

# Correlation effects in Fermion Systems

## Lecture II

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# Outlines

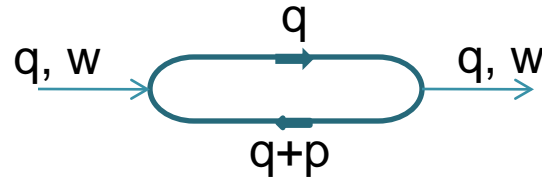
## 1) Correlations in Electron Liquid

- Introduction
- Dyson equation and self-energy
- Effective mass
- Fermi Liquid & non-Fermi Liquid
- Luttinger Liquid, Bosonization

## 2) Hubbard Model

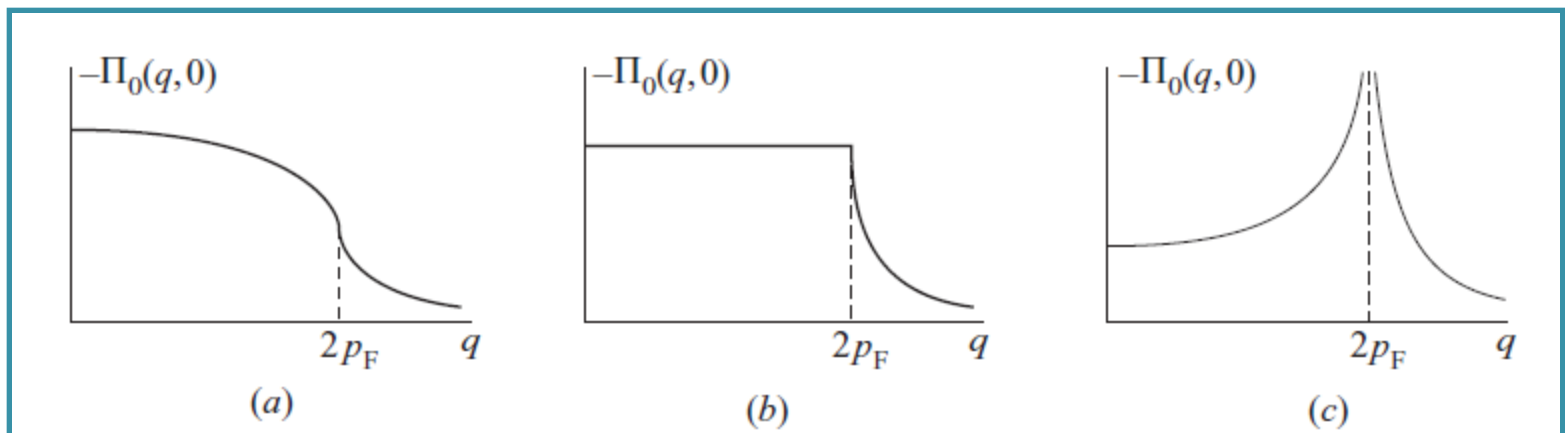
- Introduction
- Spin-charge separation
- Mott insulators
- Phase diagram

# Non-interacting response functions



Bubble

$$\begin{aligned} \Pi_0(\mathbf{q}, \omega) &= \sum_{\alpha\beta} \frac{n_\alpha - n_\beta}{\hbar\omega + \varepsilon_\alpha - \varepsilon_\beta + i\hbar\eta} A_{\alpha\beta} B_{\beta\alpha} \\ &= -\frac{2i}{(2\pi)^4} \int \frac{d^3 p d\omega'}{\left[ \omega + \omega' - \varepsilon(\mathbf{p} + \mathbf{q}) + i\delta \operatorname{sign}(\varepsilon(\mathbf{p} + \mathbf{q}) - \mu) \right] \left[ \omega' - \varepsilon(\mathbf{p}) + i\delta \operatorname{sign}(\varepsilon(\mathbf{p}) - \mu) \right]} \end{aligned}$$



# Minimum theory: RPA

$$\begin{aligned}
 \tilde{\Pi} &= \text{diagram with cross-hatch} = \overset{\Pi_0}{\text{diagram with empty oval}} + \overset{V(q)}{\text{diagram with empty oval connected to empty oval}} + \dots \\
 &= \text{diagram with empty oval} + \overset{V(q)}{\text{diagram with empty oval connected to cross-hatch}}
 \end{aligned}$$

$$\tilde{\Pi} = \Pi_0 + \Pi_0 V \tilde{\Pi} .$$

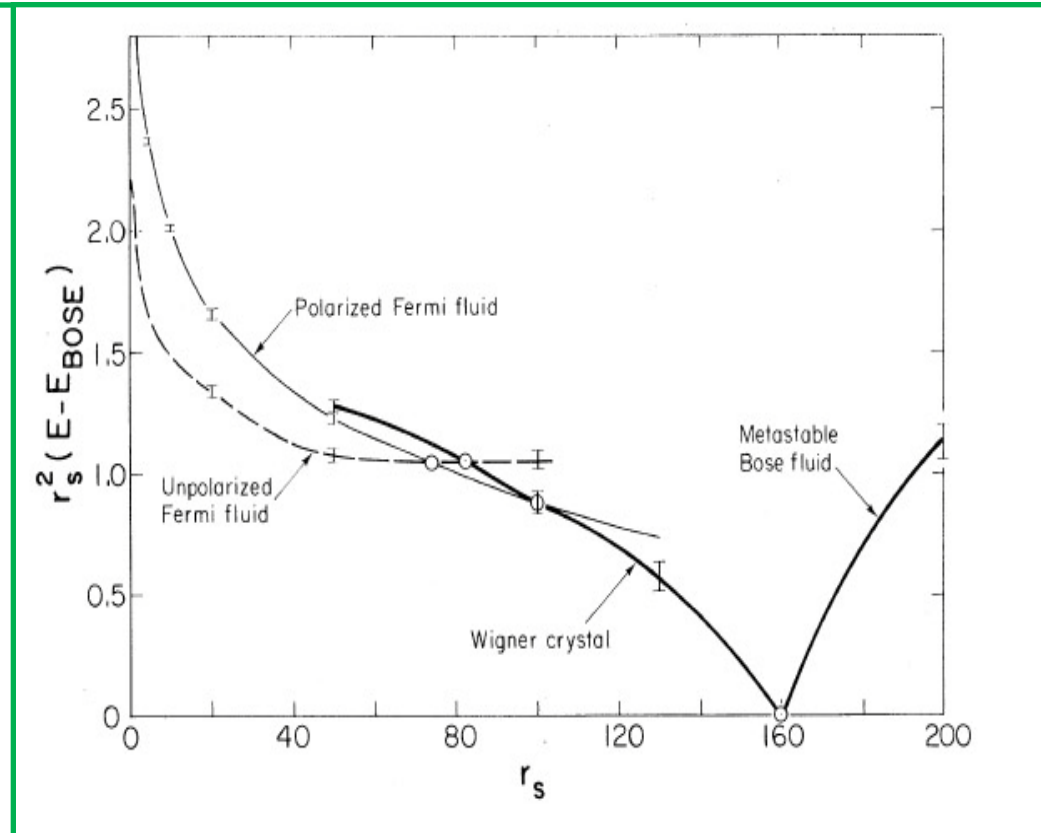
$$\tilde{\Pi} = \frac{\Pi_0}{1 - \Pi_0 V}$$

Effective potential

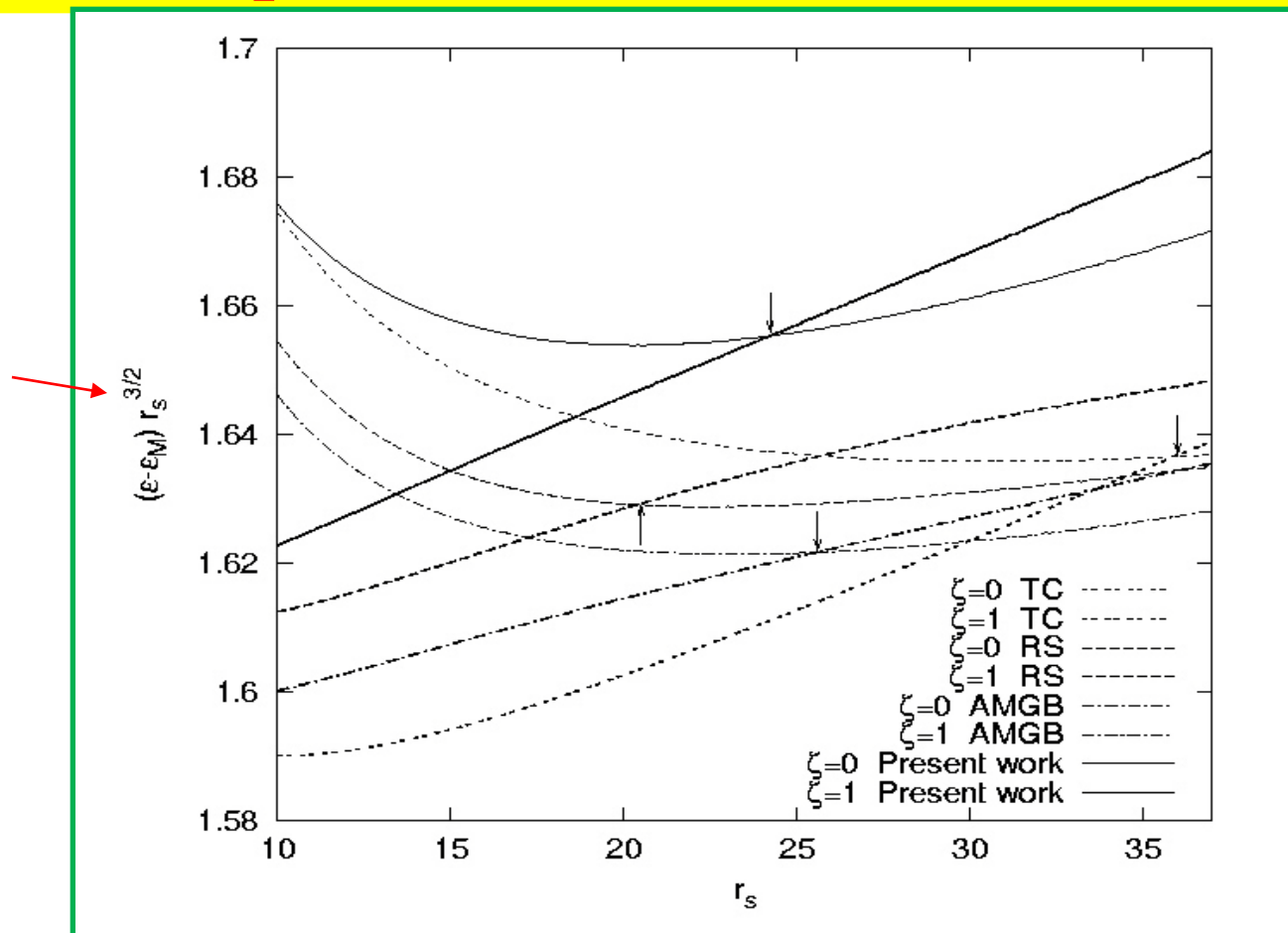
$$v = \frac{V}{1 - \Pi_0 V} = \frac{V}{\epsilon}$$

# Jellium Model: High density region, Correlation

$$\epsilon(r_s) \simeq \begin{cases} \left( \frac{2.210}{r_s^2} - \frac{0.916}{r_s} + 0.062 \ln r_s - 0.093 + \mathcal{O}(r_s \ln r_s) \right) Ry, & (3D) \\ \left( \frac{1}{r_s^2} - \frac{1.20}{r_s} - 0.17 r_s \ln r_s - (0.38 \pm 0.04) + \mathcal{O}(r_s) \right) Ry, & (2D) \end{cases}$$



# paramagnetic to fully spin-polarized quantum phase transition of a 2D EL



**TC:** B. Tanatar and D. M. Ceperley, Phys. Rev. B **39**, 5005 (1989)

**RS:** F. Rapisarda and G. Senatore, Aust. J. Phys. **49**, 161 (1996)

**AMGB :** C. Attaccalite, *et al.*, Phys. Rev. Lett **88**, 256601 (2002)

R. Asgari, B. Davoudi and M. P. Tosi, Solid State Communication 131,1(2004)

# Spectral function

$$G(\mathbf{p}, \omega) = \int_{-\infty}^{\infty} \frac{A(\mathbf{p}, \omega') d\omega'}{\omega - \omega' + i\omega'\delta}$$

$$\text{Im } G(\mathbf{p}, \omega + \mu) = \begin{cases} -\pi A(\mathbf{p}, \omega) & \text{for } \omega > 0 \\ \pi A(\mathbf{p}, \omega) & \text{for } \omega < 0 \end{cases}$$

$$A(\mathbf{p}, \omega) = A^*(\mathbf{p}, \omega) > 0$$

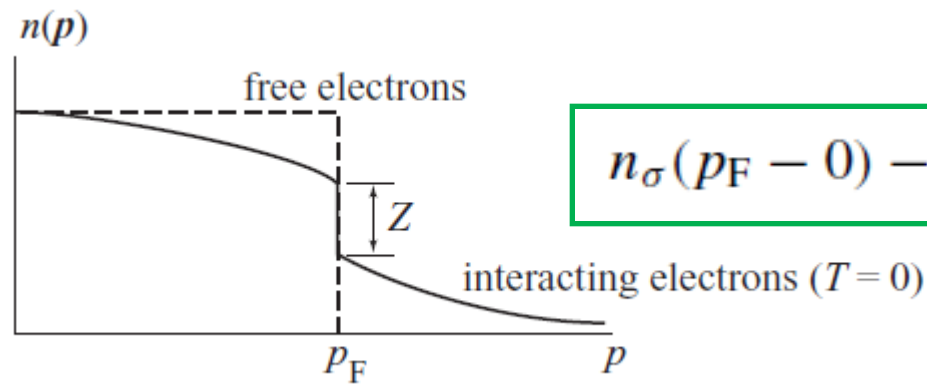
Sum rule  $\int_{-\infty}^{\infty} A(\mathbf{p}, \omega) d\omega = 1$

Non-interacting  $A^0(\mathbf{p}, \omega) = \delta(\omega - [\varepsilon_{\mathbf{p}} - \mu]) = \delta(\omega - \xi_{\mathbf{p}})$

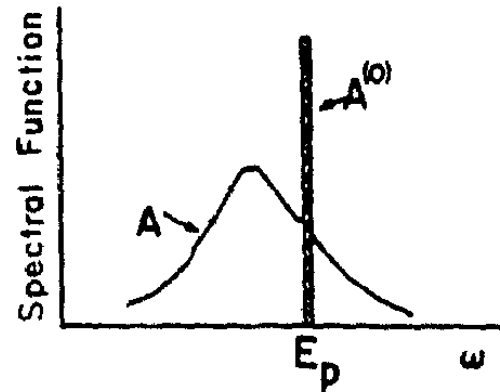
$$A(\mathbf{p}, \omega) = \sum_m |(c_{\mathbf{p}}^\dagger)_{m,0}|^2 \delta(\omega - \omega_m^{N+1}) + \sum_m |(c_{\mathbf{p}})_{m,0}|^2 \delta(\omega - \omega_m^{N-1})$$

# Renormalization constant

$$G(p, \omega) \simeq \frac{Z_p}{\omega - \tilde{\varepsilon}_p + i\gamma}$$



$$n_\sigma(p_F - 0) - n_\sigma(p_F + 0) = Z$$



$$A(p, \omega) \simeq Z_p \delta(\omega - \varepsilon_p)$$



# Dyson equation

$$G = G_0 + G_0 \Sigma_0 G_0 + G_0 \Sigma_0 G_0 \Sigma_0 G_0 + \dots$$

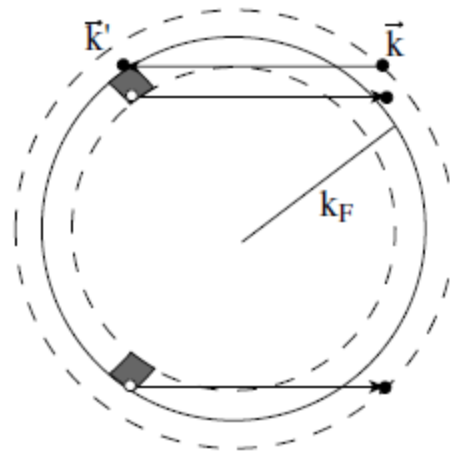
Hartree term  $\Sigma_0(p, \omega) = \int V(q, \omega') G_0(p - q, \omega - \omega') \frac{d^3 q d\omega'}{(2\pi)^4}$



$$G(p, \omega) = \frac{G_0}{1 - \Sigma_0 G_0} = \frac{1}{G_0^{-1} - \Sigma_0} = \frac{1}{\omega - \varepsilon_p - \Sigma_0(p, \omega) + i\delta_p}$$

$$G = G_0 + G_0 \Sigma_0 G ;$$

# Fermi Liquid Theory



$$E[\{N_{\vec{k}\sigma}\}] = E_0 + \sum_{\vec{k}\sigma} \mathcal{E}_{\vec{k}\sigma} \delta N_{\vec{k}\sigma} + \frac{1}{2} \sum_{\vec{k}\sigma, \vec{k}'\sigma'} f_{\vec{k}\sigma, \vec{k}'\sigma'} \delta N_{\vec{k}\sigma} \delta N_{\vec{k}'\sigma'}$$

$$\mathcal{E}_{\vec{k}\sigma} = \left( \frac{\delta E}{\delta N_{\vec{k}\sigma}} \right)_{N_{\vec{k}\sigma} = N_{\vec{k}\sigma}^{(0)}} \quad N_{\vec{k}\sigma}^{(0)} = \Theta(k_F - k)$$

$$\mathcal{E}_{\vec{k}\sigma} \simeq \mu + \hbar v_F^* (k - k_F) \quad v_F^* = \frac{1}{\hbar} \left| \frac{\partial \mathcal{E}_{\vec{k}\sigma}}{\partial \vec{k}} \right|_{k=k_F}$$

$$\tilde{\mathcal{E}}_{\vec{k}\sigma} = \frac{\delta E}{\delta N_{\vec{k}\sigma}} = \mathcal{E}_{\vec{k}\sigma} + \sum_{\vec{k}'\sigma'} f_{\vec{k}\sigma, \vec{k}'\sigma'} \delta N_{\vec{k}'\sigma'}$$

# Fermi Liquid Theory

$$F_\ell^{s,a} = \frac{L^d N^*(0)}{2} \int \frac{d\Omega_d}{\Omega_d} [f_{\uparrow\uparrow}(\cos\theta) \pm f_{\uparrow\downarrow}(\cos\theta)] \begin{cases} P_\ell(\cos\theta), & 3D, \\ \cos\ell\theta, & 2D, \end{cases}$$

$$f_{\uparrow\uparrow}(\cos\theta) \pm f_{\uparrow\downarrow}(\cos\theta) = \frac{2}{L^d N^*(0)} \sum_{\ell=0}^{\infty} F_\ell^{s,a} \begin{cases} (2\ell+1)P_\ell(\cos\theta), & 3D, \\ (2-\delta_{\ell 0})\cos\ell\theta, & 2D, \end{cases}$$

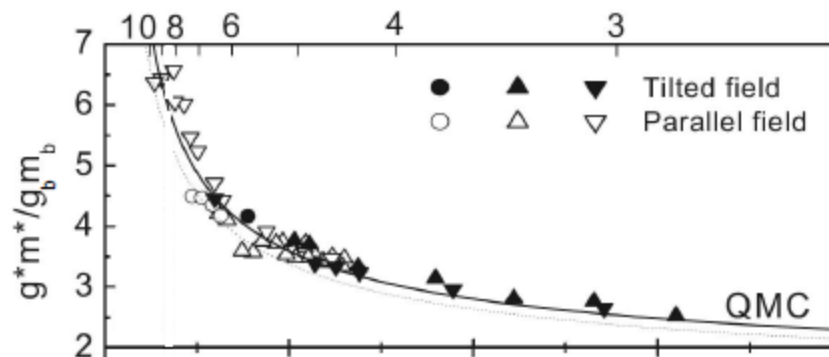
$$\frac{K}{K_0} = \frac{\frac{m^*}{m}}{1 + F_0^s}$$

$$\frac{\chi_S}{\chi_P} = \frac{\frac{m^*}{m}}{1 + F_0^a}$$

$$\frac{\chi_S}{\chi_P} = \frac{m^*}{m} \frac{g^*}{g}$$

$$\frac{m^*}{m} = 1 + F_1^s,$$

$$\vec{v} \cdot \vec{k} = \frac{m}{m^*} \vec{v} \cdot \vec{k} + \frac{m}{m^*} \vec{v} \cdot \vec{k} \frac{L^d N^*(0)}{2} \int \frac{d\Omega'_d}{\Omega_d} \cos\theta' \sum_{\sigma'} f_{\sigma\sigma'}(\cos\theta'),$$



# Effective mass

$$A_{\sigma}(\vec{k}, \omega) = -\frac{\hbar}{\pi} \frac{\Im m \Sigma_{\sigma}^{ret}(\vec{k}, \omega)}{[\hbar\omega - \varepsilon_{\vec{k}\sigma} - \Re e \Sigma_{\sigma}^{ret}(\vec{k}, \omega)]^2 + [\Im m \Sigma_{\sigma}^{ret}(\vec{k}, \omega)]^2}$$

$$\hbar\omega - \varepsilon_{\vec{k}\sigma} - \Re e \Sigma_{\sigma}^{ret}(\vec{k}, \omega) \approx \frac{(\hbar\omega - \varepsilon_{\vec{k}\sigma})}{Z_{\vec{k}\sigma}}, \quad Z_{\vec{k}\sigma} \equiv \left( 1 - \frac{1}{\hbar} \frac{\partial}{\partial \omega} \Re e \Sigma_{\sigma}^{ret}(k, \omega) \Big|_{\hbar\omega = \varepsilon_{\vec{k}\sigma}} \right)^{-1}$$

$$A_{\sigma}(k, \omega) \simeq \frac{Z_{\vec{k}\sigma}}{\pi} \frac{\frac{1}{2\tau_{\vec{k}\sigma}}}{\left( \omega - \frac{\varepsilon_{\vec{k}\sigma}}{\hbar} \right)^2 + \left( \frac{1}{2\tau_{\vec{k}\sigma}} \right)^2}$$

$$\varepsilon_{k\sigma} \simeq \mu + \frac{\hbar^2 k_F (k - k_F)}{m^*} \quad \frac{\hbar^2 k_F}{m^*} = \frac{d\varepsilon_{k\sigma}}{dk} \Big|_{k=k_F}$$

$$\frac{m^*}{m} = \frac{1}{Z_{k_F\sigma} \left[ 1 + \frac{m}{\hbar^2 k_F} \frac{\partial}{\partial k} \Re e \Sigma_{\sigma}^{ret}(k, \mu) \Big|_{k=k_F} \right]}$$

# Effective mass

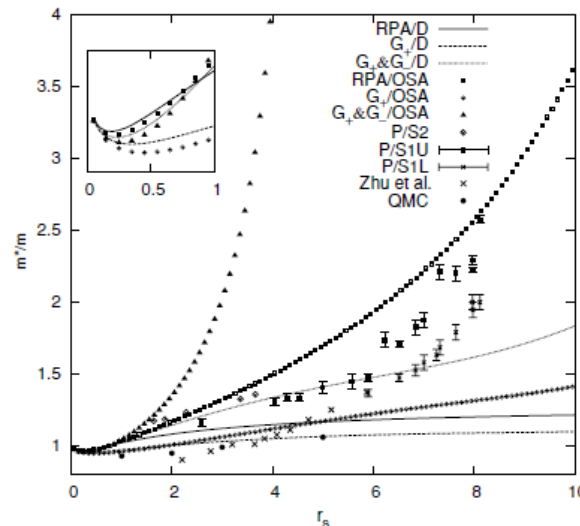
$$\Sigma_{\sigma}(\vec{k}, i\omega_n) = -\frac{1}{\hbar\beta} \int \frac{d\vec{q}}{(2\pi)^d} \sum_{m=-\infty}^{\infty} \mathcal{G}_{\sigma}(\vec{k} - \vec{q}, i\omega_n - i\Omega_m) W(q, i\Omega_m)$$

$$W(q, \omega) = v_q + v_n^2(q) \chi_{nn}(q, \omega) + 3v_s^2(q) \chi_{s_z s_z}(q, \omega)$$

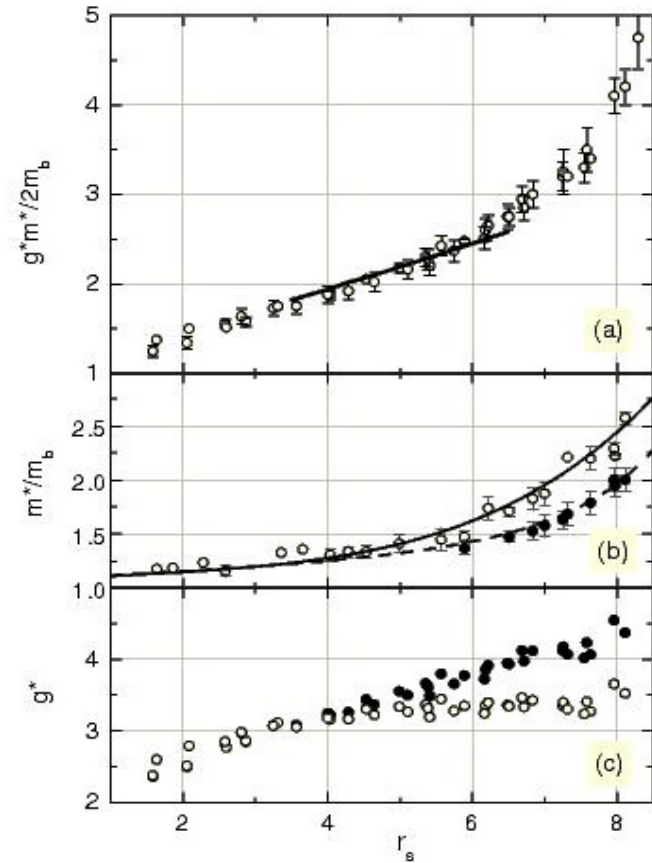
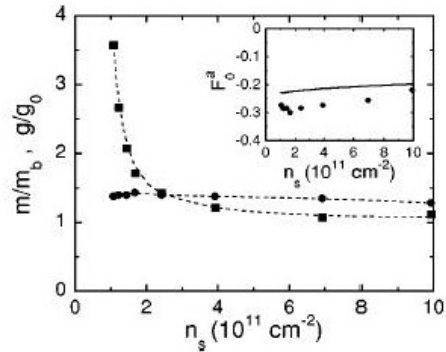
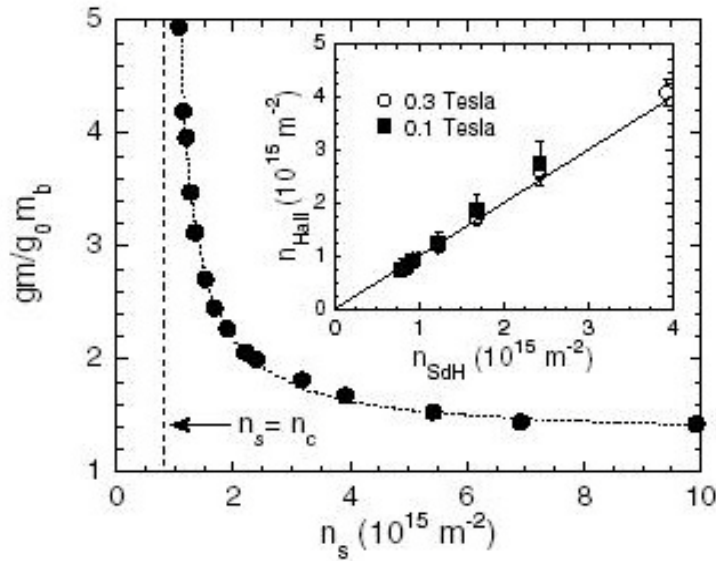
$$V_{\uparrow\uparrow}^{eff}(q) = v_q + \{v_q[1 - G_+(q)]\}^2 \chi_{nn}(q) + \{v_q G_-(q)\}^2 \chi_{s_z, s_z}(q),$$

$$V_{\uparrow\downarrow}^{eff}(q) = \{v_q G_T(q)\}^2 \frac{1}{4} \chi_{s_+ s_-}(q),$$

Local field factors



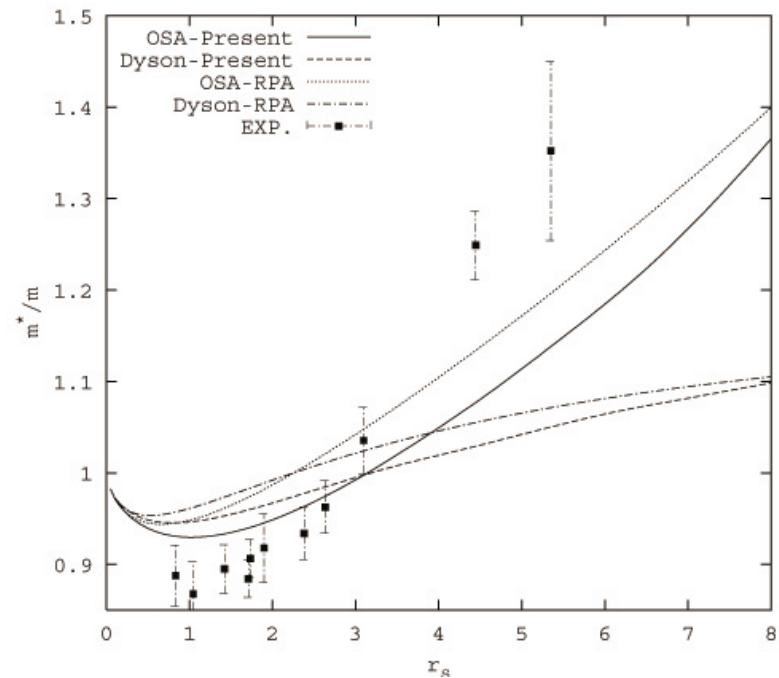
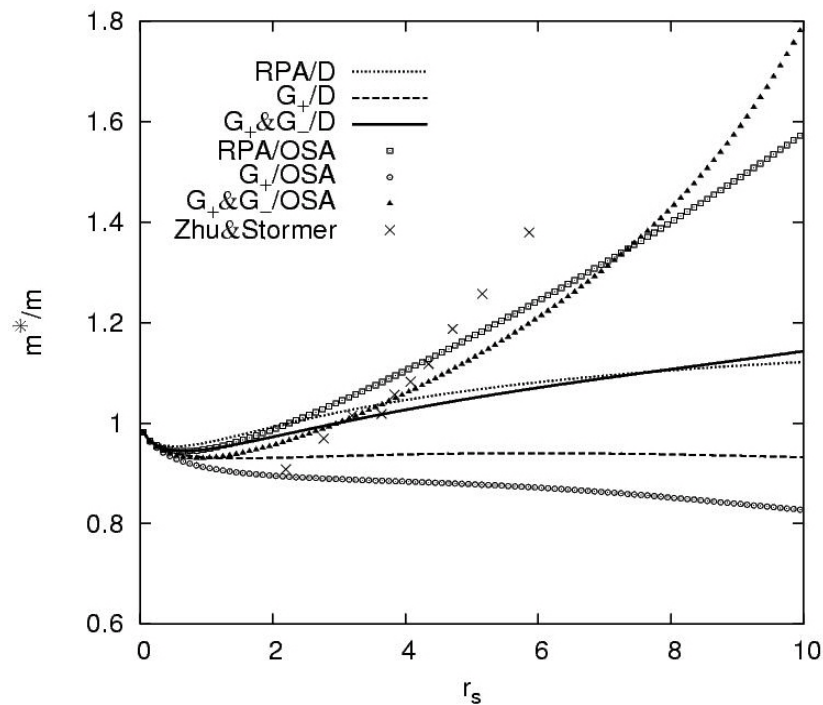
# Some of the recent experimental results in Silicon MOSFETs



V. M. Pudalov et al. Phys. Rev. L **88**, 196404 (2002)

A. Shashkin et al. Phys. Rev. L **87**, 086801 (2001)

# Numerical results: theory versus experiment

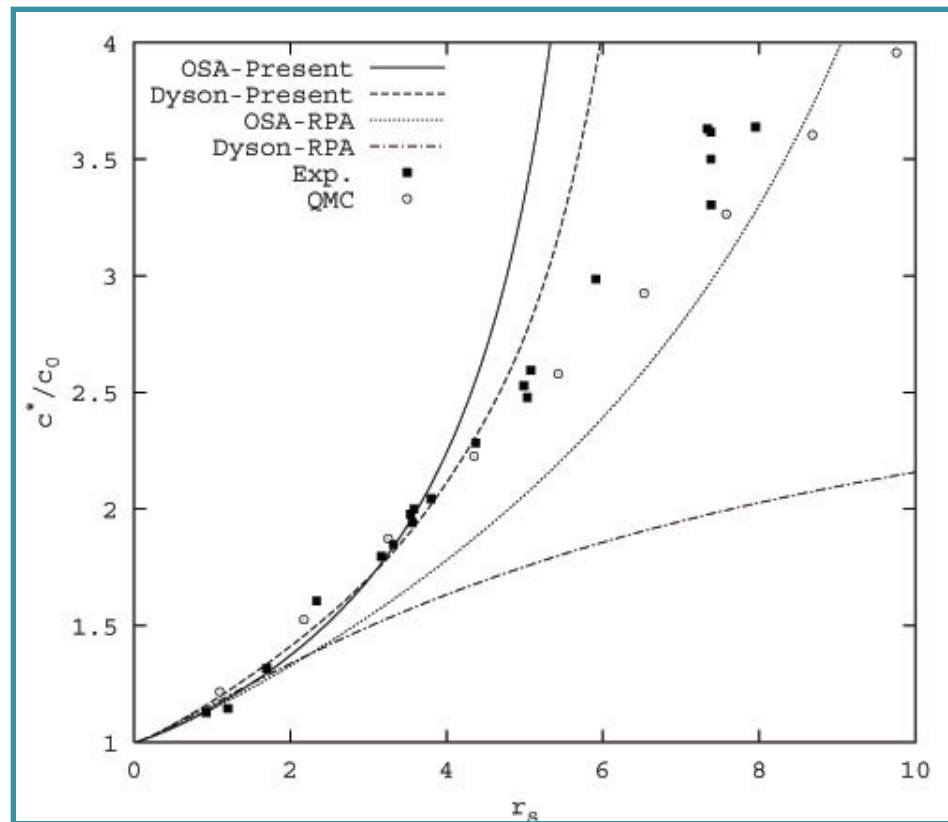


R. Asgari, B. Tanatar Phys. Rev. B 74, 075301 (2006)

R. Asgari, B. Davoudi and B. Tanatar, Solid State Commun. **130**,13(2004)

R. Asgari, B. Davoudi, M. Polini, G. Giuliani, M.P. Tosi and G. Vignale, Phys. Rev. B, **71** 045323 (2005)

# Spin Susceptibility: theory versus experiment



R. Asgari, B. Tanatar Phys. Rev. B 74, 075301 (2006)

R. Asgari, L. Subasi, A. Sabori and B. Tanatar Phys. Rev. B 74, 155319 (2006)



# Marginal Fermi Liquid

$$\text{Im } \Pi(\mathbf{q}, \omega) = \begin{cases} -\rho(\varepsilon_F) \omega / T & \text{for } \omega \ll T \\ -\rho(\varepsilon_F) & \text{for } T \ll \omega \ll \omega_c. \end{cases}$$

$$\text{Re } \Pi(\mathbf{q}, \omega) \sim \rho(\varepsilon_F) \ln \left( \frac{\omega}{T} \right)$$

$$\Sigma = \text{---} \overbrace{\text{---}}^{\text{---}} \text{---} = \text{---} \overbrace{\text{---}}^{\text{---}} \text{---}$$

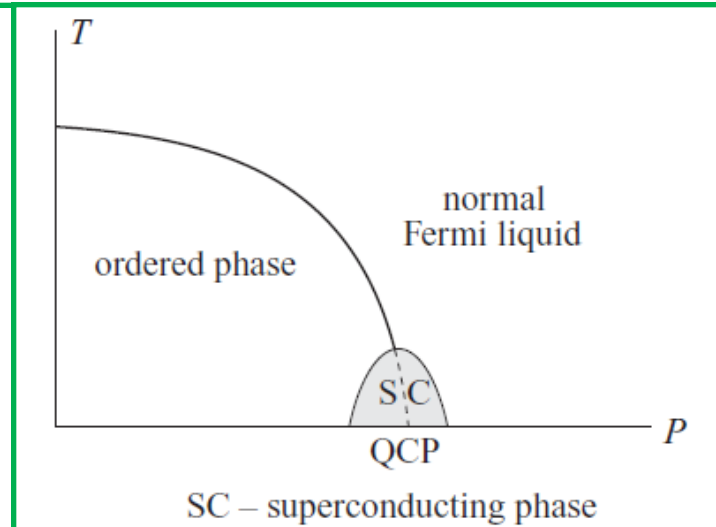
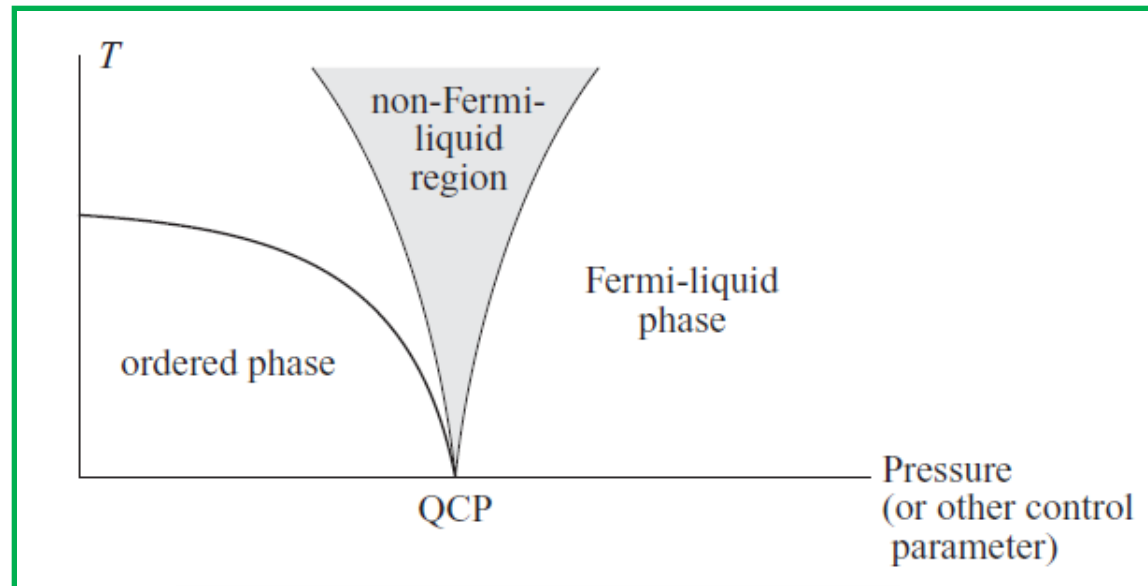
$$\Sigma(\omega) = \lambda \omega \left[ \ln \frac{x}{\omega_c} + i \frac{\pi}{2} x \text{sign}(\omega) \right] \quad x = \max(\omega, T).$$

$$Z = \frac{1}{1 - \partial \text{Re } \Sigma / \partial \omega} \sim \frac{1}{1 - \lambda \ln(y/\omega_c)}$$

$$y = \max[(\varepsilon_p - \mu), T]$$

$$\tau^{-1} \sim (\varepsilon_p - \mu).$$

# Non-Fermi Liquid close to a QCP



# Luttinger Liquid

$$n(p) - n(p_F) \sim |p - p_F|^\delta \text{sign}(p - p_F).$$

