

Andreev Reflection in a Monolayer MoS₂

Leyla Majidi

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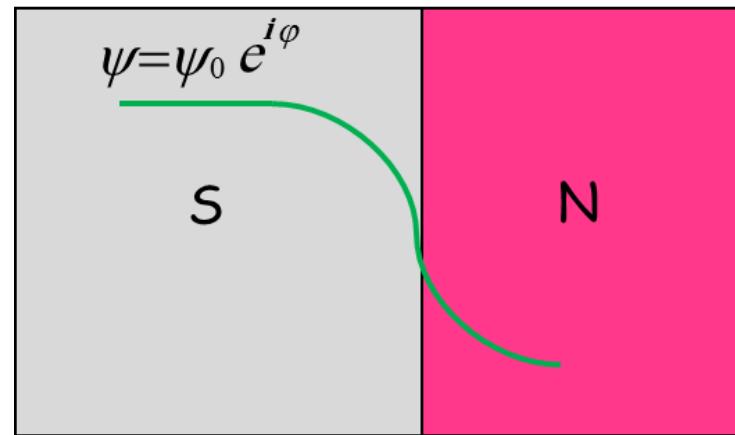
Habib Rostami and Dr. Reza Asgari

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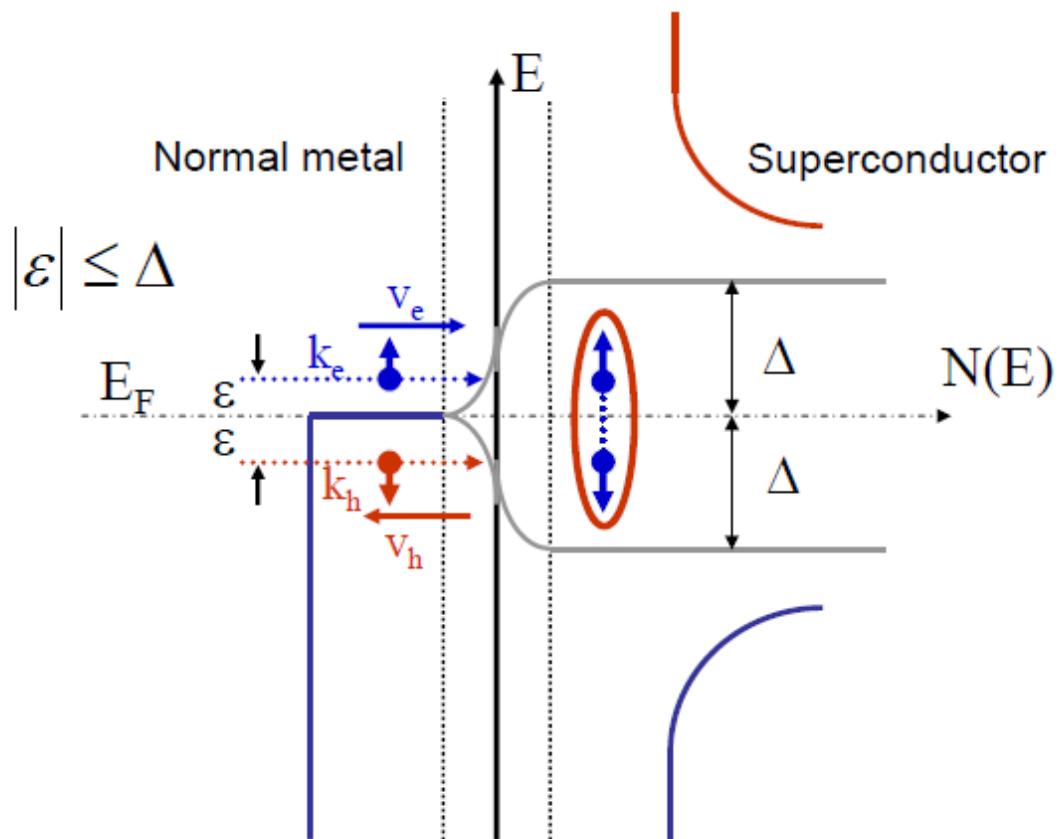
Workshop on Electronic Properties of MoS₂

IPM 2013, October 2

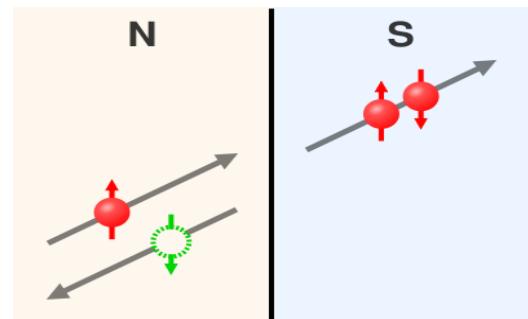
Superconducting Proximity Effect



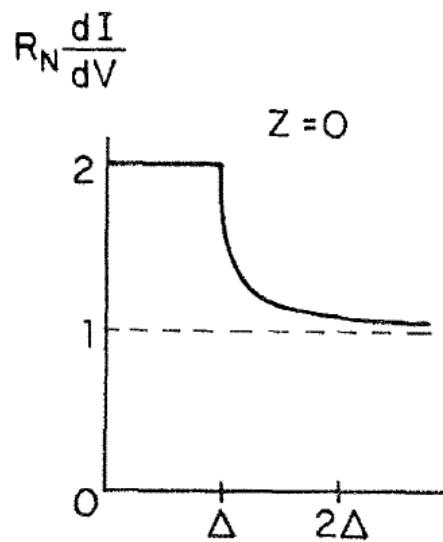
Andreev Reflection in S/N Junctions



Retro Andreev Reflection



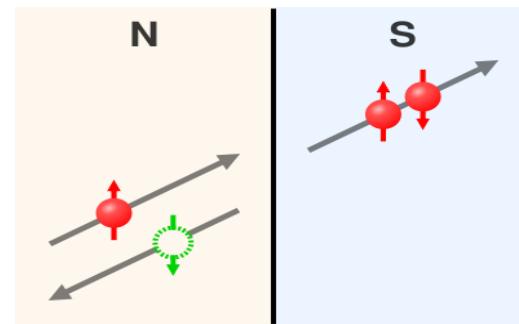
Andreev Reflection in S/N Junctions



Differential conductance vs voltage
at $T=0$

BTK, PRB 25, 4515 (1982)

Retro Andreev Reflection



Conversion of **normal current** to **supercurrent** :
charge transfer by $2e$

Andreev Reflection in F/S Junctions

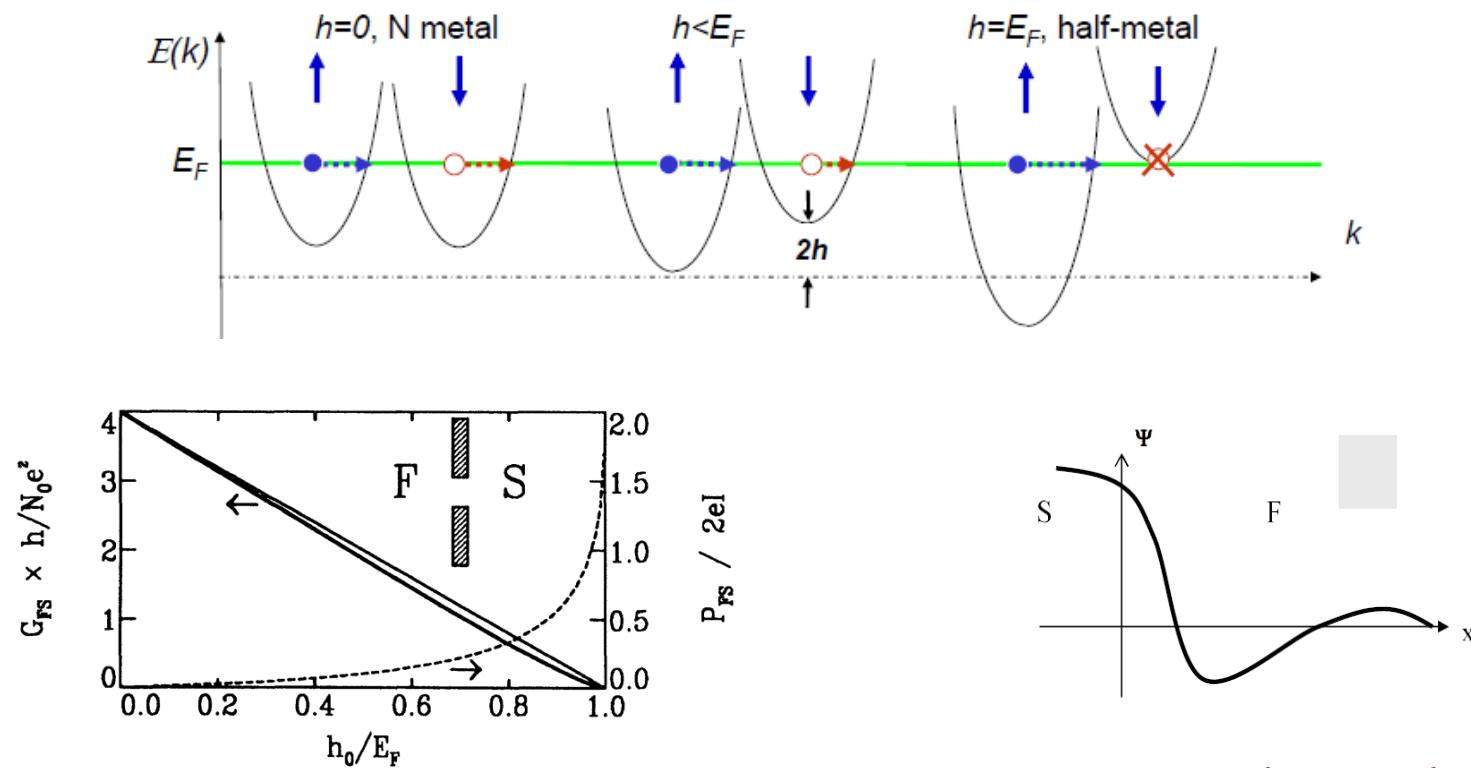


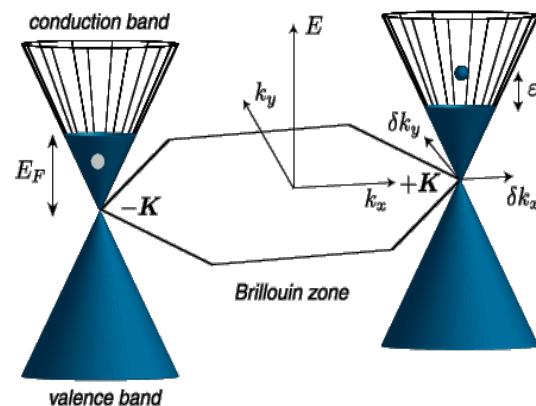
FIG. 1. The conductance G_{FS} (full curves) and the shot-noise power P_{FS} (dashed) of a ballistic point contact in a ferromagnet-superconductor junction (see inset), as a function of the exchange energy h_0 . The thick line represents the exact result (7) for G_{FS} , the thin line the estimation (2).

Ferromagnetism and singlet-spin superconductivity are exclusive phenomena

Andreev Reflection in Atomically Thin 2D Crystals:

Graphene, Transition-Metal Dichalcogenids (TDMs), ...

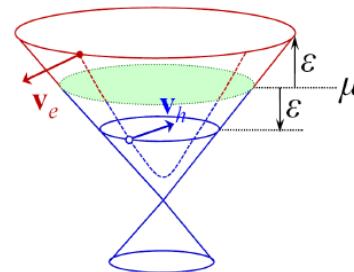
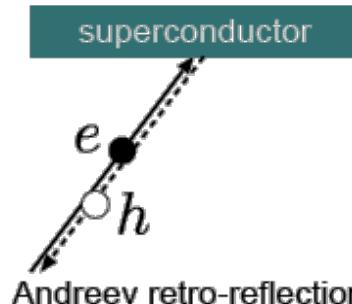
Andreev Reflection in Graphene S/N Junction



Superconducting correlation of massless Dirac electrons with opposite spin and different valley

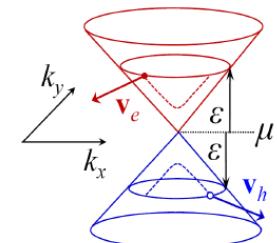
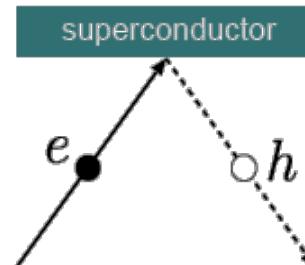
Highly doped Graphene

$$E_F \gg \Delta$$



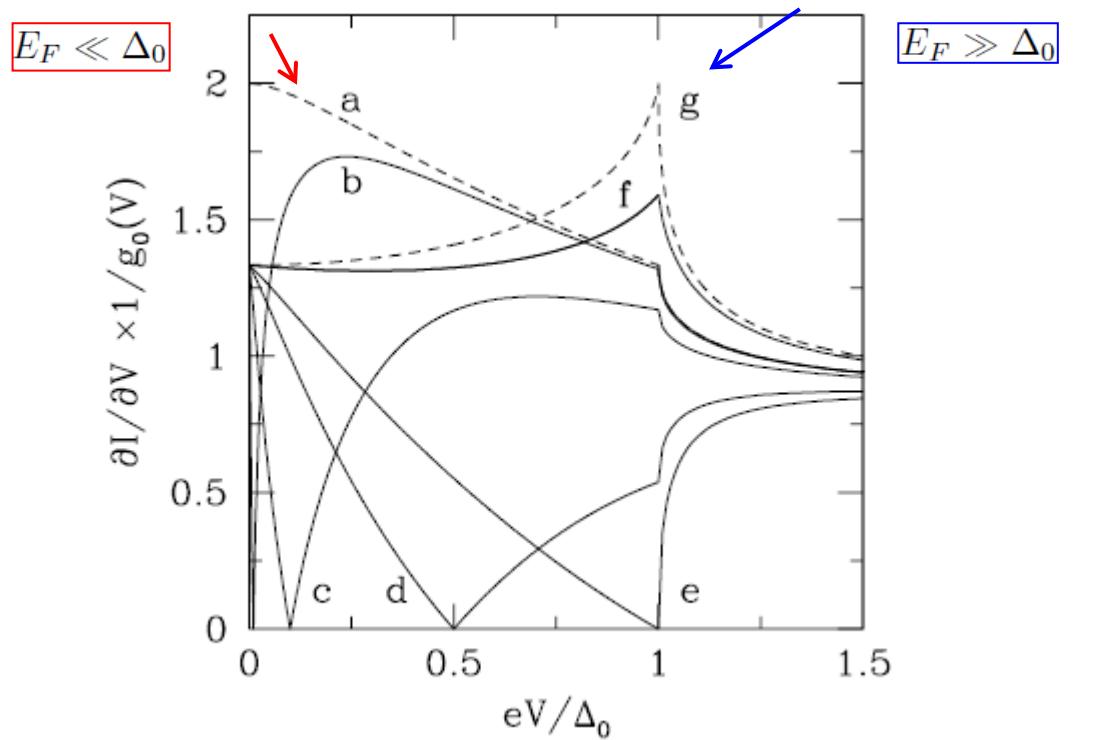
Undoped Graphene

$$E_F \ll \Delta$$



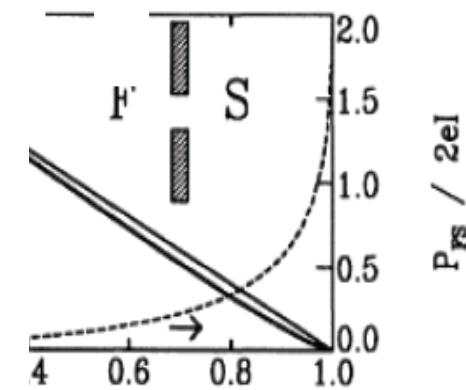
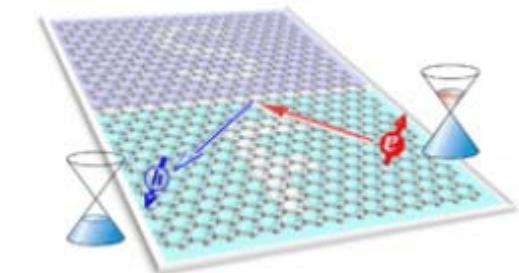
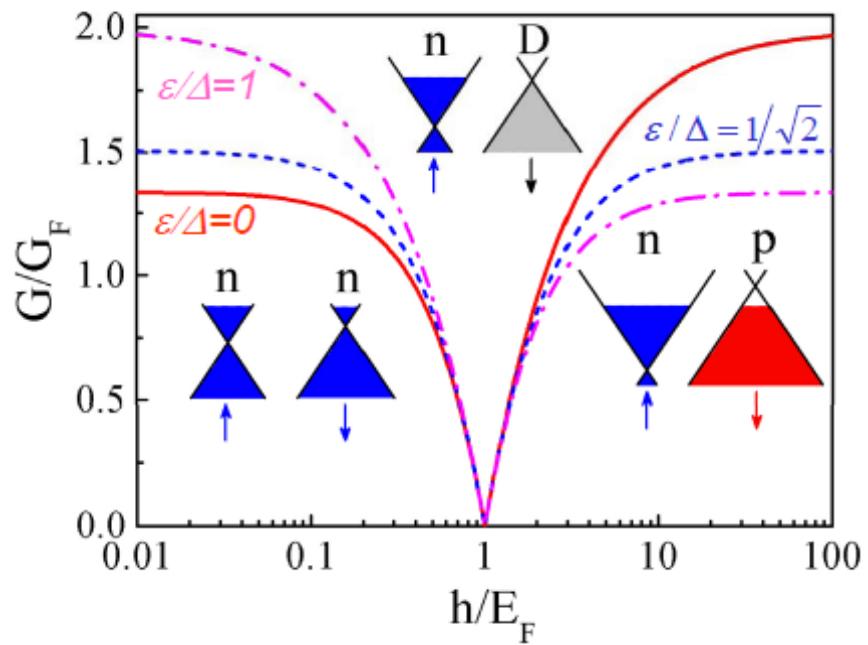
Andreev Reflection in Graphene S/N Junction

Specular reflection



Differential conductance of the NS interface for $E_F/\Delta_S = 0.01, 0.1, 0.5, 1, 10$ (solid curves labeled b,c,d,e,f, respectively). The dashed curves are the asymptotes for $E_F/\Delta_S \rightarrow 0, \infty$ (labeled a,g, respectively).

Andreev Reflection in Graphene S/F Junction

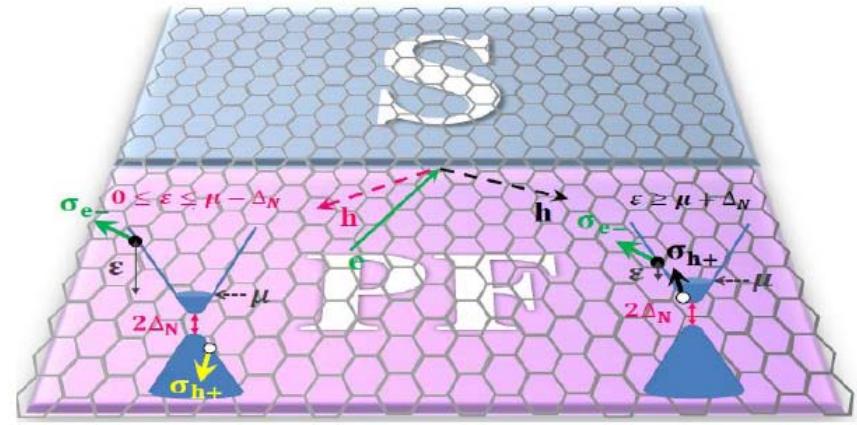
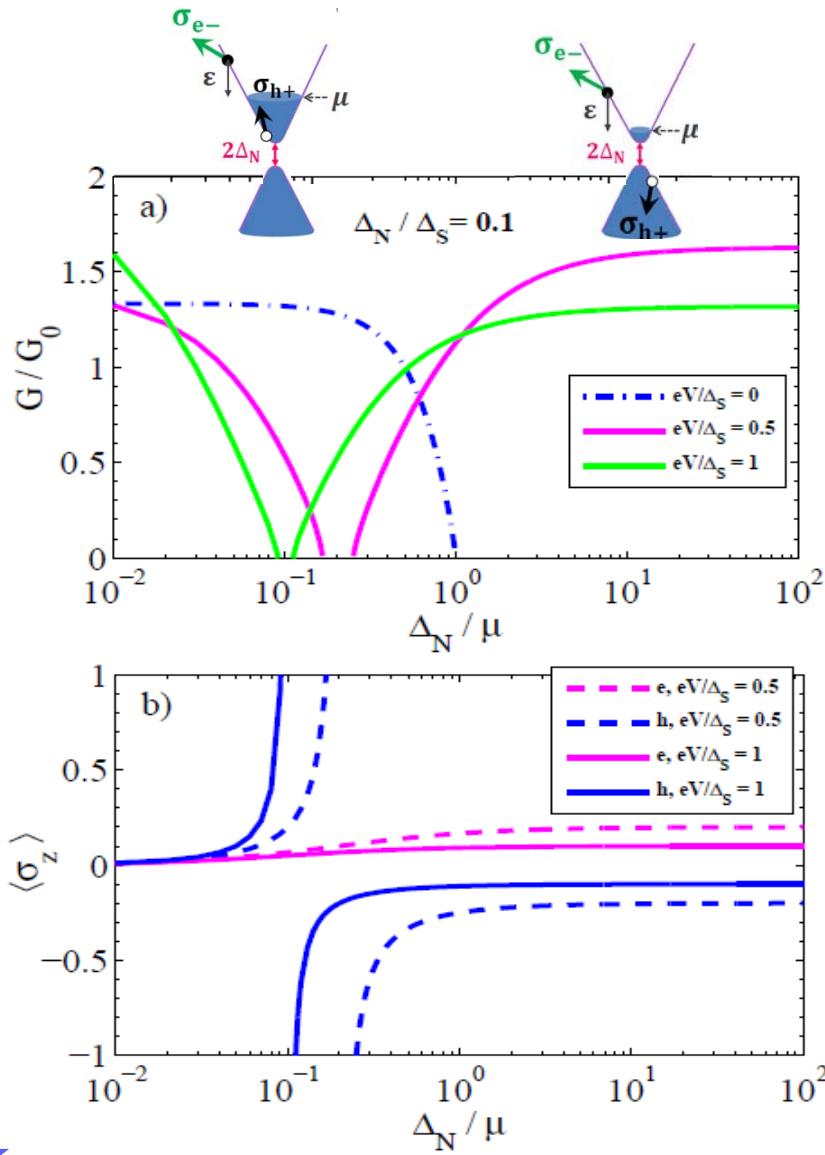


g and C. W. J. Beenakker,

for a highly doped F, $E_F \gg \Delta$ ys. Rev. Lett. **74**, 1657 (1995).

M. Zareyan, H. Mohammadpour, and A. G. Moghaddam,
Phys. Rev. B 78, 193406 (2008)

Andreev Reflection in Graphene S/PF Junction

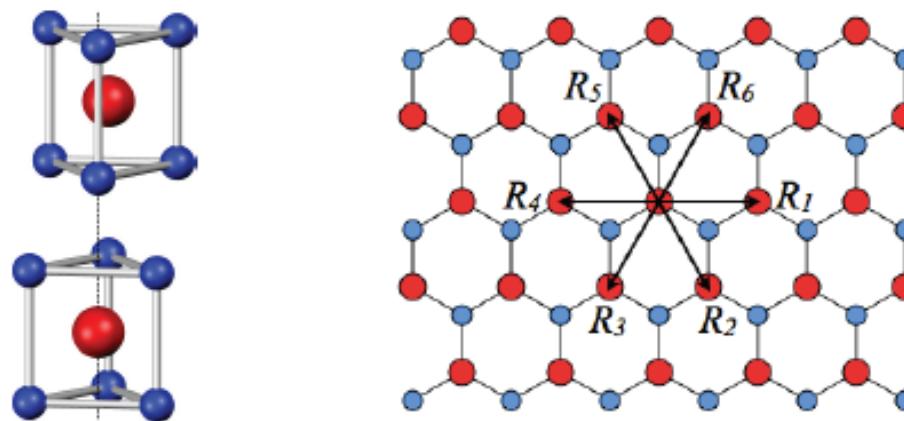
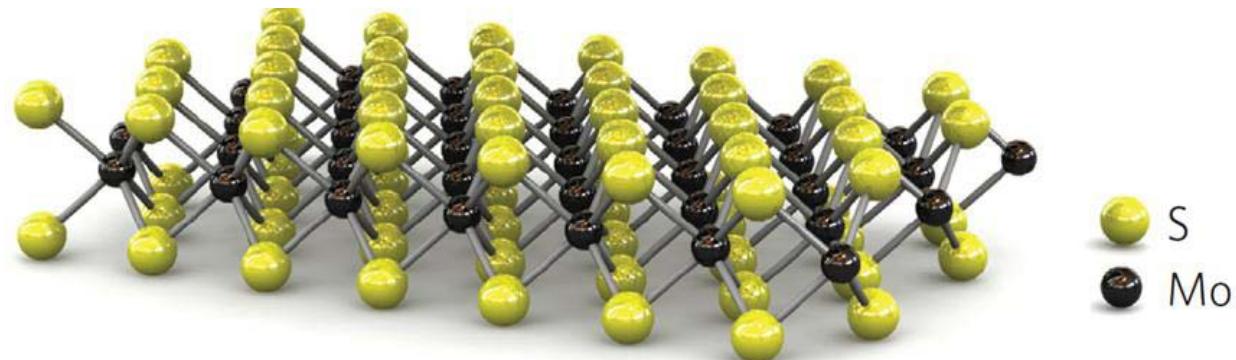


$\Delta_N < \mu/(1 + eV/\Delta_N)$ AR without $\langle \sigma_z \rangle$ inversion
 $\mu/(1 + eV/\Delta_N) < \Delta_N < \mu/(eV/\Delta_N - 1)$ Gap
 $\Delta_N \geq \mu/(eV/\Delta_N - 1)$ AR with $\langle \sigma_z \rangle$ inversion

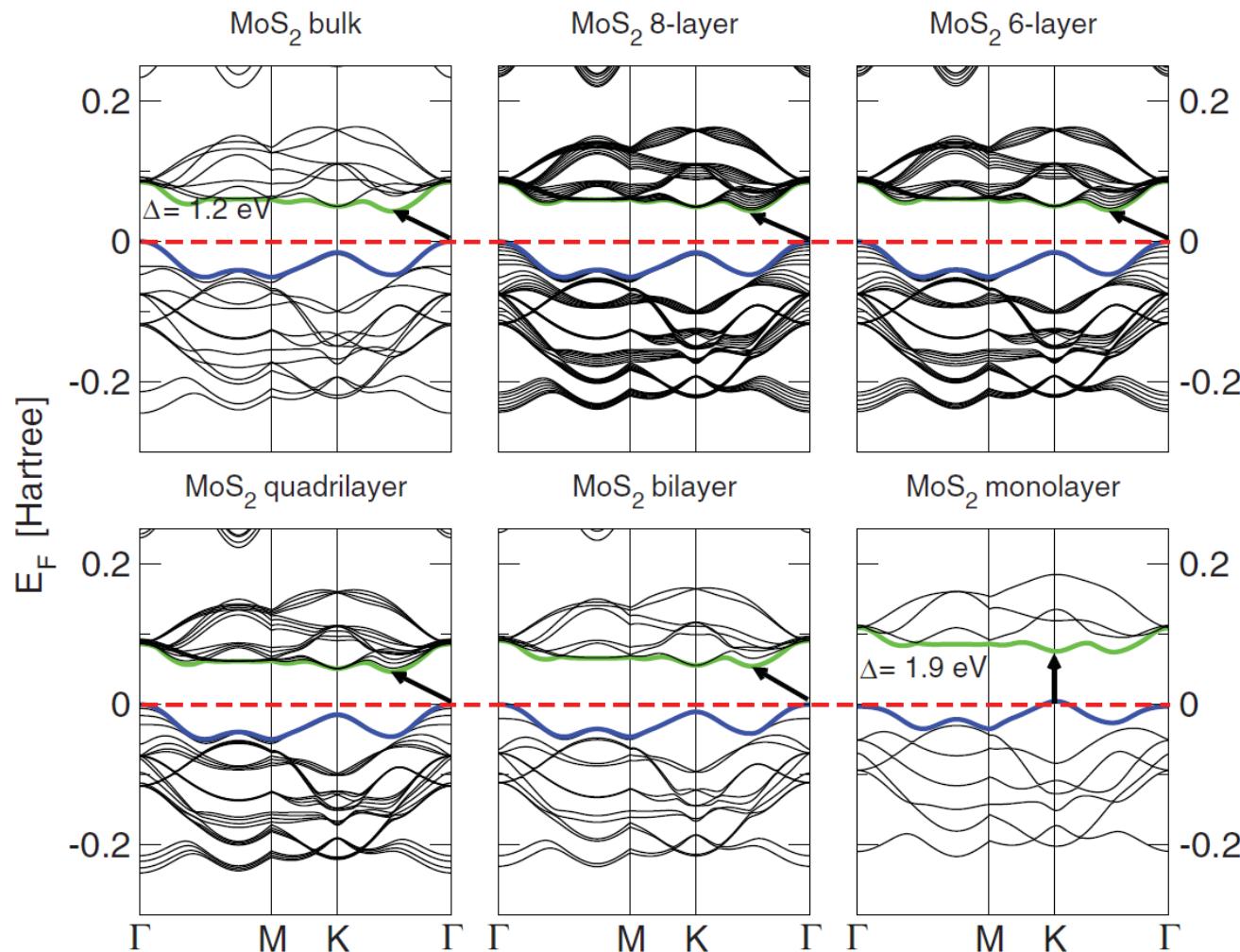
FIG. 2. (Color online) (a) Dependence of Andreev conductance of graphene S/PF contact on the gap Δ_N/μ (in units of the chemical potential) at three bias voltages $eV/\Delta_S = 0, 0.5, 1$. (b) The behavior of the z component of the pseudospin vector $\langle \sigma_z \rangle$ vs Δ_N/μ for the incident electron from the conduction band and the reflected hole from the conduction (before the gap) or valence band (after the gap) at two values of the bias voltage $eV/\Delta_S = 0.5, 1$, when $\Delta_N/\Delta_S = 0.1$.



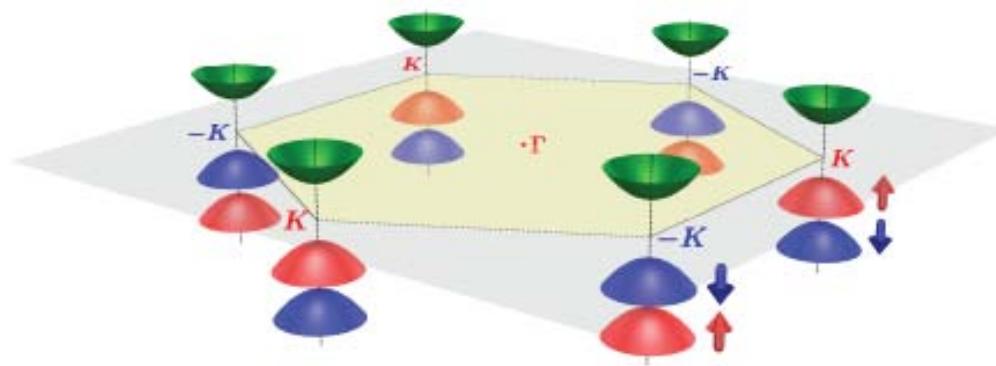
TMDCs: MoS_2 Crystal



Band Structure



Low-energy Hamiltonian



$$\hat{H} = at(\tau k_x \hat{\sigma}_x + k_y \hat{\sigma}_y) + \frac{\Delta}{2} \hat{\sigma}_z - \lambda \tau \frac{\hat{\sigma}_z - 1}{2} \hat{s}_z$$

$$\begin{aligned}\Delta &= 1.9\text{eV} \\ \lambda &= 80\text{meV}\end{aligned}$$

Di Xiao, et al , Phys. Rev. Lett. **108**, 196802 (2012)

The difference between electron and hole masses
recently reported by using ab initio calculations
Phys. Rev. B 86, 241401(R) (2012).

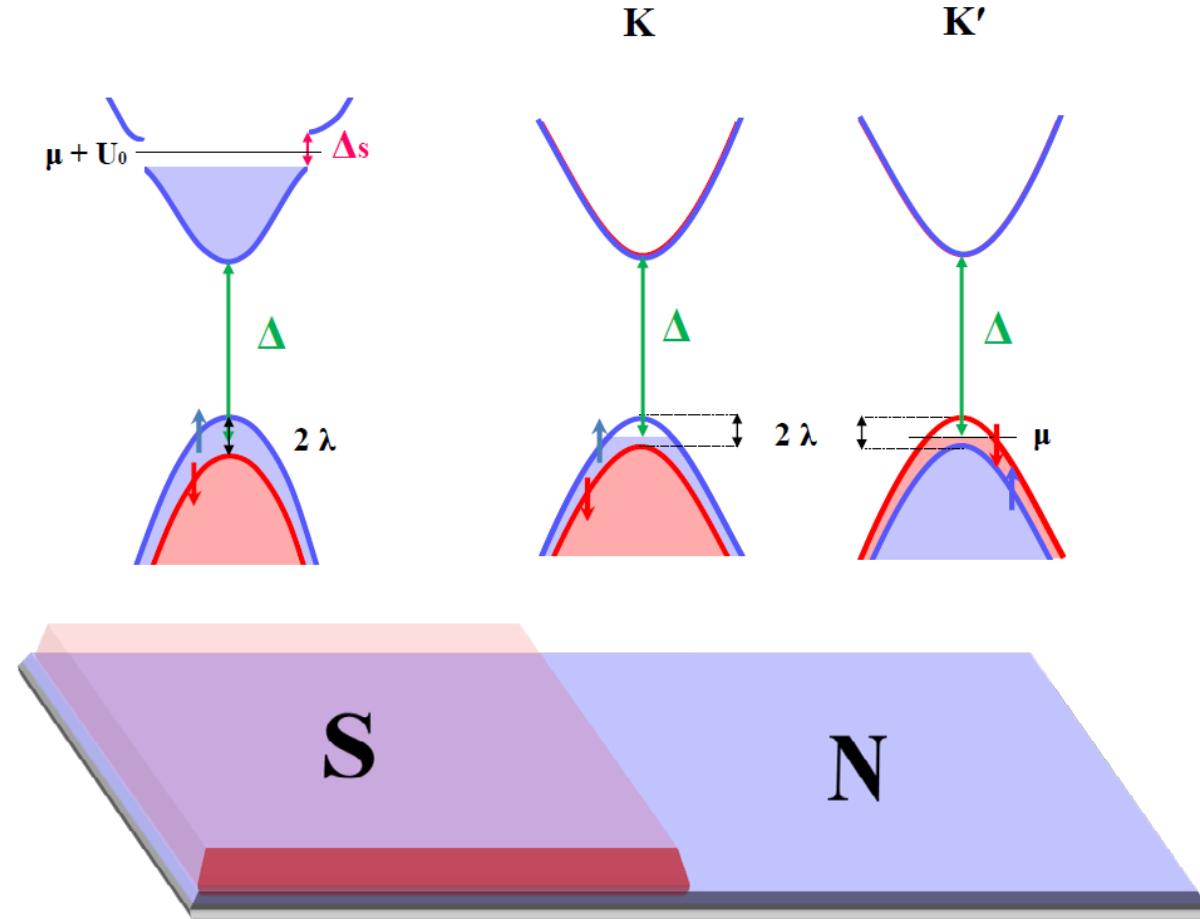
$$H_{\tau s} = \frac{\Delta}{2} \sigma_z + \lambda \tau s \frac{1 - \sigma_z}{2} + t_0 a_0 \mathbf{q} \cdot \boldsymbol{\sigma}_\tau + \frac{\hbar^2 |\mathbf{q}|^2}{4m_0} (\alpha + \beta \sigma_z) + t_1 a_0^2 \mathbf{q} \cdot \boldsymbol{\sigma}_\tau^* \sigma_x \mathbf{q} \cdot \boldsymbol{\sigma}_\tau^*$$

$$\alpha = 0.43$$

$$\beta = 2.21$$

H. Rostami, A. G. Moghaddam, R. Asgari, Phys. Rev. B 88, 085440 (2013)

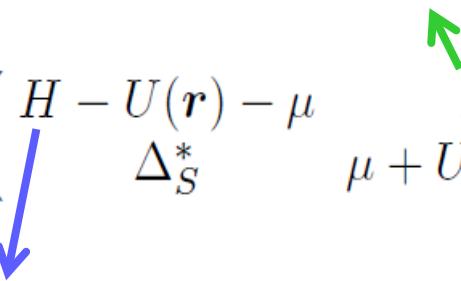
MoS₂-based S/N junction



Dirac-Bogoliubov-de Gennes (DBdG) equation

DBdG equation describes the superconducting correlations between relativistic electrons and holes with **opposite spins** and **different valley indices**

**Electron-hole conversion:
superconducting correlations**

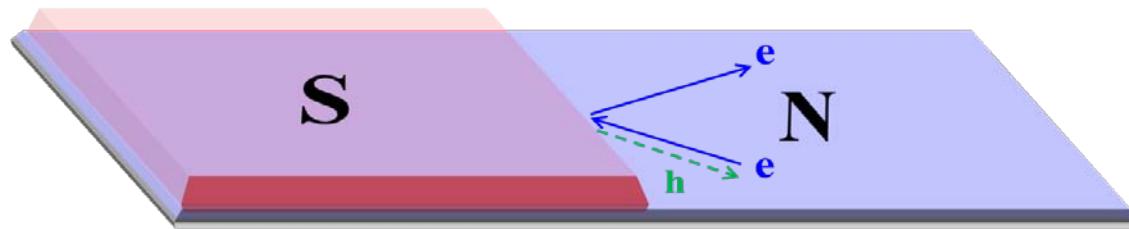
$$\begin{pmatrix} H - U(\mathbf{r}) - \mu & \Delta_S \\ \Delta_S^* & \mu + U(\mathbf{r}) - H \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = \varepsilon \begin{pmatrix} u \\ v \end{pmatrix}$$


2D Hamiltonian of a monolayer MoS_2

$$H = v_F(\boldsymbol{\sigma} \cdot \mathbf{p}) + \frac{\Delta}{2}\sigma_z + \lambda s\tau \left(\frac{1 - \sigma_z}{2} \right) - \frac{\hbar^2 \nabla^2}{4m_0} (\alpha + \beta \sigma_z)$$

$$\begin{cases} \Delta = 1.9\text{eV} & s = \pm 1 \\ \lambda = 80\text{meV} & \tau = \pm 1 \\ \alpha = 0.43 \text{ and } \beta = 2.21 \end{cases}$$

Andreeev Reflection in MoS₂



$$-\Delta/2 - \lambda - \varepsilon < \mu_N < -\Delta/2 + \lambda - \varepsilon$$

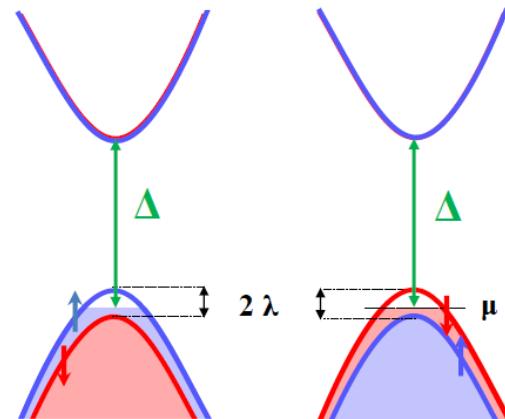
$$\mu_N \leq -\Delta/2 - \lambda - \varepsilon$$

Spin-valley polarized AR process

$$s\tau = 1$$

K

K'



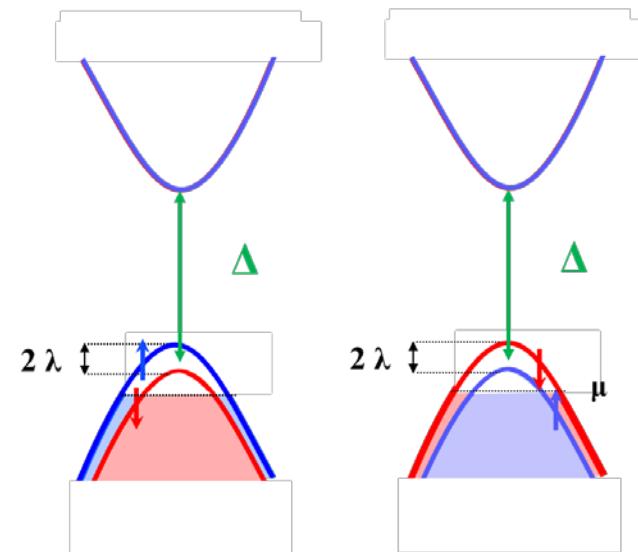
$$s = \tau = 1$$

$$s = \tau = -1$$

$$s\tau = \pm 1$$

K

K'



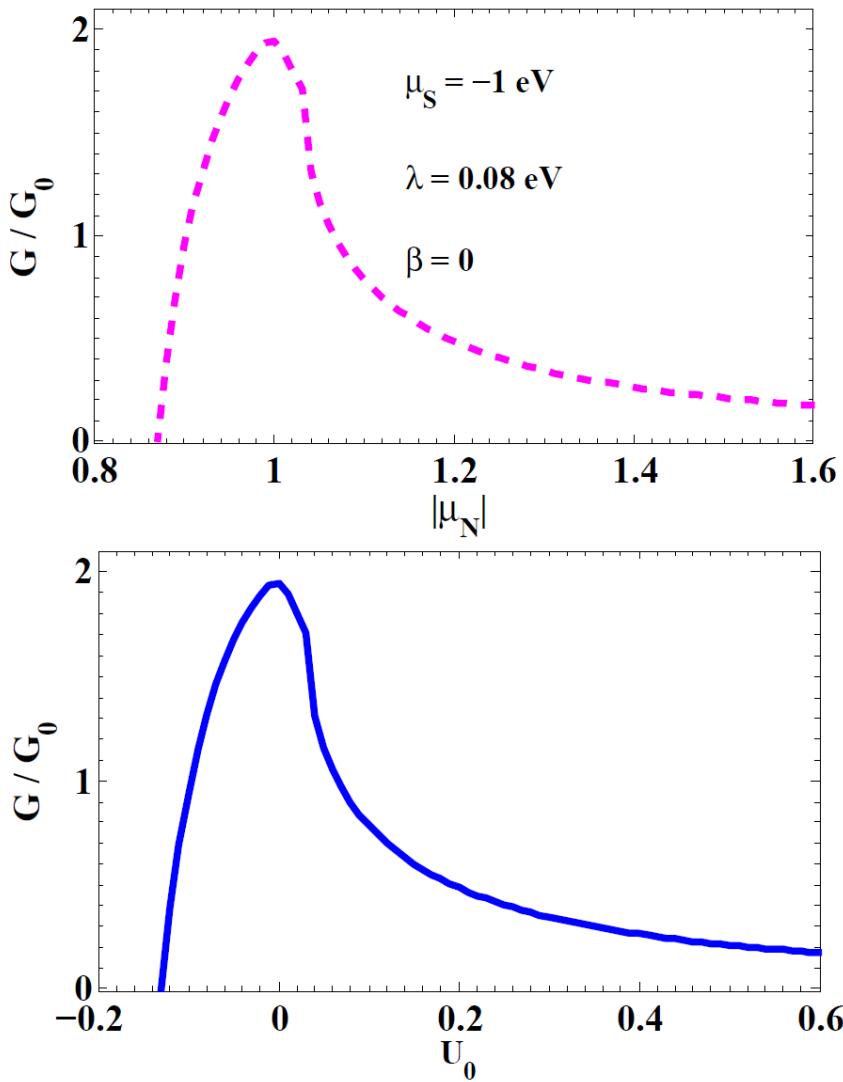
normal reflection	Andreev reflection
$\left\{ \begin{array}{l} \psi_N = \psi^{e-} + r^{s,\tau} \overset{\nearrow}{\psi}{}^{e+} + r_A^{s,\tau} \overset{\nearrow}{\psi}{}^{h+}, \\ \psi_S = t \psi^{S+} + t' \psi^{S-}. \end{array} \right.$	

Andreev conductance (BTK formula)

$$G = \sum_{\tau,s=\pm 1} G_0^{\tau,s} \int_0^{\theta_c} (1 - |r^{\tau,s}|^2 + |r_A^{\tau,s}|^2) \cos \theta_e \, d\theta_e$$

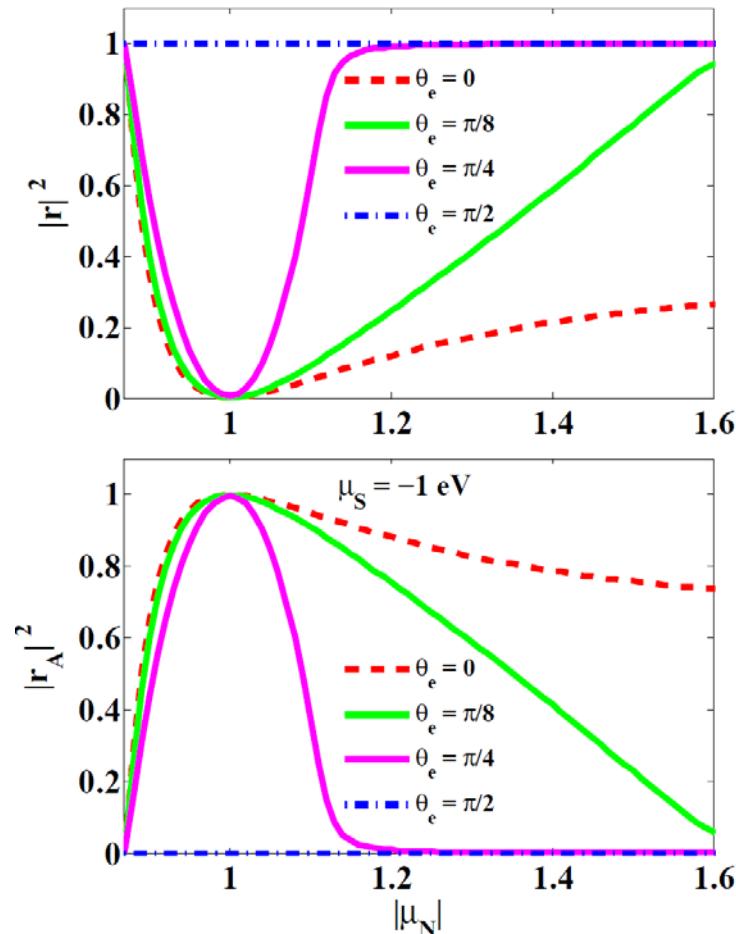
$$G_0^{\tau,s} = \frac{e^2}{h} N_{\tau,s}(eV)$$

Andreev Reflection in S(p)/N(p) Junction

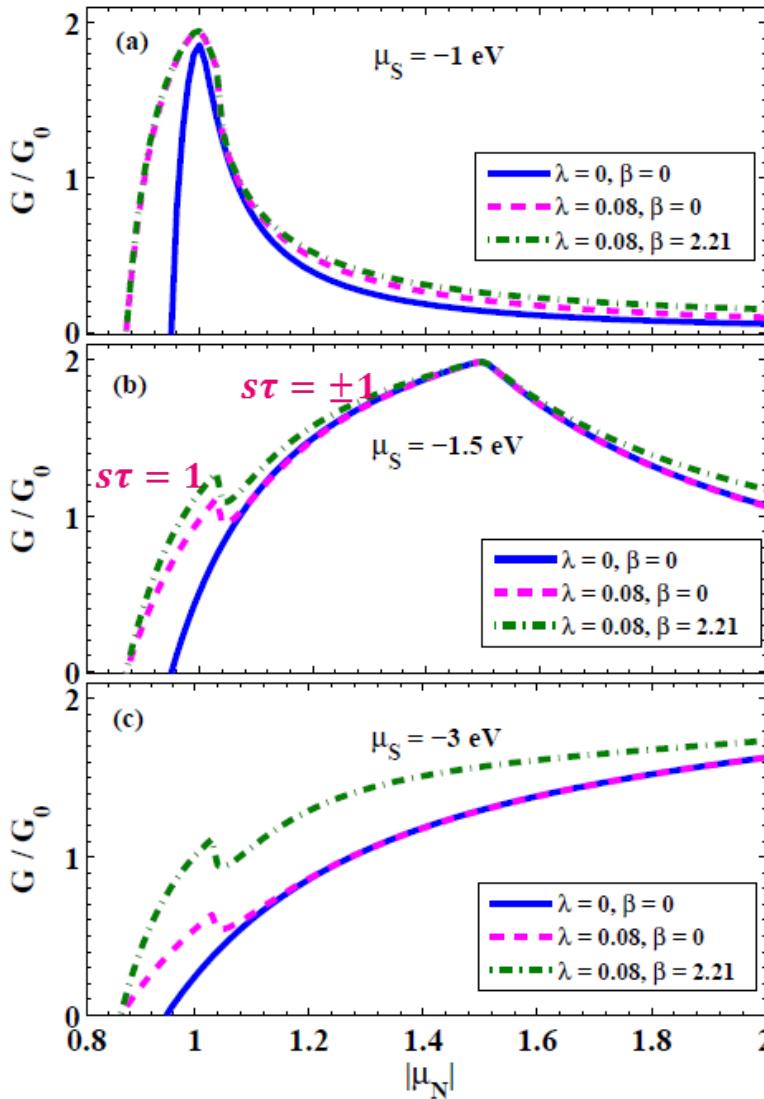


$$\mu_S = \mu + U_0$$

$$\mu_N = \mu \quad (\mu_N < 0)$$



Andreev Reflection in S(p)/N(p) Junction



❖ Gapped Graphene ($\lambda=0, \beta=0$)

❖ MoS₂ ($\lambda=0.08$ eV) with $\beta=0$

❖ MoS₂ ($\lambda=0.08$ eV) with $\beta=2.21$

$$\alpha = 0$$

Amplification of the proximity inducing superconductivity in N region

FIG. 2: Dependence of the Andreev conductance of S/N contact with *p*-doped S region on the chemical potential $|\mu_N|$ (in units of eV) at three chemical potentials $\mu_S = -1, -1.5, -3$ eV, for gapped graphene ($\lambda = 0, \beta = 0$), MoS₂ ($\lambda = 0.08$ eV) with $\beta = 0$ and $\beta = 2.21$, when $\alpha = 0$, $\Delta_S = 0.01$ eV and $eV/\Delta_S = 0$.

Andreev Reflection in S(p)/N(p) Junction

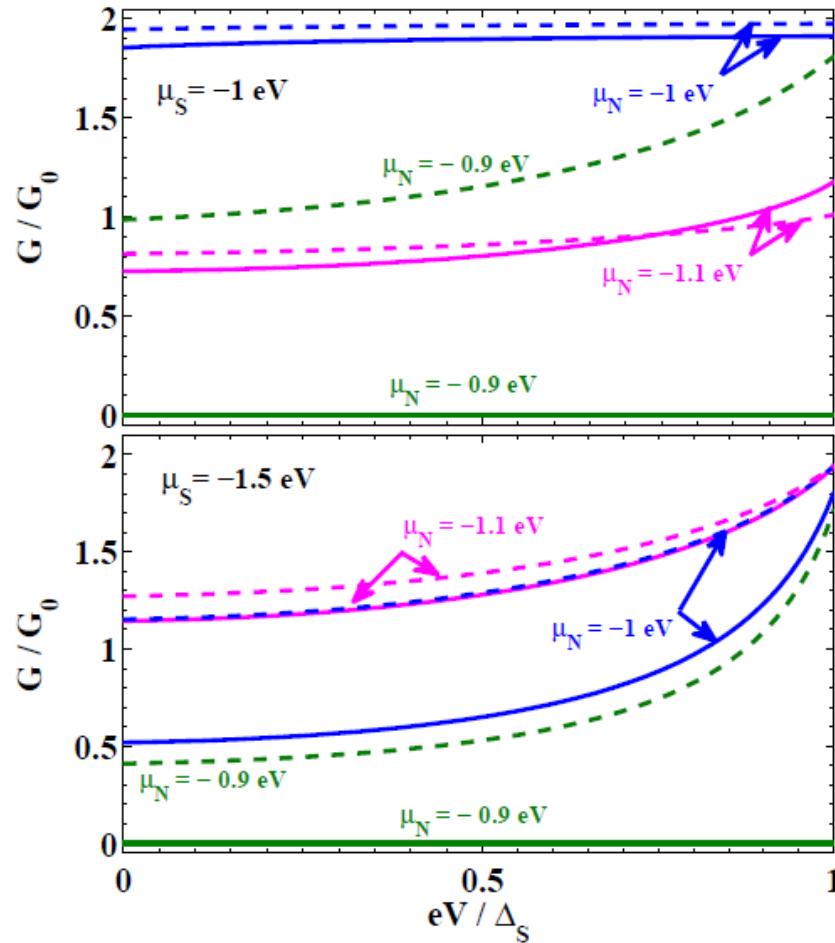
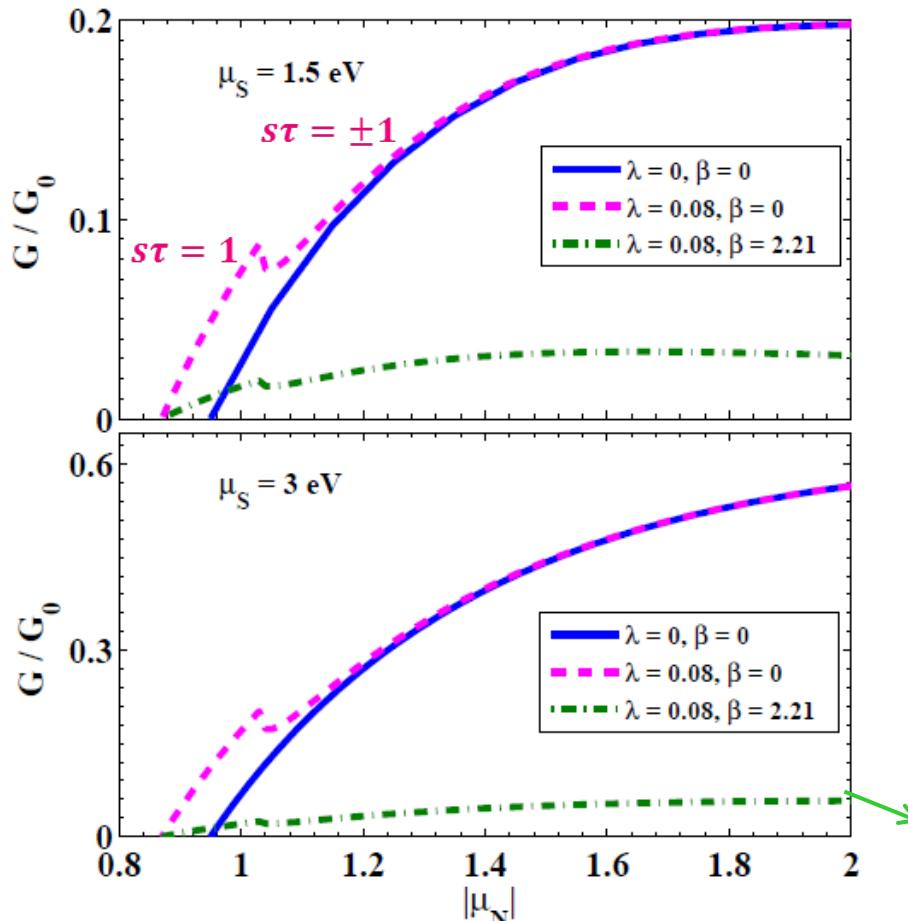


FIG. 3: Plot of the Andreev conductance of S/N contact (with *p*-doped S region) versus the bias voltage eV/Δ_S (in units of the superconducting gap Δ_S) at two values of $\mu_S = -1, -1.5 \text{ eV}$, for gapped graphene with $\lambda = 0, \beta = 0$ (solid lines) and MoS_2 with $\lambda = 0.08 \text{ eV}, \beta = 2.21$ (dashed lines), when $\mu_N = -0.9, -1, -1.1 \text{ eV}$ and $\alpha = 0$.

Andreev Reflection in S(n)/N(p) Junction



$$\mu_S = \mu + U_0$$

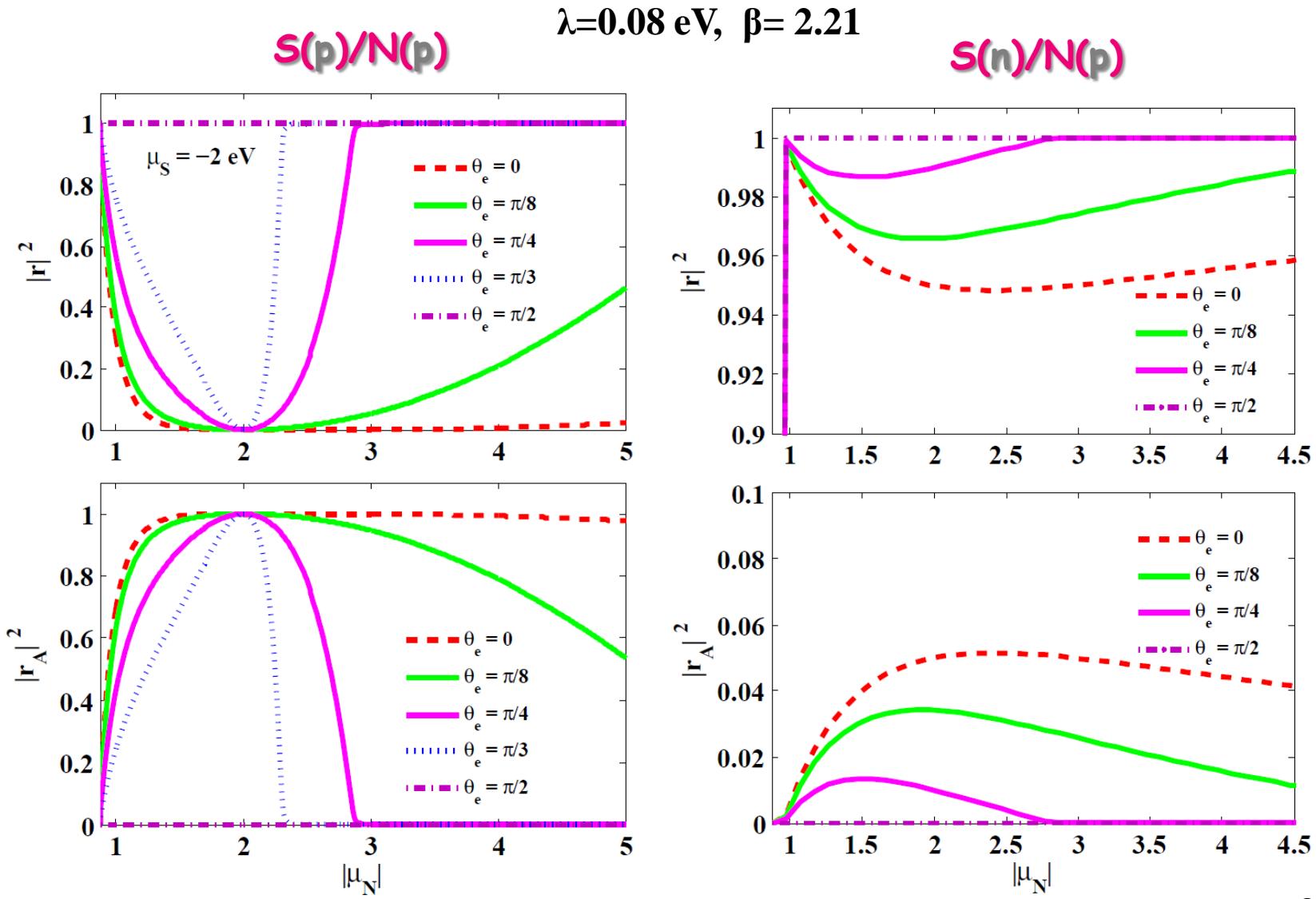
$$\mu_N = \mu \quad (\mu_N < 0)$$

- ❖ Gapped Graphene ($\lambda=0, \beta= 0$)
- ❖ MoS₂ ($\lambda=0.08$ eV) with $\beta= 0$
- ❖ MoS₂ ($\lambda=0.08$ eV) with $\beta= 2.21$

Attenuation of the proximity
inducing superconductivity
in N region

FIG. 4: Andreev conductance of S/N contact with n -doped S region versus $|\mu_N|$ at two chemical potentials $\mu_S = 1.5, 3$ eV, for gapped graphene ($\lambda = 0, \beta = 0$), MoS₂ ($\lambda = 0.08$ eV) with $\beta = 0$ and $\beta = 2.21$, when $\alpha = 0$, $\Delta_S = 0.01$ eV and $eV/\Delta_S = 0$.

Andreev Reflection in S/N Junction



Conclusion

- ❖ Presence of strong spin-orbit interaction in MoS₂ enhances the amplitude of Andreev reflection and the resulting Andreev conductance in *MoS₂*-based S/N structures with n-(p-) doped S and p-doped N regions.
- ❖ Presence of β term enhances the Andreev conductance in *MoS₂*-based S/N structures with p-doped S, while it decreases the Andreev conductance in S/N structure with n-doped S region.

$$H = v_F(\boldsymbol{\sigma} \cdot \mathbf{p}) + \frac{\Delta}{2}\sigma_z + \lambda s\tau \left(\frac{1 - \sigma_z}{2} \right) - \frac{\hbar^2 \nabla^2}{4m_0} (\alpha + \beta \sigma_z)$$

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