₂ , Ca ₂ Ta ₂ TiO ₁₀ and so on		Hydroxides: Ni(OH) ₂ , Eu(OH) ₂ and so on
	Layered Cu oxides	TiO ₂ , MnO ₂ , V ₂ O ₅ , TaO ₅ , RuO ₂ and so on				Others

Electronic band gaps

2D sheets	theoretical E_g (eV)	experimental E_g (eV)
graphene	0	0
bilayer graphene	0	0
bulk <i>h</i> -BN		5.97 [ref 52]
monolayer <i>h</i> -BN		6.07 [ref 65]
fully hydrogenized <i>h</i> -BN	3.05 [ref 66]	
2–5 layers <i>h</i> -BN		5.92 [ref 105.]
bulk MoS ₂	1.20 (indirect ^b) [refs 35, 139]	1.0–1.29 (indirect) [refs 35, 139]
monolayer MoS ₂ ^a	~1.90 (direct ^b) [ref 140]	~1.90 (direct) [ref 140]
bulk WS ₂	~1.30 (indirect ^b) [refs 35, 147]	~1.35 (indirect) [refs 35, 147]
monolayer WS ₂ ^a	~2.10 (direct ^b) [ref 147] ~1.80 (direct ^c) [ref 148]	
monolayer MoSe ₂	~1.44 (direct ^c) [ref 148]	
monolayer MoTe ₂	~1.07 (direct ^c) [ref 148]	

Electronic structures:

From insulator to metal

Include the topological insulator effect, superconductivity, and thermoelectricity

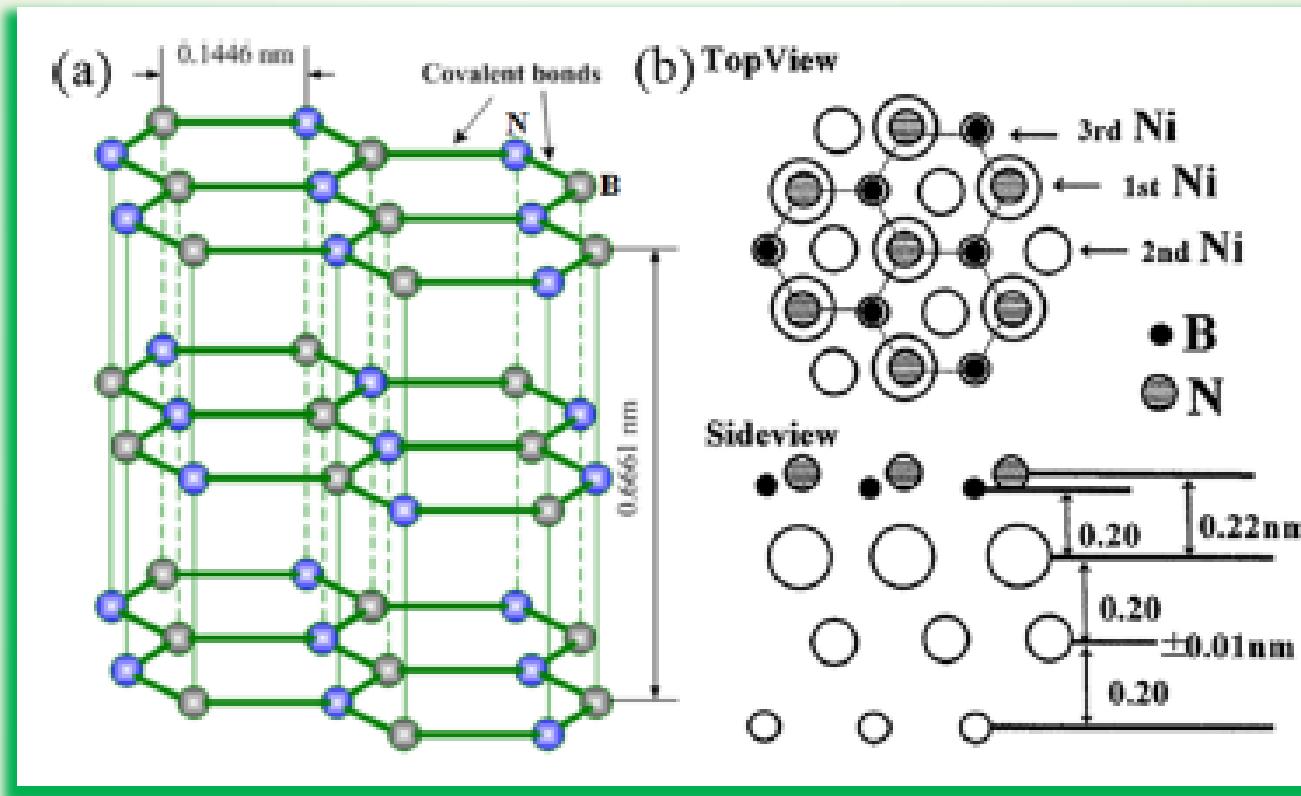
Applications:

optoelectronics, spintronics, catalysts, chemical and biological sensors, supercapacitors, solar cells, and lithium ion batteries.

Specific examples

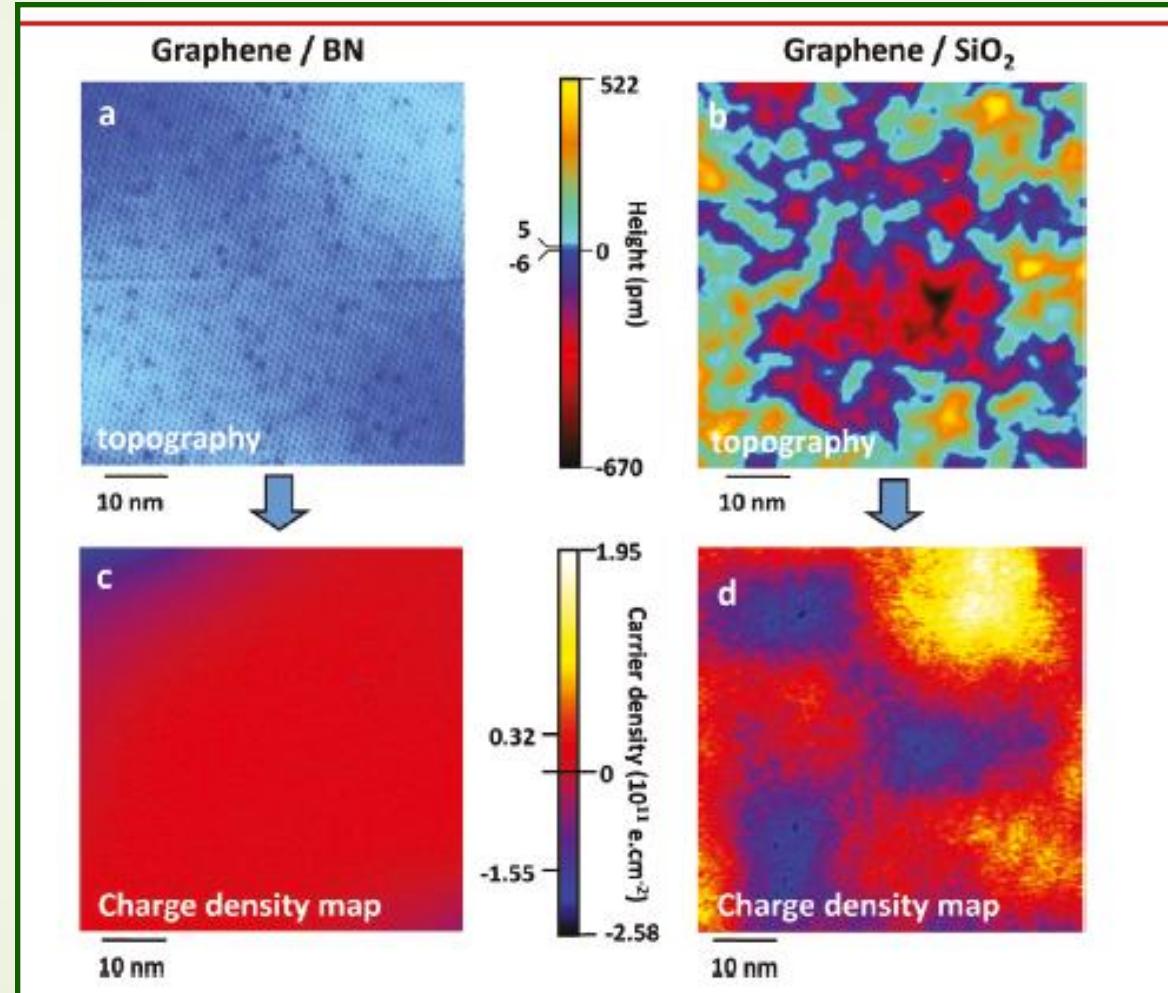
Boron Nitride

Hexagonal Boron Nitride



$$\begin{array}{ll} \textit{Bulk h-BN} & E_g = 5.97 \text{ eV} \\ \textit{monolayer h-BN} & E_g = 6.07 \text{ eV} \end{array}$$

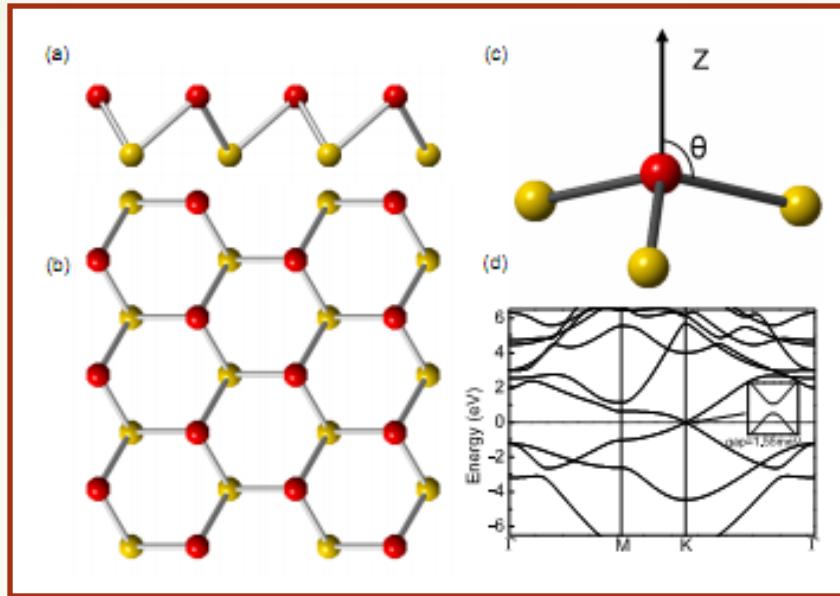
Topography and charge density



Specific examples

Silicene

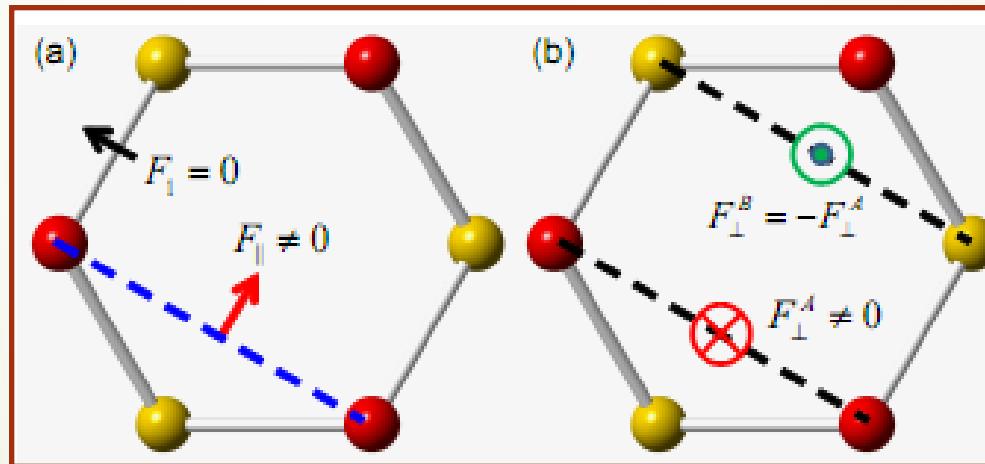
Silicene, Germanium and Tin



The lattice geometry of low-buckled silicene. Note that A sublattice (red) and B sublattice (yellow) are not coplanar. The angle is defined as being between the Si-Si bond and the z direction normal to the plane. Inset: zooming in the energy dispersion near the K point and the gap induced by SOC.

Silicene: Spin-orbit interaction

$$H_{so} = \frac{\hbar}{4m_0^2c^2} (\nabla V \times \vec{p}) \cdot \vec{\sigma} = -\frac{\hbar}{4m_0^2c^2} (\vec{F} \times \vec{p}) \cdot \vec{\sigma},$$



Cheng-Cheng Liu, Hua Jiang, and Yugui Yao, Phys. Rev. B (2011)

Low-energy model Hamiltonian

$$\begin{aligned}
H_K^{eff}(\theta) &= H_K \otimes I_2 + H_{so}^{1st} + H_{so}^{2st} + H_R(k) \\
&= (\varepsilon_1 - \lambda_{so}^{2st})I_4 + \begin{pmatrix} h_{11} & v_F k_+ \\ v_F k_- & -h_{11} \end{pmatrix},
\end{aligned}$$

$$h_{11} \equiv -\lambda_{so}\sigma_z - a\lambda_R(k_y\sigma_x - k_x\sigma_y),$$

$$k_+ = k_x + ik_y, k_- = k_x - ik_y,$$

$$\lambda_{so}^{2st} \equiv \left(\frac{\xi_0}{2}\right)^2 \frac{2}{9} \frac{-\Delta}{\sin^2 \theta V_{sp\sigma}^2}.$$

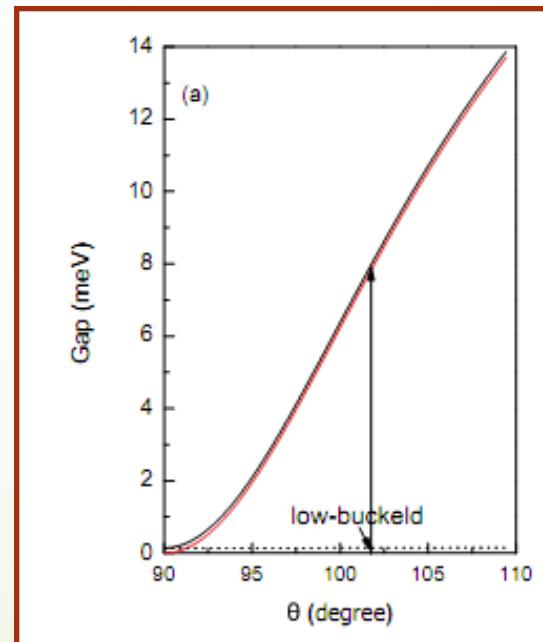
$$\begin{aligned}
v_F &= \frac{-\sqrt{3}a}{2} [u_{11}^2 (V_{pp\pi} \sin^2 \theta + V_{pp\sigma} \cos^2 \theta) - u_{21}^2 V_{ss\sigma} \\
&\quad + 2u_{11}u_{21} \cos \theta V_{sp\sigma} - \frac{1}{2}|u_{31}|^2 \sin^2 \theta (V_{pp\sigma} - V_{pp\pi})],
\end{aligned}$$

$$\begin{aligned}
\lambda_R &= \frac{i\xi_0}{\sqrt{2}} \frac{u_{11}u_{32} - u_{31}u_{12}}{(\varepsilon_2 - \varepsilon_1)a} \times \\
&\quad [(u_{12}u_{21} + u_{22}u_{11})v_5 + u_{22}u_{21}v_7 + u_{12}u_{11}v_4 - 2u_{32}u_{31}v_1] \\
&\quad + \frac{i\xi_0}{\sqrt{2}} \frac{u_{11}u_{33} - u_{31}u_{13}}{(\varepsilon_3 - \varepsilon_1)a} \times \\
&\quad [(u_{13}u_{21} + u_{23}u_{11})v_5 + u_{23}u_{21}v_7 + u_{13}u_{11}v_4 - 2u_{33}u_{31}v_1] \\
&\quad + \xi_0 \frac{u_{11} \left(u_{11}v_3 - u_{21}v_6 - \frac{i}{\sqrt{2}}u_{31}v_2 \right)}{2(V_1 + \varepsilon_1)a} \\
&\quad - \xi_0 \frac{u_{11} \left(-u_{11}v_3 + u_{21}v_6 - \frac{i}{\sqrt{2}}u_{31}v_2 \right)}{2(V_1 - \varepsilon_1)a}.
\end{aligned}$$

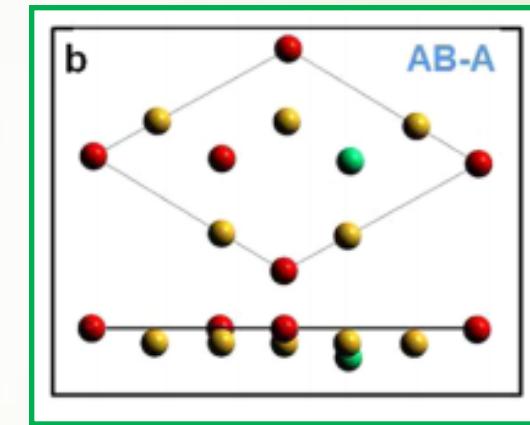
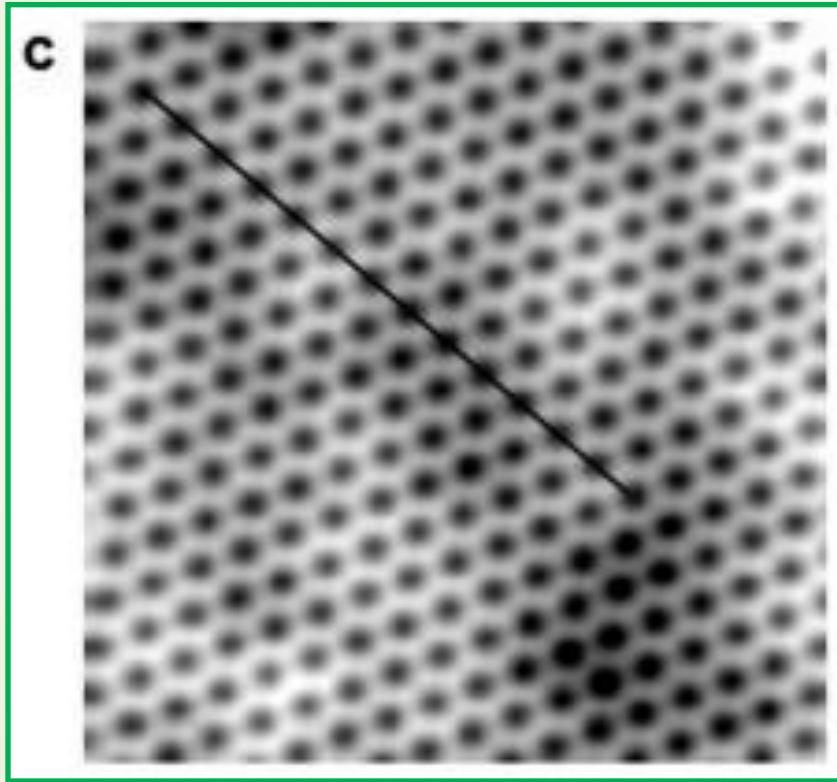
$$\begin{aligned}
\lambda_{so}^{1st} &= \frac{\xi_0}{2} \frac{\varepsilon_1^2}{\alpha_1^2 V_3^2} \approx \frac{\xi_0}{2} \frac{9}{9} \frac{\Delta^2 (V_{pp\pi} - V_{pp\sigma})^2}{V_{sp\sigma}^4} \\
&\quad \times \frac{\cot^2 \theta}{1 + \frac{\cos^2 \theta (V_{pp\pi} - V_{pp\sigma})^2}{V_{sp\sigma}^2} \left(1 + \frac{2}{9} \frac{\Delta^2}{\sin^2 \theta V_{sp\sigma}^2} \right)}.
\end{aligned}$$

Silicene: Spin-orbit interaction

System	$a(\text{\AA})$	θ	λ_{so}^{1st} (meV)	λ_{so}^{2st}	λ_R	Gap (meV) (TB)	Gap(FP)	$v_F (10^5 m/s)$ (TB)	v_F (FP)
graphene	2.46	90°	0	1.3×10^{-3}	0	2.6×10^{-3}	0.8×10^{-3} ^a	9.80	8.46
silicene	3.86	101.7°	3.9	7.3×10^{-2}	0.7	7.9	1.55 ^b	5.52	5.42
ge(licene)	4.02	106.5°	43	3.3	10.7	93	23.9 ^b	4.57	5.24
sn(licene)	4.70	107.1°	29.9	34.5	9.5	129	73.5	4.85	4.70



Evidence for Dirac fermions: Silicene



Silicene: Topological phase transition

$$\begin{aligned}
 H_0 = & -t \sum_{\langle i,j \rangle \alpha} c_{i\alpha}^\dagger c_{j\alpha} + i \frac{\lambda_{SO}}{3\sqrt{3}} \sum_{\langle\langle i,j \rangle\rangle \alpha\beta} \nu_{ij} c_{i\alpha}^\dagger \sigma_{\alpha\beta}^z c_{j\beta} \\
 & - i \frac{2}{3} \lambda_{R2} \sum_{\langle\langle i,j \rangle\rangle \alpha\beta} \mu_i c_{i\alpha}^\dagger (\boldsymbol{\sigma} \times \hat{\mathbf{d}}_{ij})_{\alpha\beta}^z c_{j\beta},
 \end{aligned}$$

$\nu_{ij} = +1$ if the second-nearest-neighboring hopping is anticlockwise and -1 if it is clockwise with respect to the positive z axis. The third term represents the Rashba SO coupling with $\lambda_R = 0.7\text{meV}$, where $\mu_i = \pm 1$ for the A (B) site, and $d_{ij} = d_{ij}/|d_{ij}|$ with d_{ij} the vector connecting two sites i and j in the same sublattice.

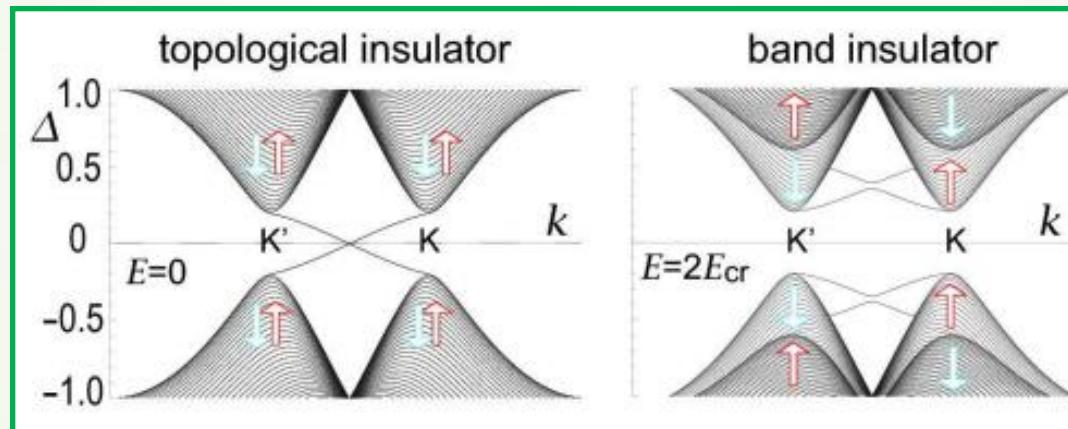
$$\ell = 0.23\text{\AA}$$

$$\begin{aligned}
 H_E = & i \lambda_{RI}(E_z) \sum_{\langle i,j \rangle \alpha\beta} c_{i\alpha}^\dagger (\boldsymbol{\sigma} \times \hat{\mathbf{d}}_{ij})_{\alpha\beta}^z c_{j\beta} \\
 & + \ell \sum_{i\alpha} \mu_i E_z^i c_{i\alpha}^\dagger c_{i\alpha},
 \end{aligned}$$

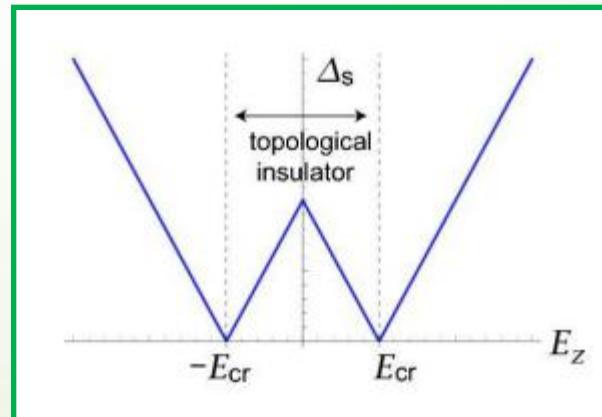
$\lambda_{RI}(E_z) \propto E_z$
 $10\mu\text{eV at } E_z = \lambda_{SO}/\ell = 17\text{meV\AA}^{-1}$

Silicene: Topological phase transition

$$\Delta_s(E_z) = -s\lambda_{SO} + \frac{1}{2}\ell E_z + \frac{1}{2}\sqrt{(\ell E_z)^2 + \lambda_{RI}^2}.$$



$$E_{cr} = \frac{s\lambda_{SO}}{\ell} \left[1 - \left(\frac{\lambda_{RI}}{2\lambda_{SO}} \right)^2 \right] = \pm 17 \text{ meV/\AA},$$



Specific examples

Transition metal dichalcogenides

Transition metal dichalcogenides

From insulator to superconductor

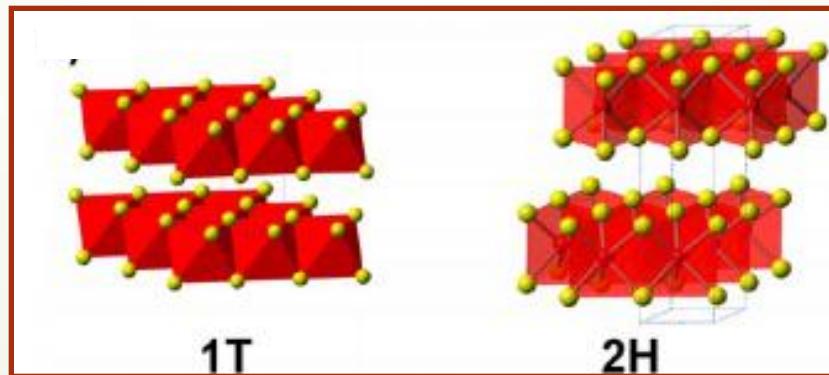
Include the semiconductivity, half-metallic magnetism, superconductivity, or charge density wave

Applications:

lubrication, catalysis, photovoltaics, supercapacitors, and rechargeable battery systems.

Examples: Layered van der Waals solids

- Graphene
- Silicene, Germanium, Tin
- Transition metal dichalcogenid materials , MX_2 ($\text{M} = \text{Ti}, \text{Zr}, \text{Hf}, \text{V}, \text{Nb}, \text{Ta}, \text{Re}; \text{X} = \text{S}, \text{Se}, \text{Te}$)



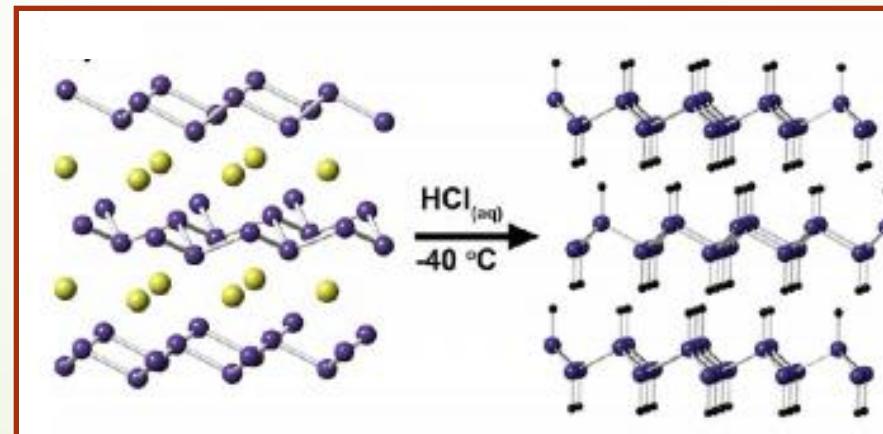
Mattheis, Phys. Rev. B **8**, 3719 (1973)
Helveg, et al Phys. Rev. Lett. **84**, 951 (2000)

Wang, Kalantar Zadeh, Kis, Coleman, Strano, Nature Nanotech. **7**, 699 (2012)

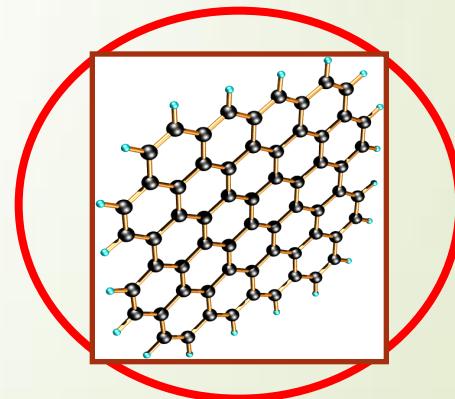
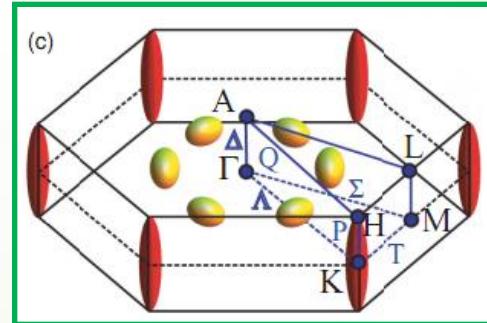
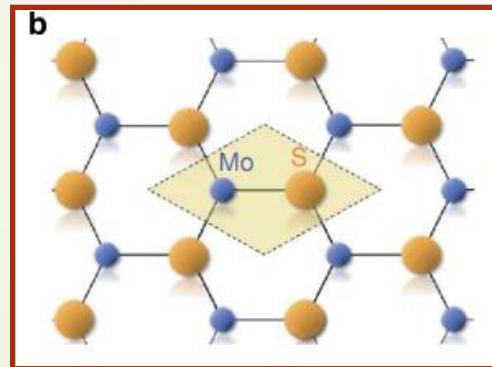
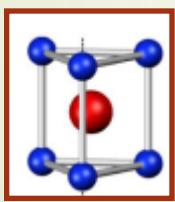
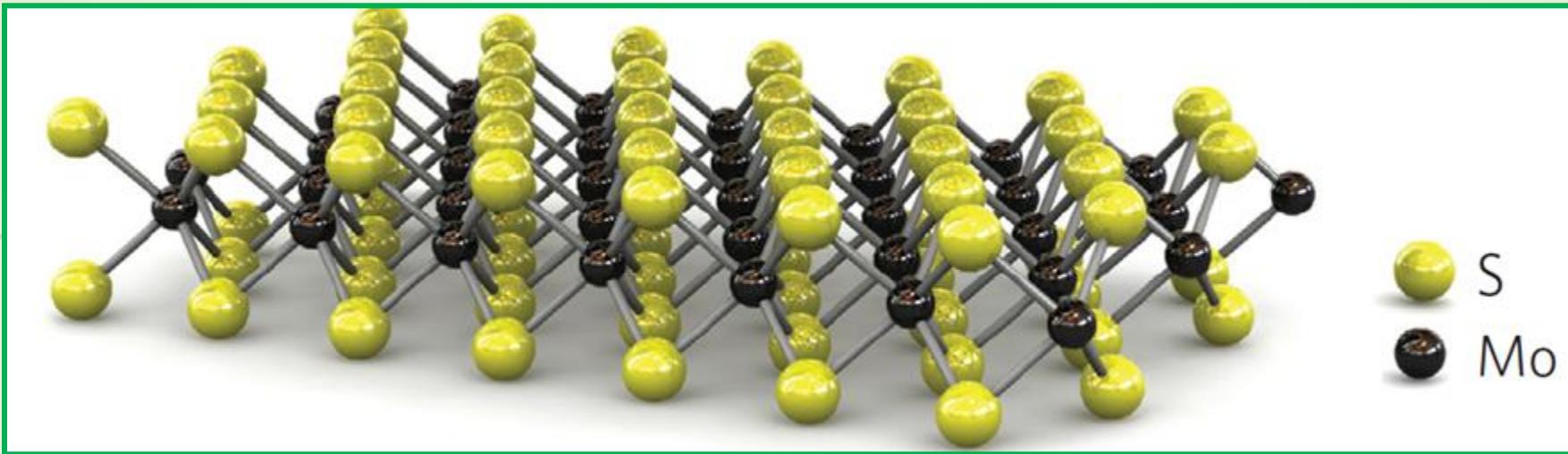
Examples: Layered van der Waals solids

Many novel van der Waals compounds can be created via precursor solids.

For example, the layered **CaGe** can be layered in aqueous HCl to produce monolayer



TMDCs: *MoS₂* crystal

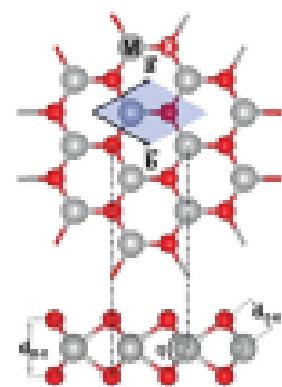


Wilson and Yaffe, Adv. Phys. **18**, 193 (1969)
 Romley, Murray, Yoffe, J. Phys. C **5** (1972)
 Mattheis, Phys. Rev. B **8**, 3719 (1973)
 Helveg, *et al.*, Phys. Rev. Lett. **84**, 951 (2000)

kobayashi, Yamauchi, Phys. Rev. B **51**, 17085 (1995)

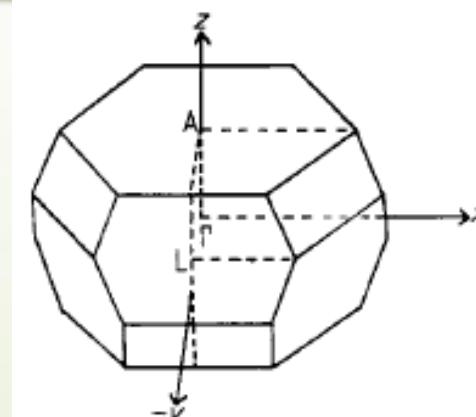
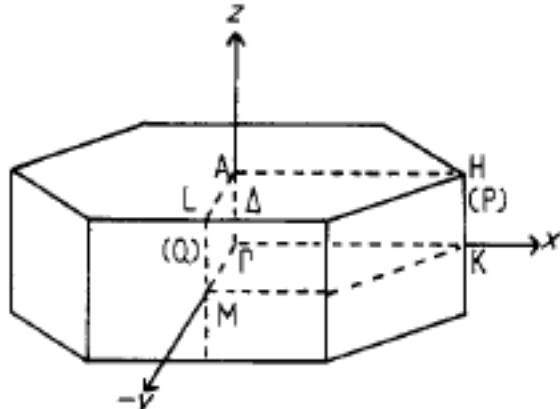
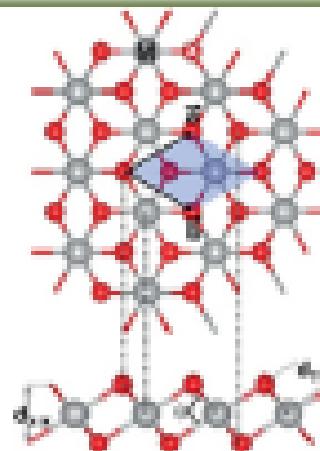
Transition metal dichalcogenides

2H-MX₂

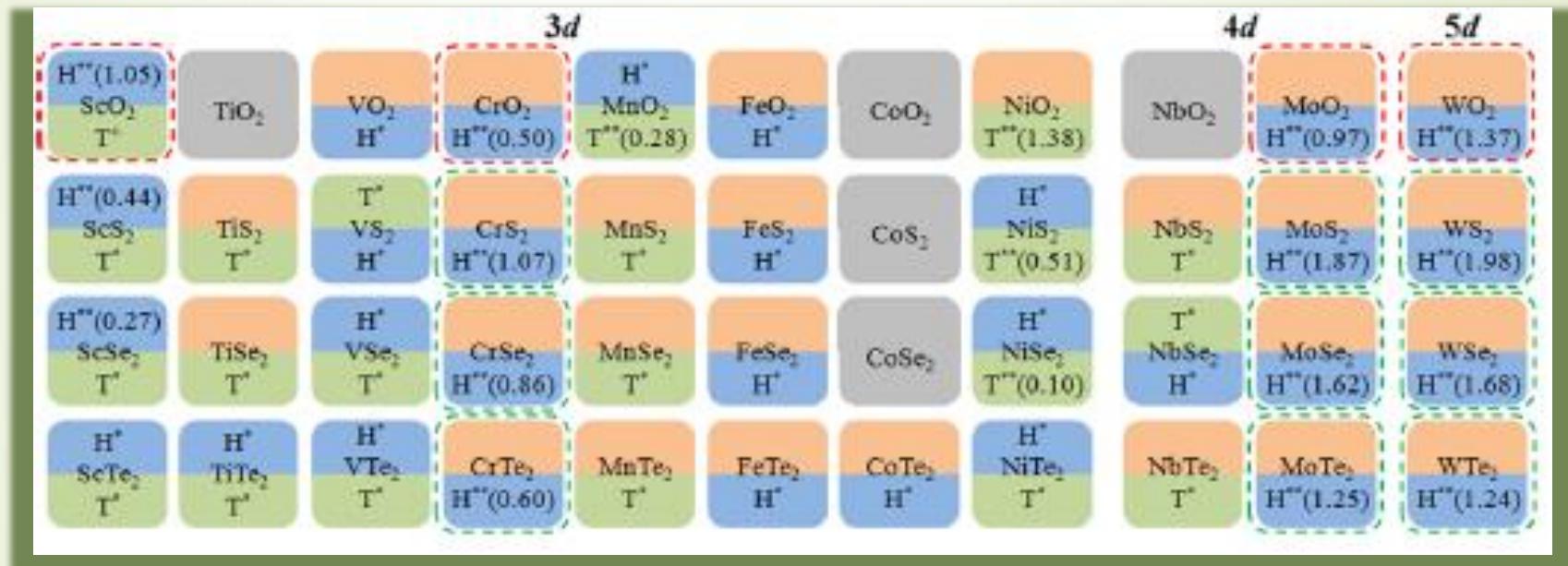


Honeycomb (H) structure

1T-MX₂



Transition metal dichalcogenides



H & T
unstable

H stable

T stable

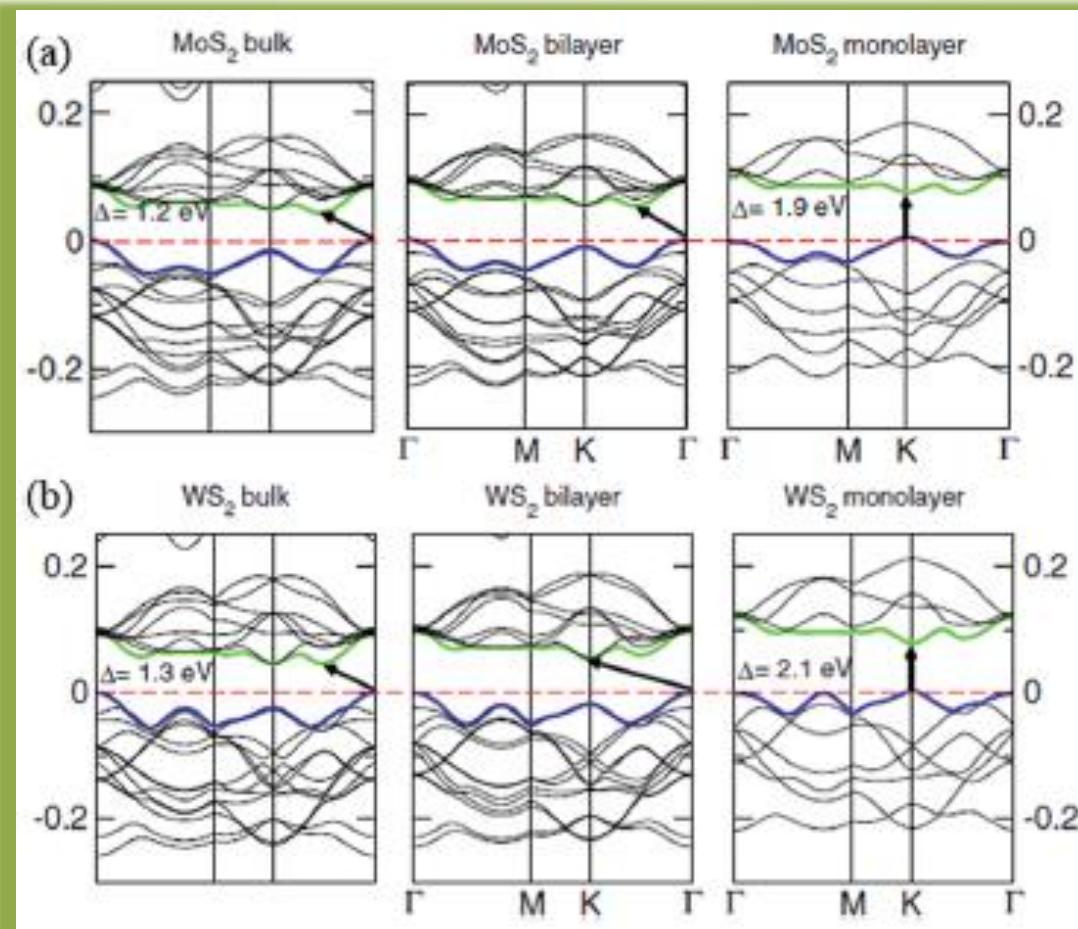
H & T
Stable
 $E_c[T] > E_c[H]$

H & T
Stable
 $E_c[T] < E_c[H]$

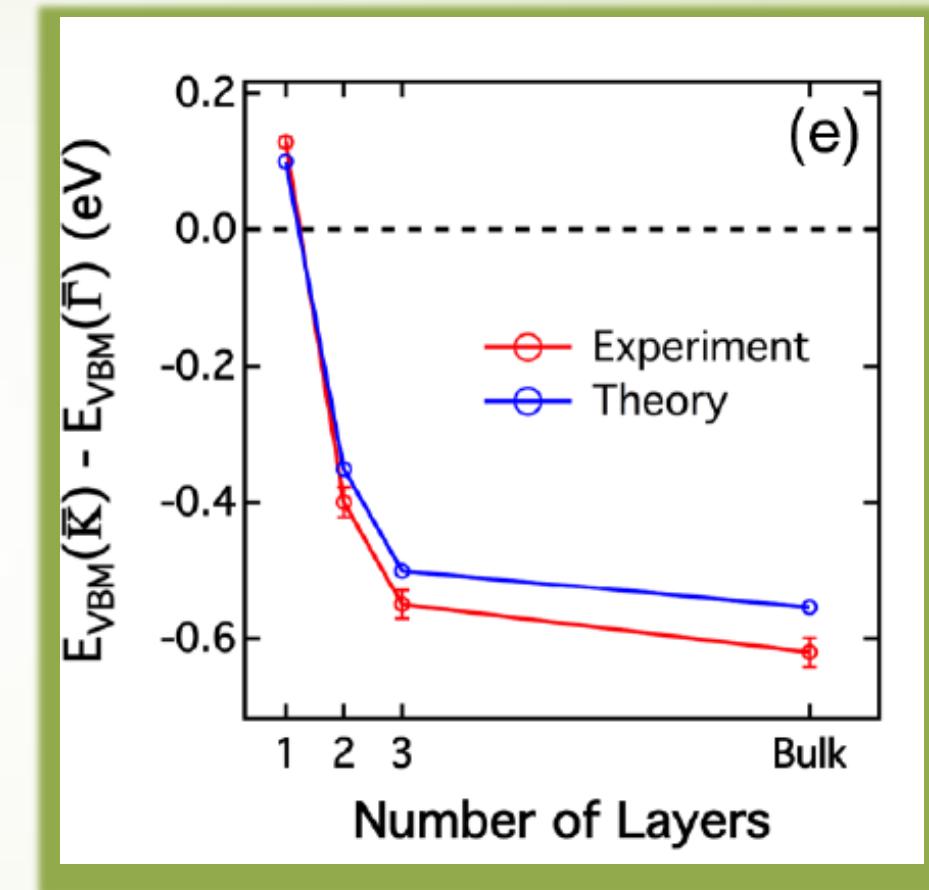
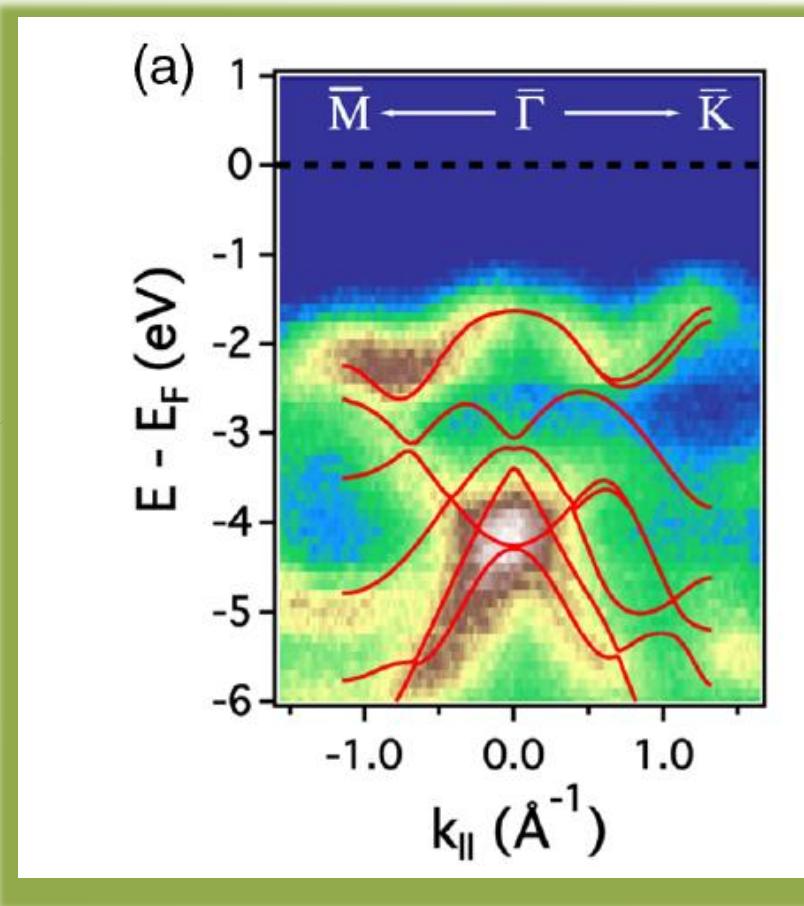
E_c : cohesive energy per MX_2 unit
 T' : half-metal; T' & H' : metal
 T'' & H'' : semiconductor (E_g /eV)

direct band gap
 indirect band gap

Band structures



ARPES measurements of the electronic structure



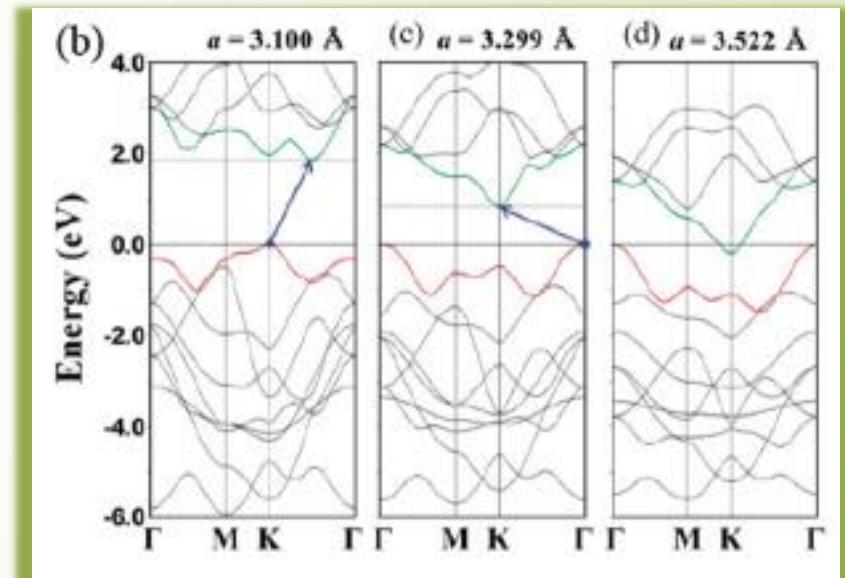
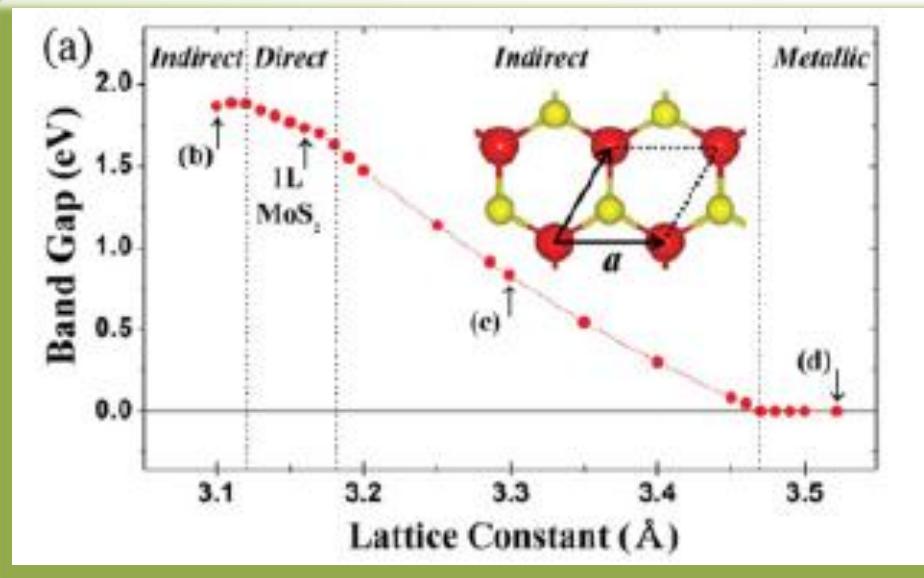
W. Jin, *et al*, Phys. Rev. Lett. **111**, 106801 (2013)

E. Cappelluti, *et al* Phys. Rev. B **88**, 075409 (2013).

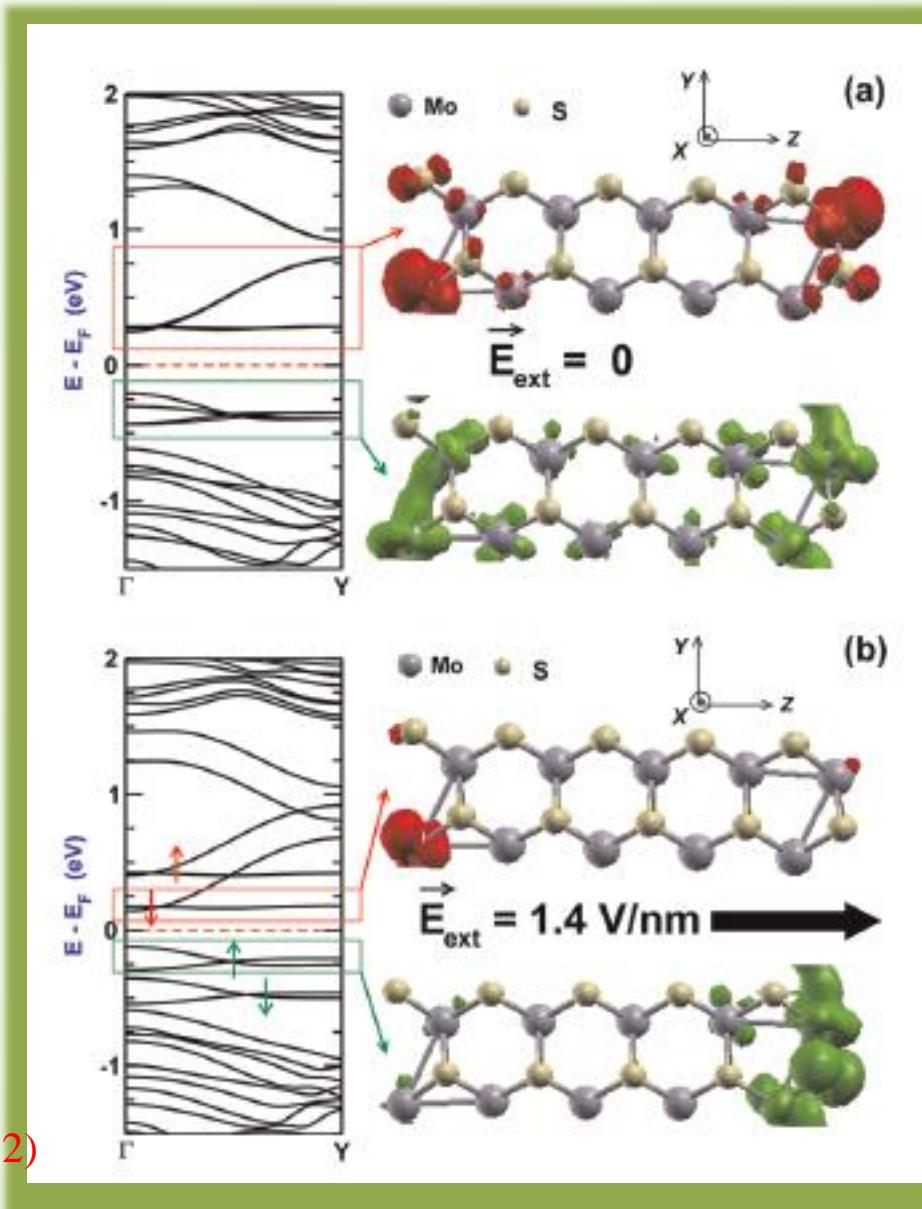
Strain effects on the electronic properties of TMDs

Strain in a crystalline solid
modifies the lattice constants and reduces the
crystal symmetry, leading to significant shifts in the energy band edges

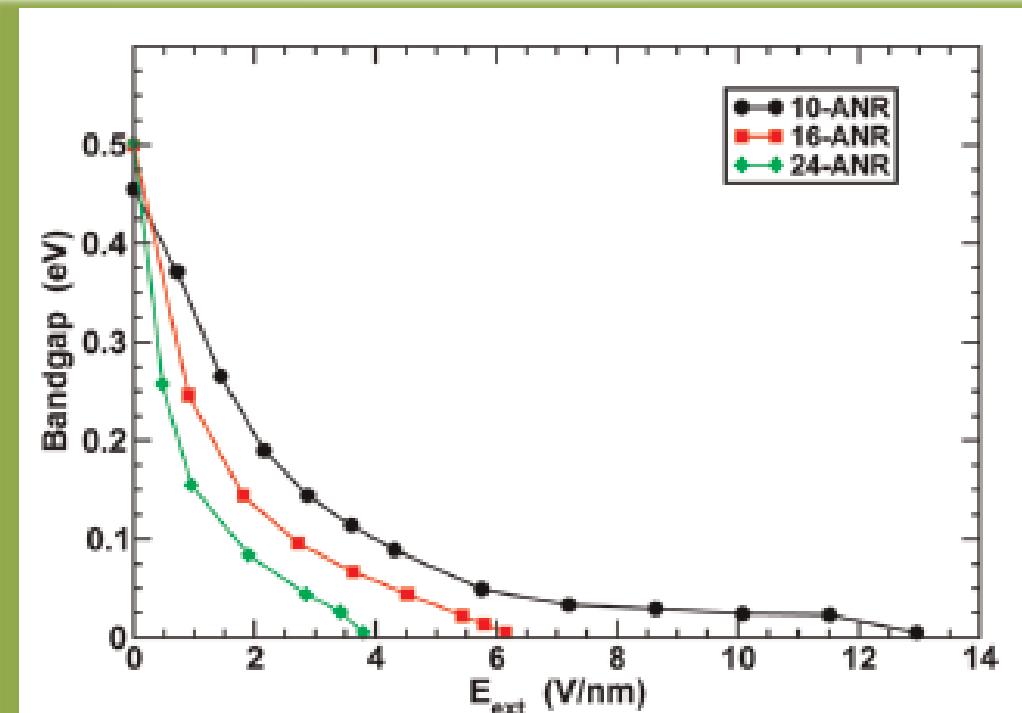
Strain effects on the electronic properties of TMDs



Electric field on Armchair MoS₂



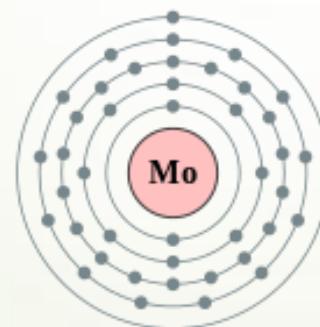
Band gap in the applied transverse electric field



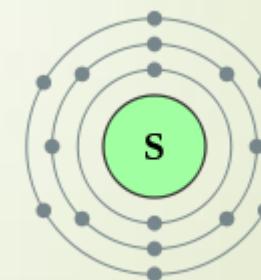
TMDCs: MoS₂ crystal

Photoluminescence properties

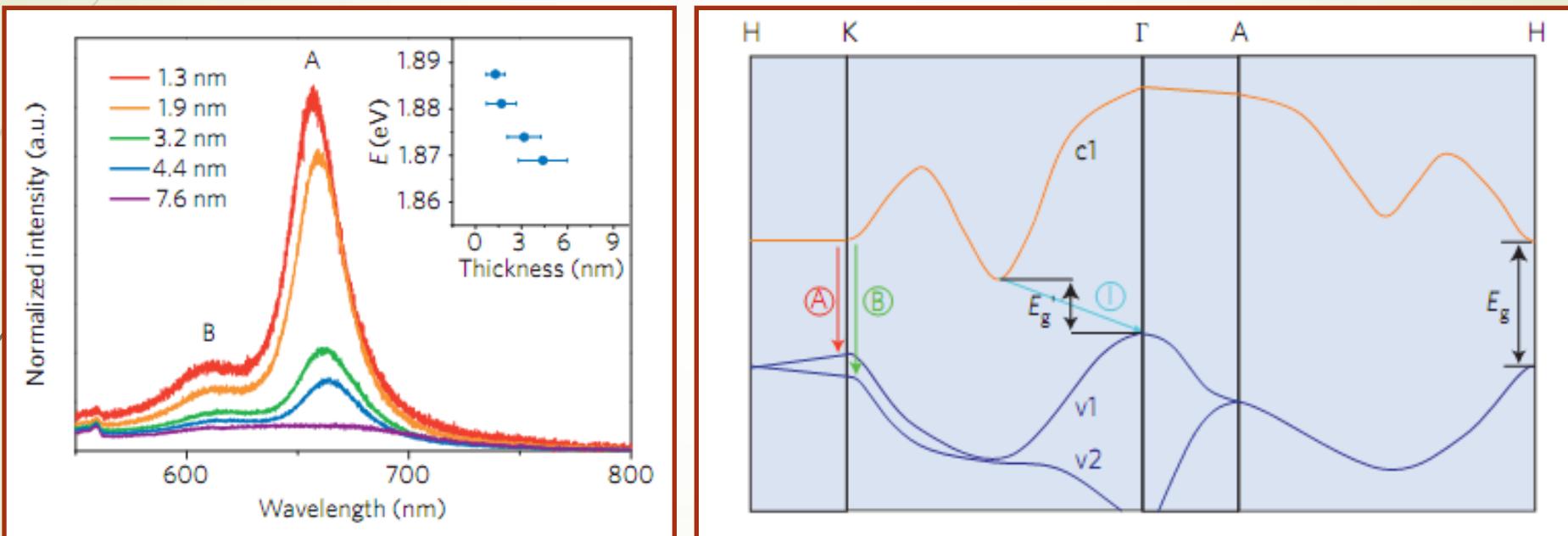
[Kr] 5s¹ 4s²p⁶d⁵
2 ,8 ,18 ,13 ,1



[Ne] 3s² 3p⁴
2 ,8 ,6



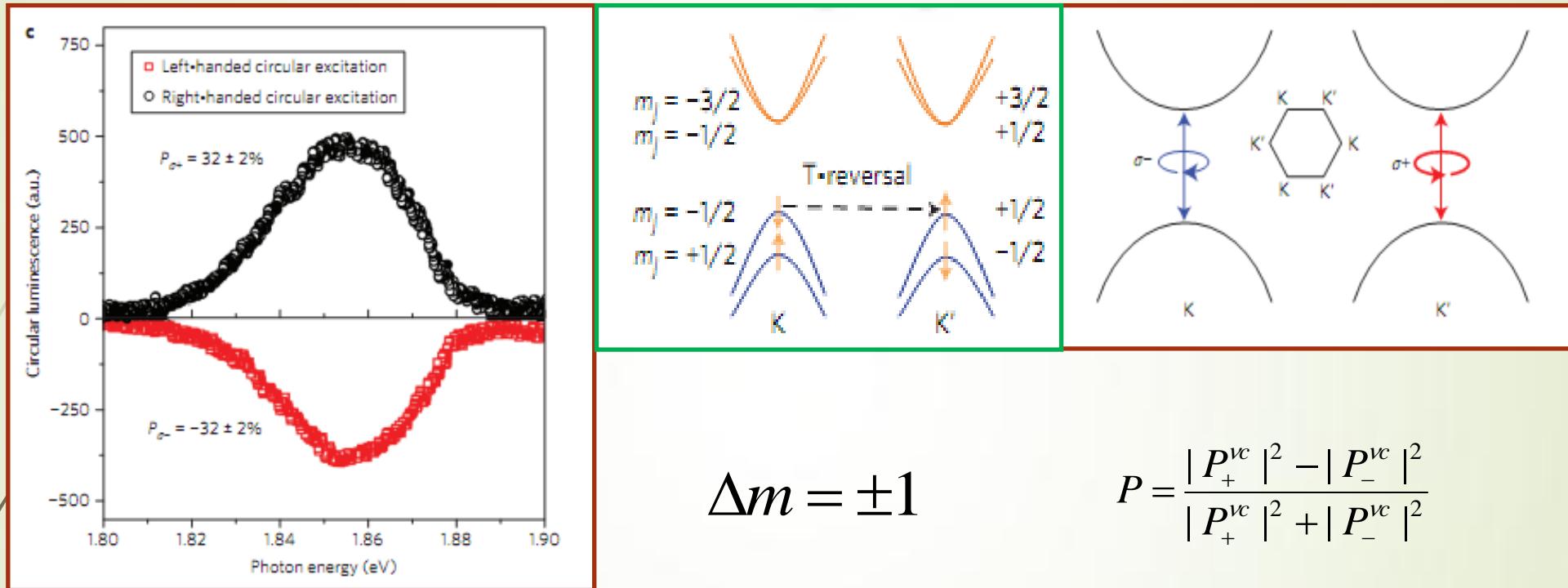
Spectroscopy: few layers



Eda, *et al*, Nano Lett. **11**, 5111 (2011)

Mak *et al*, Phys. Rev. Lett. **105**, 136805 (2010)

Valley polarization



$$\Delta m = \pm 1$$

$$P = \frac{|P_+^{vc}|^2 - |P_-^{vc}|^2}{|P_+^{vc}|^2 + |P_-^{vc}|^2}$$

$$P_{\pm}^{vc}(k) = \langle u_c(k) | p_x + ip_y | u_v(k) \rangle$$

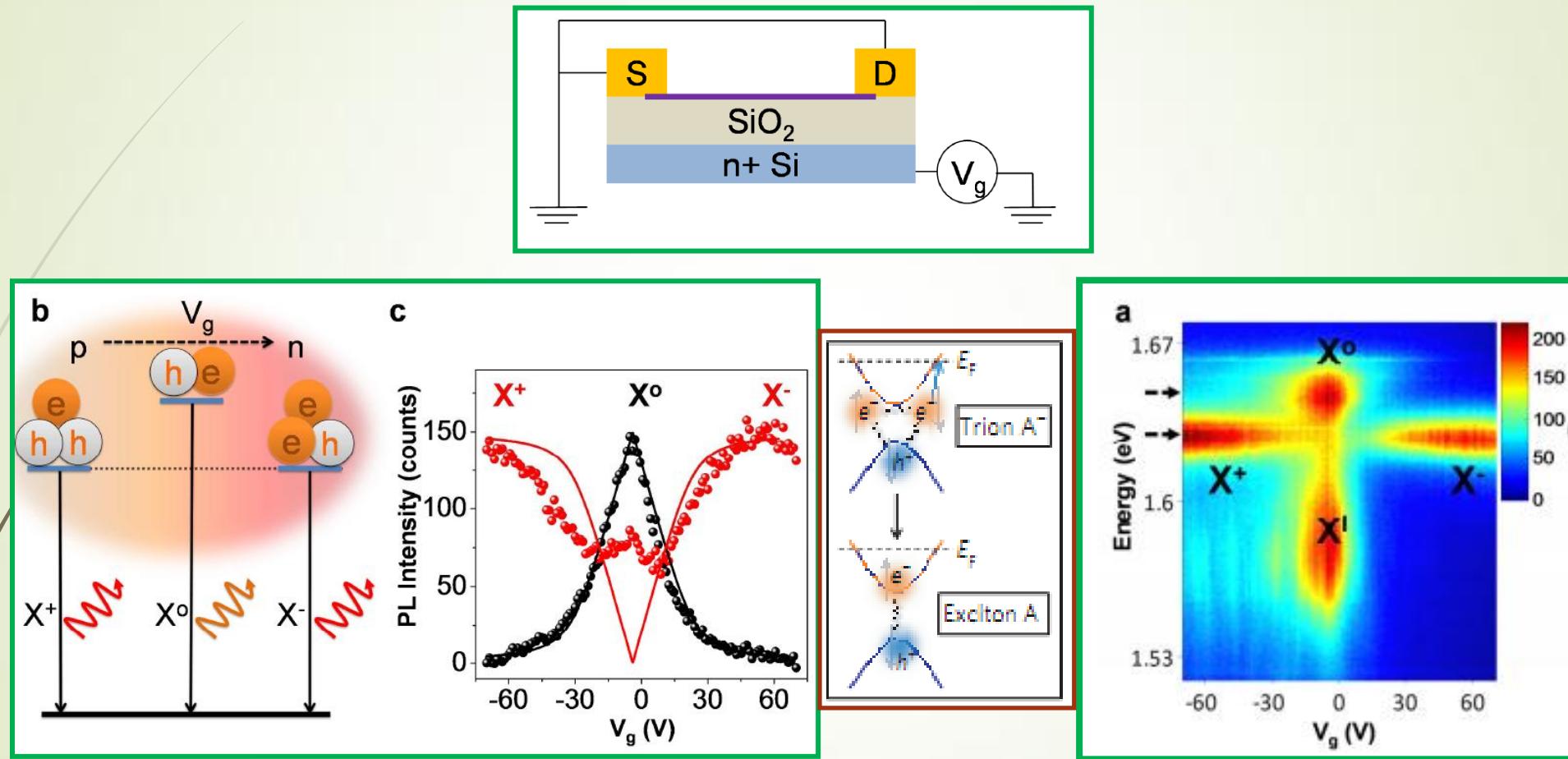
$$|P_{\pm}^{vc}(k)|^2 \propto (1 \pm \tau \frac{\Delta'}{\sqrt{\Delta'^2 + 4a^2 t^2 k^2}})^2$$

$$\Delta' = \Delta - \tau s \lambda$$

Zeng, et al, Nature Nano 7, 490 (2012)

T. Cao, et al Nature Commun. DoI: 10.1038/ncomms1882 (2012)

Charged excitons (MoSe₂: Trion)



K. F. Mak *et al*, Nature Materials **12**, 207 (2013)

J. S. Ross *et al*, Nature Communi. **4**, 1474(2013)

TMDCs: MoS₂ crystal

Nanoelectronic properties

Scattering mechanisms

- Charged impurities(Coulomb scattering)
- Neutral defects (short range scattering)
- Surface interface phonon scattering
- Ripples and roughness scattering
- Acoustic and optical phonons scattering

2D
Semiconductor

$$\mu_{charged} \approx (m_e^*)^{-1/2} T^{3/2}$$

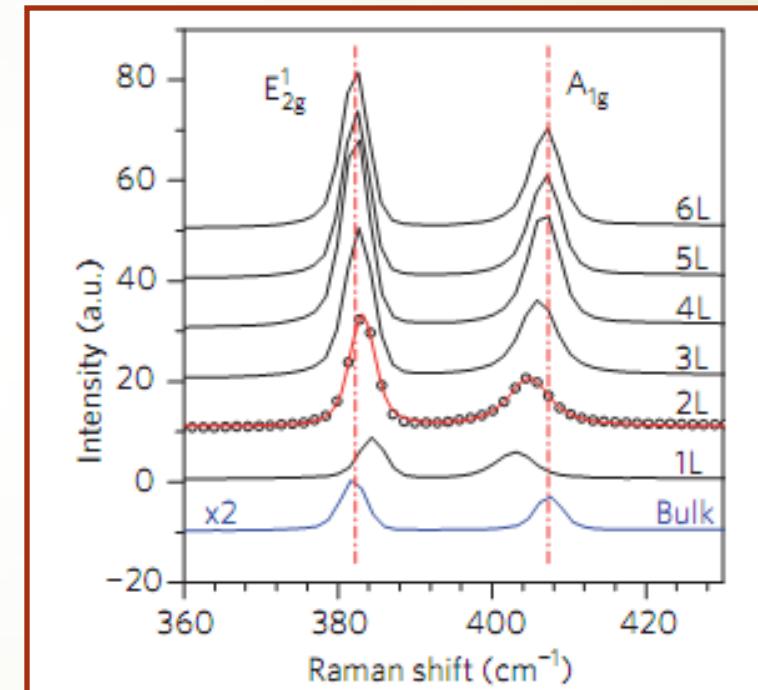
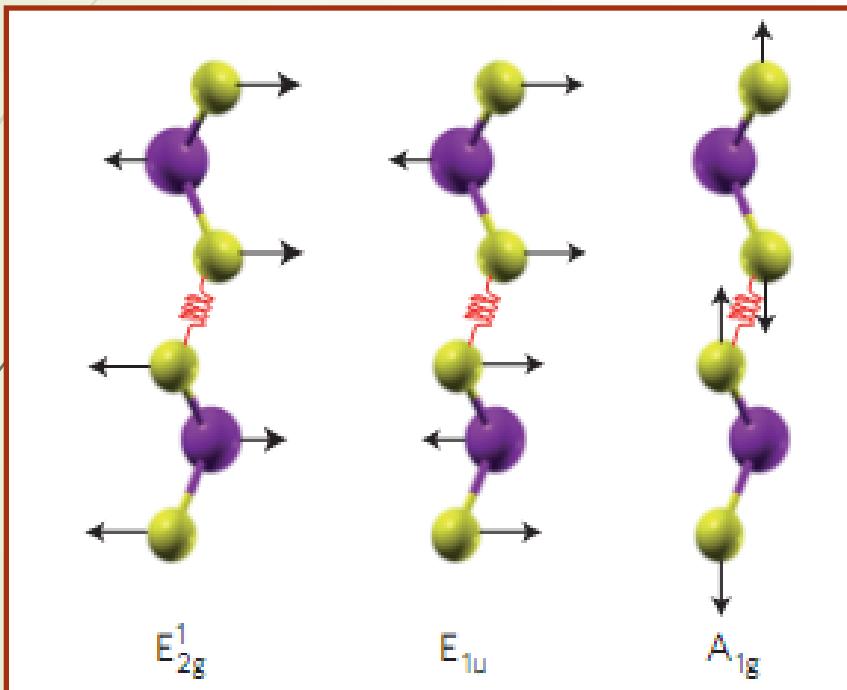
$$\mu_{optical} \approx (m_e^*)^{-5/2} T^{-1} \left[e^{\frac{\omega_{op}}{T}} - 1 \right]$$

$$\mu_{acoustic} \approx (m_e^*)^{-5/2} T^{-3/2}$$

$$\mu(T) \approx \begin{cases} T \ll 1 & 10^4 \quad (1978) \\ T \ll 1 & 3 \times 10^6 \quad (2000) \\ T = 300 & 1000 \quad (2000) \end{cases}$$

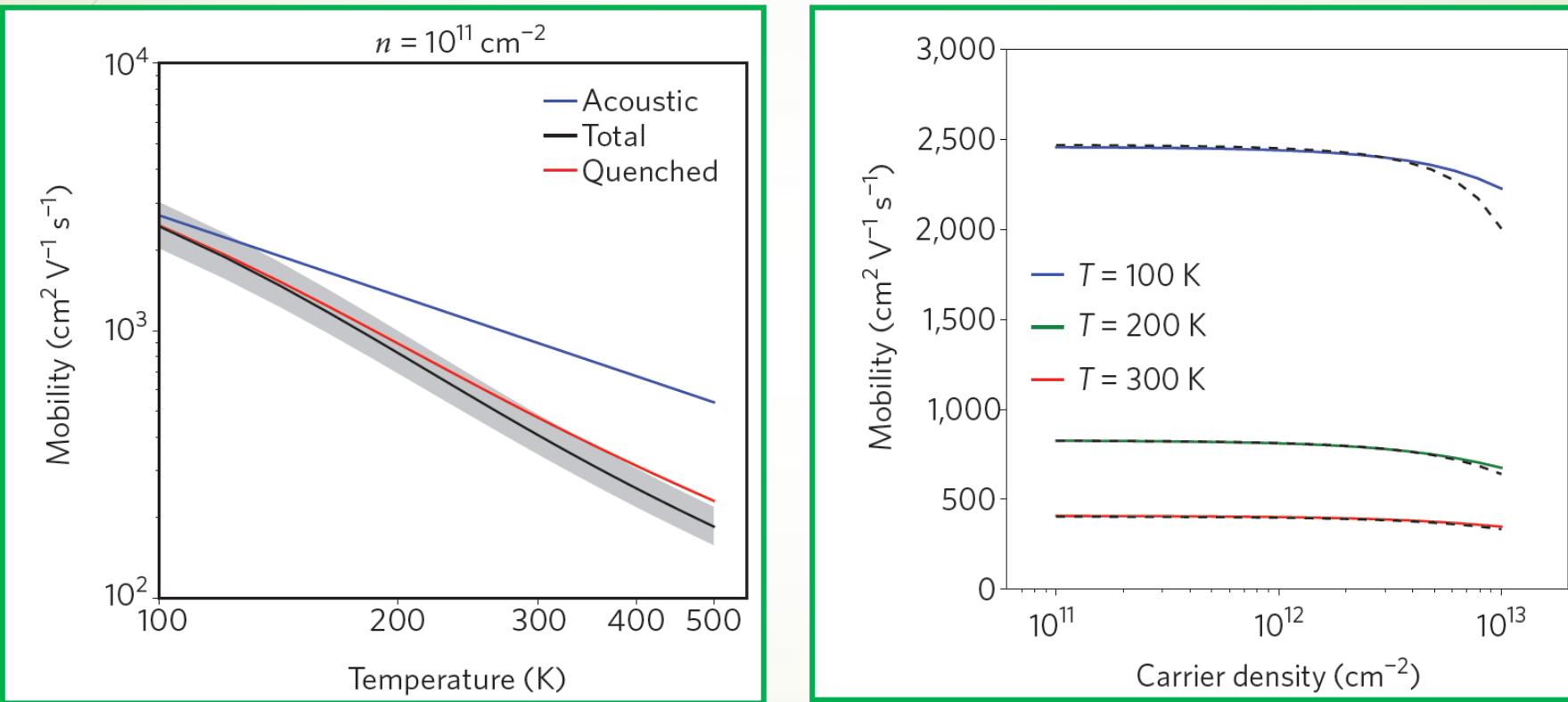
Typical electron mobility for Si at room temperature (300 K) is $1400 \text{ cm}^2/(\text{V}\cdot\text{s})$ and the hole mobility is around $450 \text{ cm}^2/(\text{V}\cdot\text{s})$.

Phonons, Raman Spectroscopy



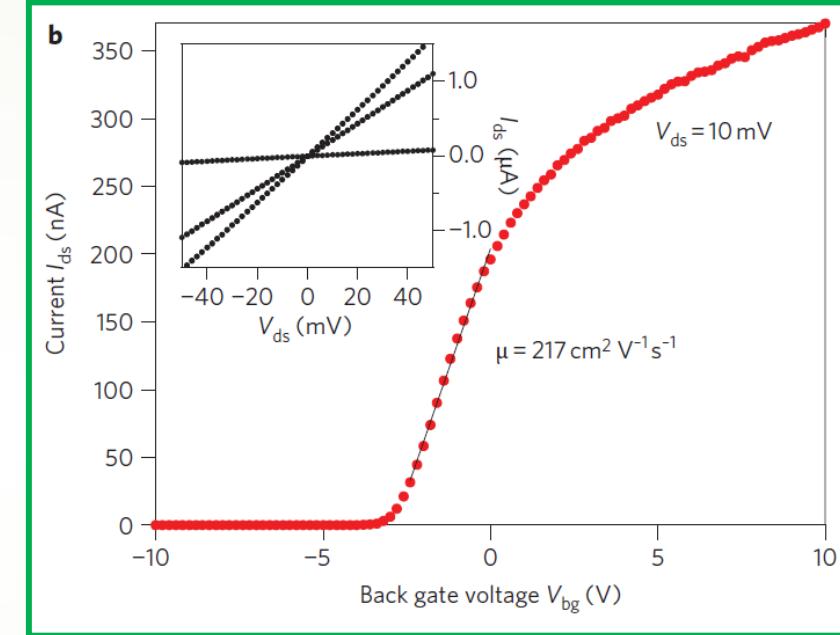
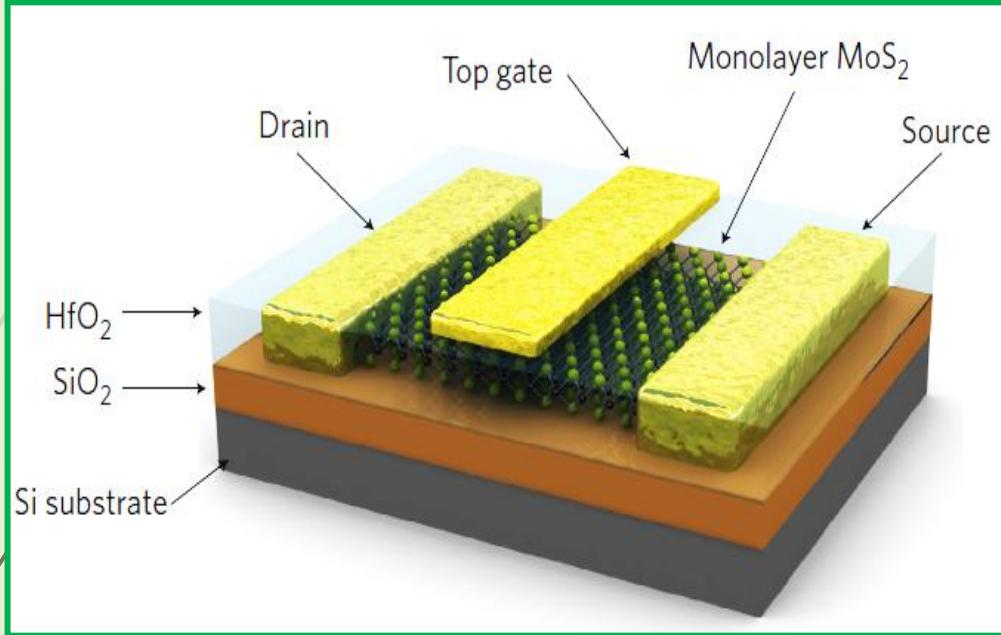
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Mobility: evidence



MoS₂ Transistor

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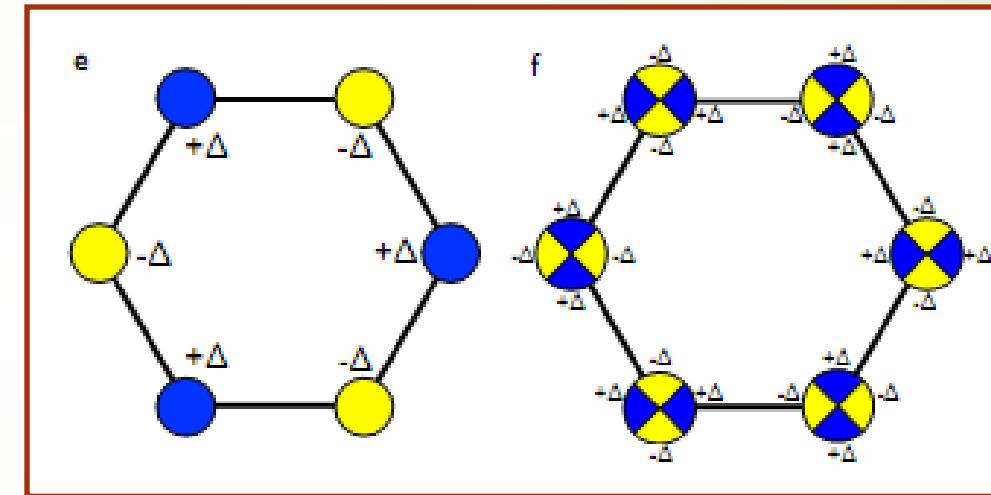
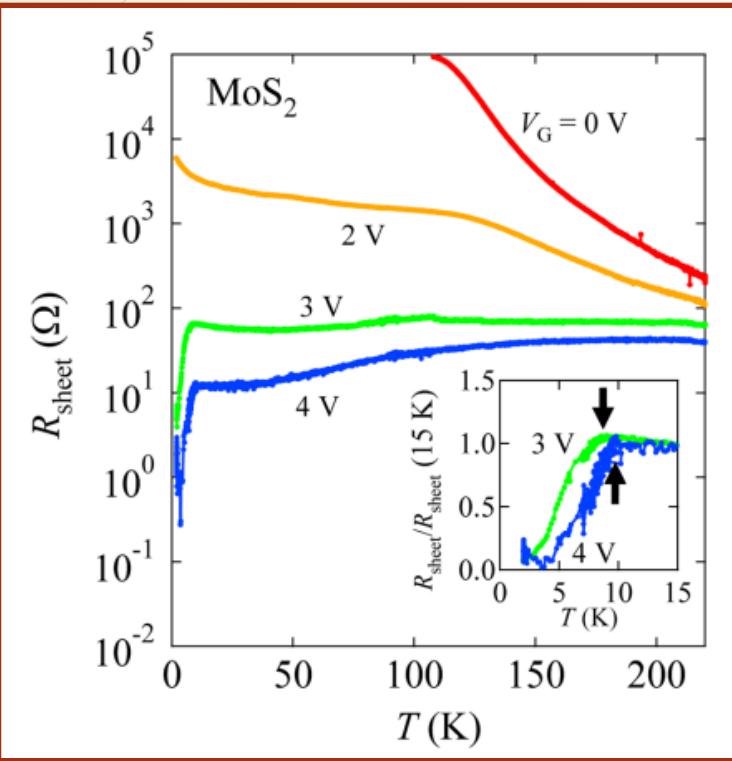
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Superconductivity

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Thanks for your attention

