

Reflections on Black Mirrors

A one-to-one correspondence between
moving mirrors and black holes.

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Astana, Kazakhstan

IPM, Tehran, Iran, May 2016

NSF Grants to WFU: PHY-0856050, PHY-1308325, and PHY-1505875,

NSF Grants to UNC: PHY-1506182,

[arXiv:gr-qc/1507.03489](https://arxiv.org/abs/gr-qc/1507.03489).

Late Time Planck Distribution

$$T = \frac{\hbar c^3}{8\pi kGM} \longrightarrow \frac{\kappa}{2\pi}$$

$$T = \frac{\hbar \kappa}{2\pi c k} \longrightarrow \frac{\kappa}{2\pi}$$

Planck Distribution

$$\bar{n} = \frac{1}{e^{\frac{\omega}{T}} - 1}$$

$$T = \frac{\kappa}{2\pi}$$

Correspondences

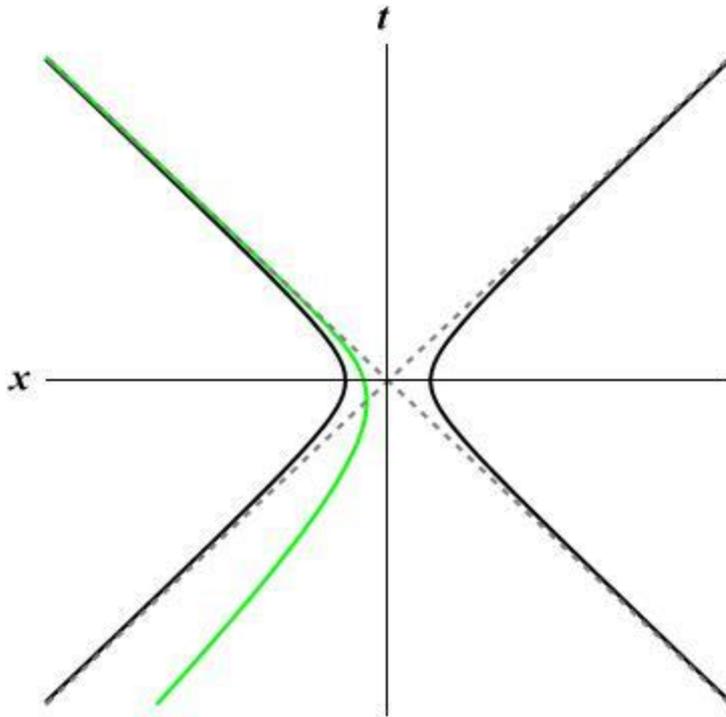
- Hawking \leftrightarrow Davies-Fulling
 - Thermal late times
- Unruh \leftrightarrow Carlitz-Willey
 - Thermal all times (exactly analytic)
- Black Hole \leftrightarrow Moving Mirror
 - All times. (even the non-thermal early times)

Thermal Equilibrium: Unruh Effect vs Thermal Mirror

Unruh world-line: Black

Thermal mirror: Green

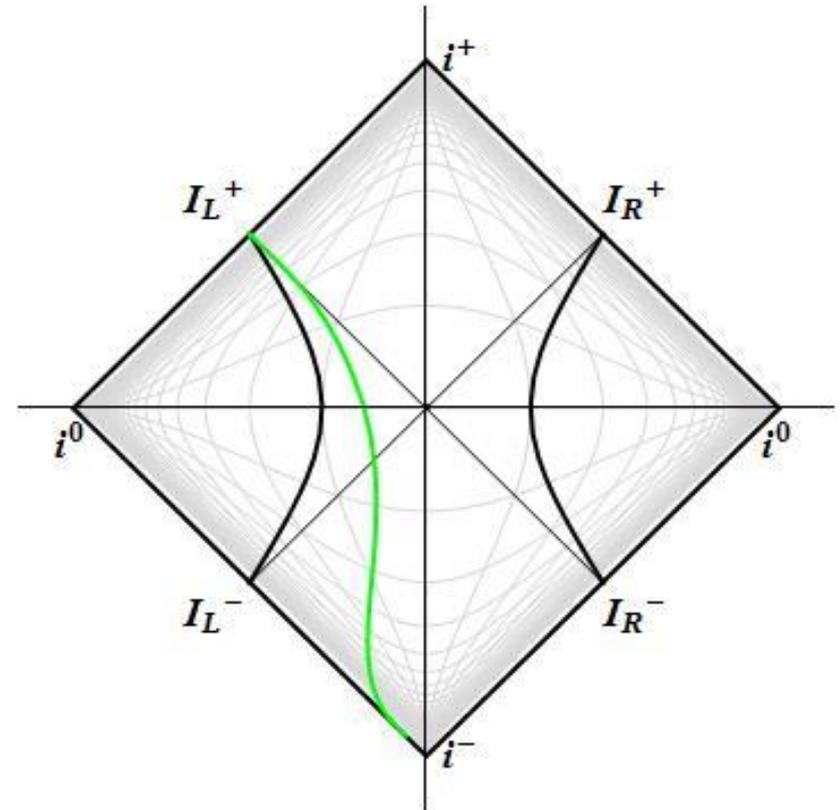
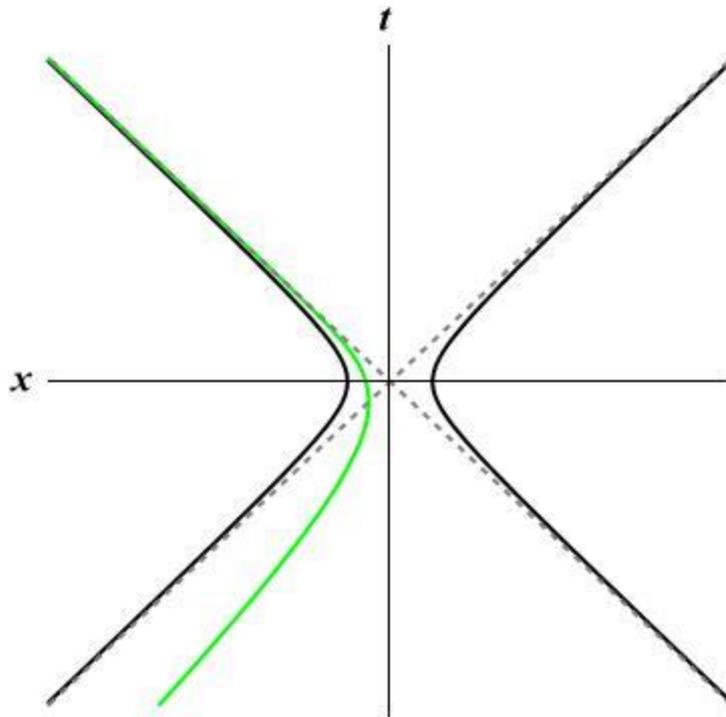
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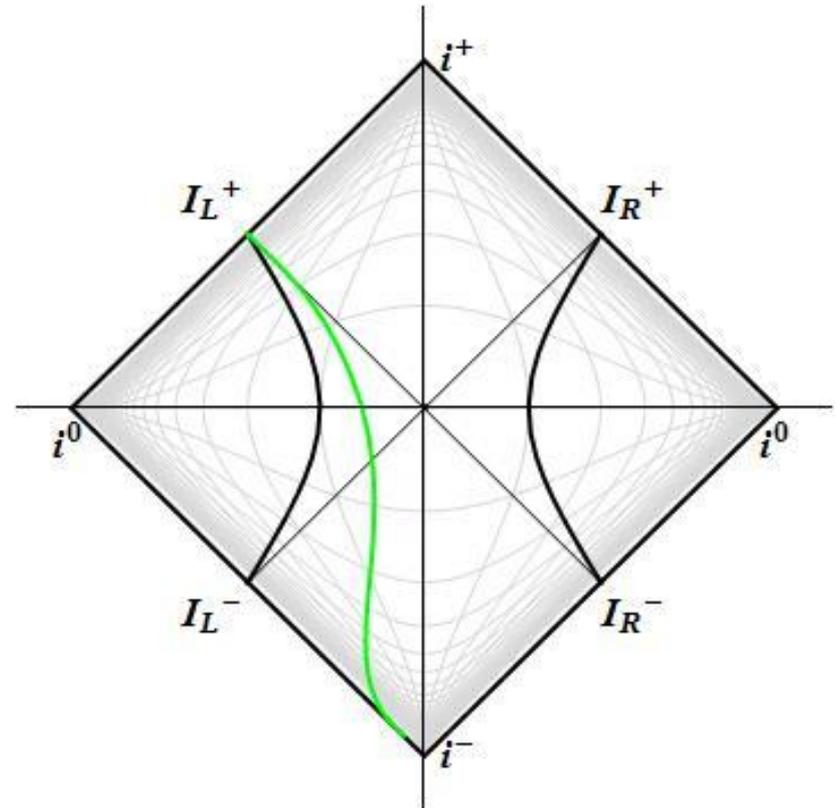
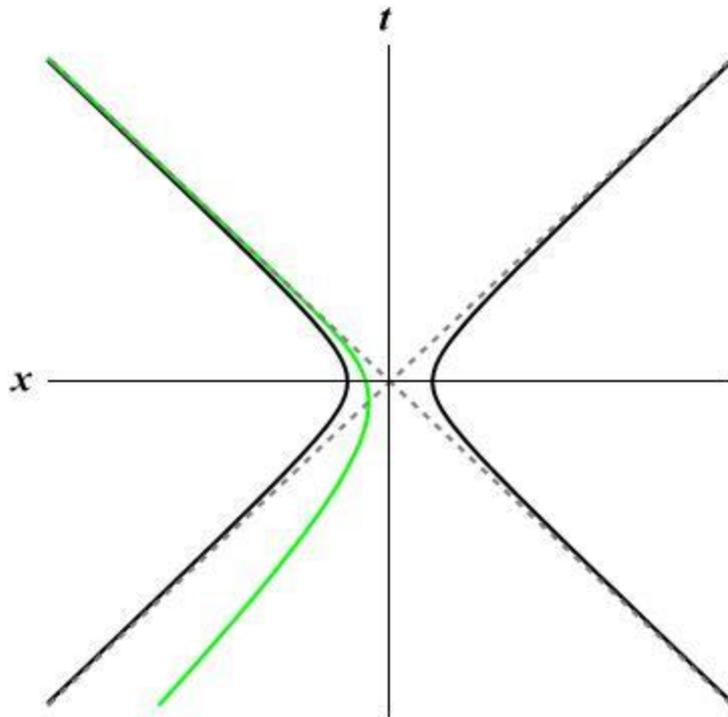
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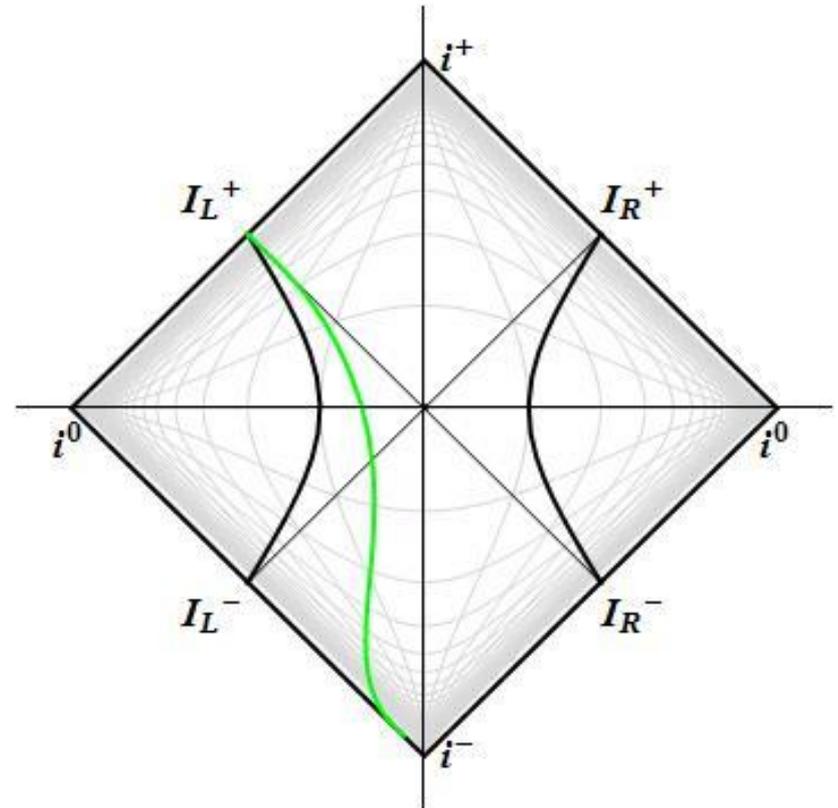
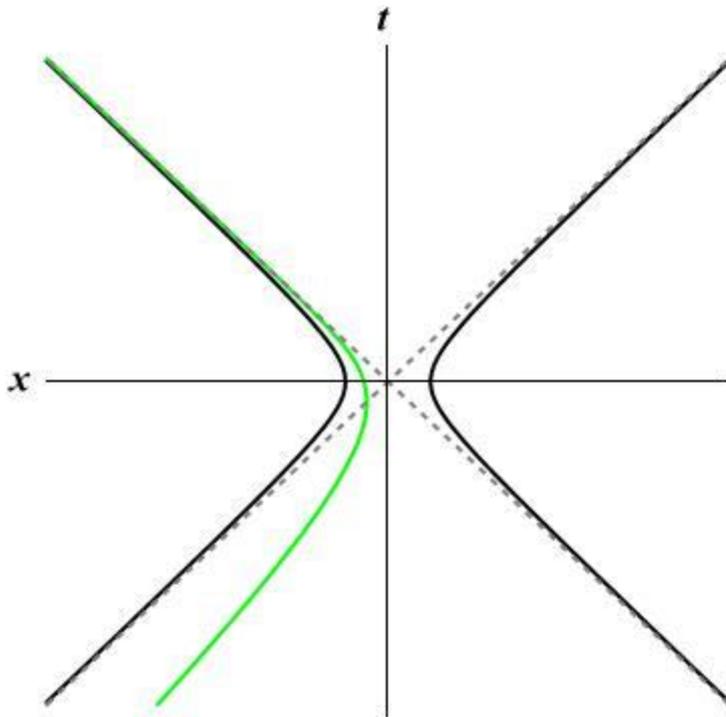
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$$z_U^2 = \kappa^{-2} + t^2$$

Thermal Equilibrium: Unruh Effect vs Thermal Mirror



Unruh world-line: Black

Thermal mirror: Green

$$z_U^2 = \kappa^{-2} + t^2$$

$$z_M = -t - \kappa^{-1} W(e^{-2\kappa t})$$

Product Log

$$xe^x = 1$$

Product Log

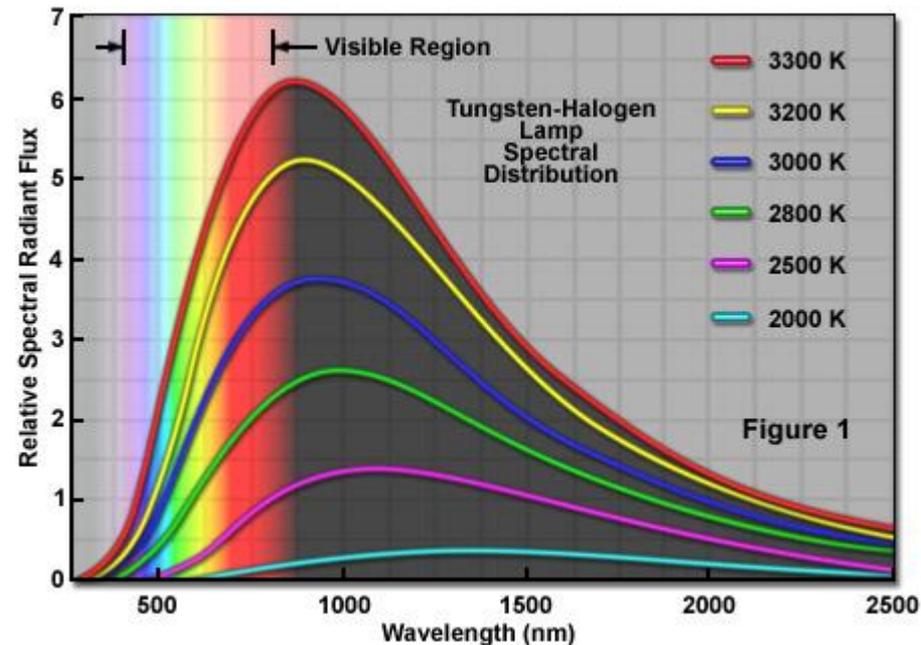
$$xe^x = 1$$

$$x = W(1) = 0.567$$

$$E(\omega)d\omega = \frac{V\hbar}{\pi^2 c^3} \frac{\omega^3}{e^{\beta\hbar\omega} - 1} d\omega$$

$$\frac{d}{d\omega} E(\omega) = 0$$

$$\omega_{\max} = \frac{3 + W(-3e^{-3})}{\beta\hbar}$$



$$\omega_{\max} = \zeta \frac{k_B T}{\hbar}$$

$$\zeta = 2.822$$

Thermal Accelerations

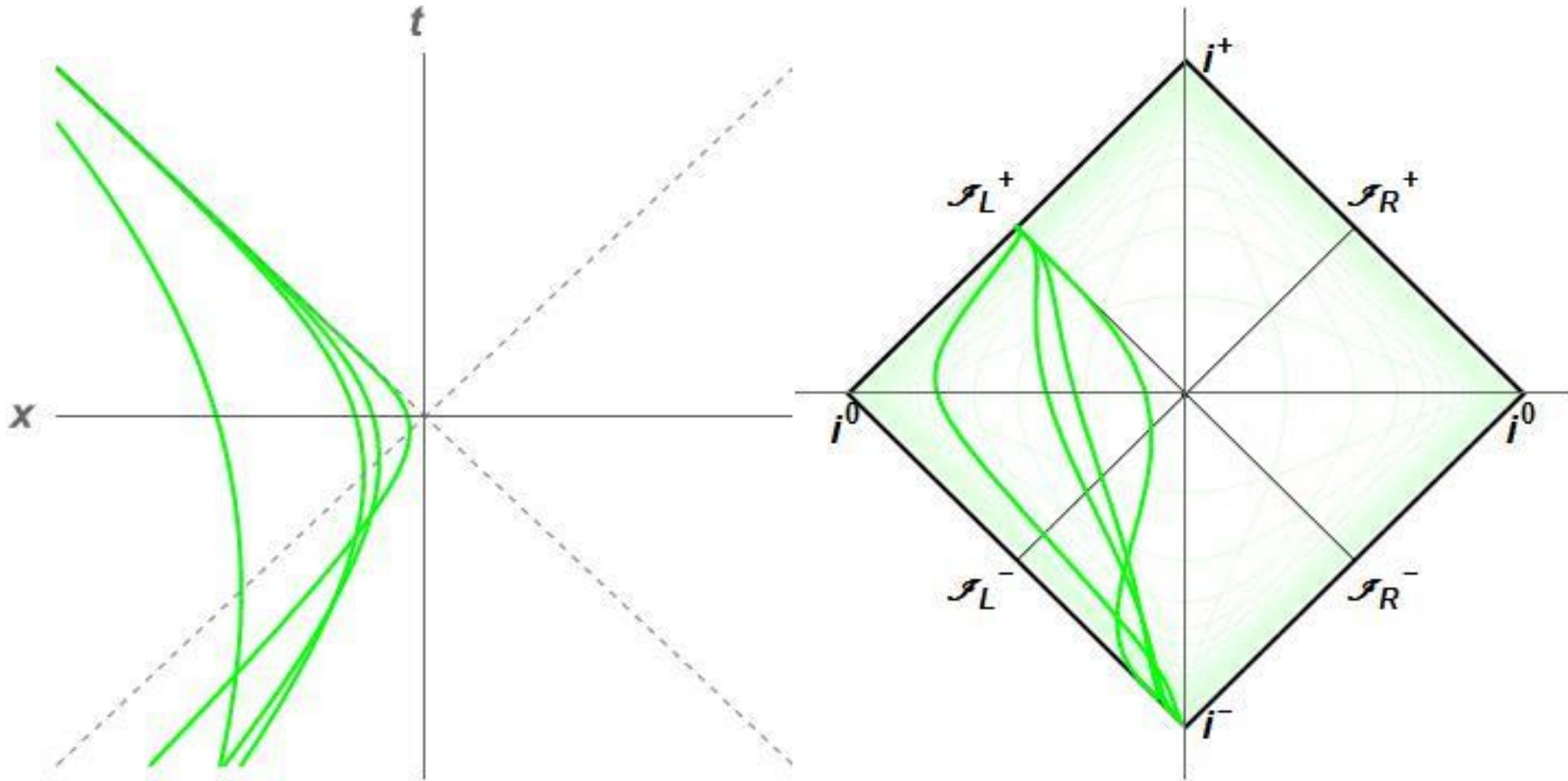
$$|\alpha_M| = \frac{\kappa}{2\sqrt{W}(e^{-2\kappa t})} \quad |\alpha_U| = \kappa$$

The Unruh effect and the thermal mirror trajectory give off constant energy flux at all times.

Both radiate particles thermally.

$$\frac{1}{e^{2\pi\omega/\kappa} - 1} \quad T = \frac{\kappa}{2\pi}$$

Eternal Thermal



Carlitz R.D. and R.S. Willey,
"Reflections on Moving Mirrors". Phys. Rev. D **36**, 2327, 1987

Localization

$$j\epsilon \leq \omega \leq (j+1)\epsilon$$

$$\omega_j = (j + 1/2)\epsilon$$

$$\frac{(2\pi n - \pi)}{\epsilon} \leq u \leq \frac{2\pi n + \pi}{\epsilon}$$

$$u_n = 2\pi n / \epsilon$$

$$\beta_{jn, \omega'} = \frac{1}{\sqrt{\epsilon}} \int_{j\epsilon}^{(j+1)\epsilon} d\omega e^{2\pi i \omega n / \epsilon} \beta_{\omega \omega'}$$

$$\langle in | N_{jn}^{out} | in \rangle = \int_0^\infty d\omega' |\beta_{jn, \omega'}|^2$$

Hawking, S.W.

“Particle Creation by Black Holes”, Comm. Math. Phys. 43, 1975.

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Hawking, S.W.

“Particle Creation by Black Holes”, Comm. Math. Phys. 43, 1975.

Time Independent Particle Creation

$$\langle N_{jn} \rangle = \frac{T}{\epsilon} \ln \left(\frac{e^{(j+1)\epsilon/T} - 1}{e^{j\epsilon/T} - 1} \right) - 1$$

Good, M.R.R., Anderson, P.R., and Evans, C.R.

“Time Dependence of Particle Creation from Accelerating Mirrors”, [arXiv:gr-qc/1303.6756](https://arxiv.org/abs/1303.6756)

Physical Review D, 88, 025023, July 2013.

Choices of Packets

$$\epsilon = 1$$

Giddings, Steven B.

“Hawking Radiation, the Stefan-Boltzmann Law, and unitarization”,
[arXiv:gr-qc/1511.08221](https://arxiv.org/abs/1511.08221) December 2015

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$$\epsilon = 1$$

$$\epsilon = \pi^{-1} = 0.32$$

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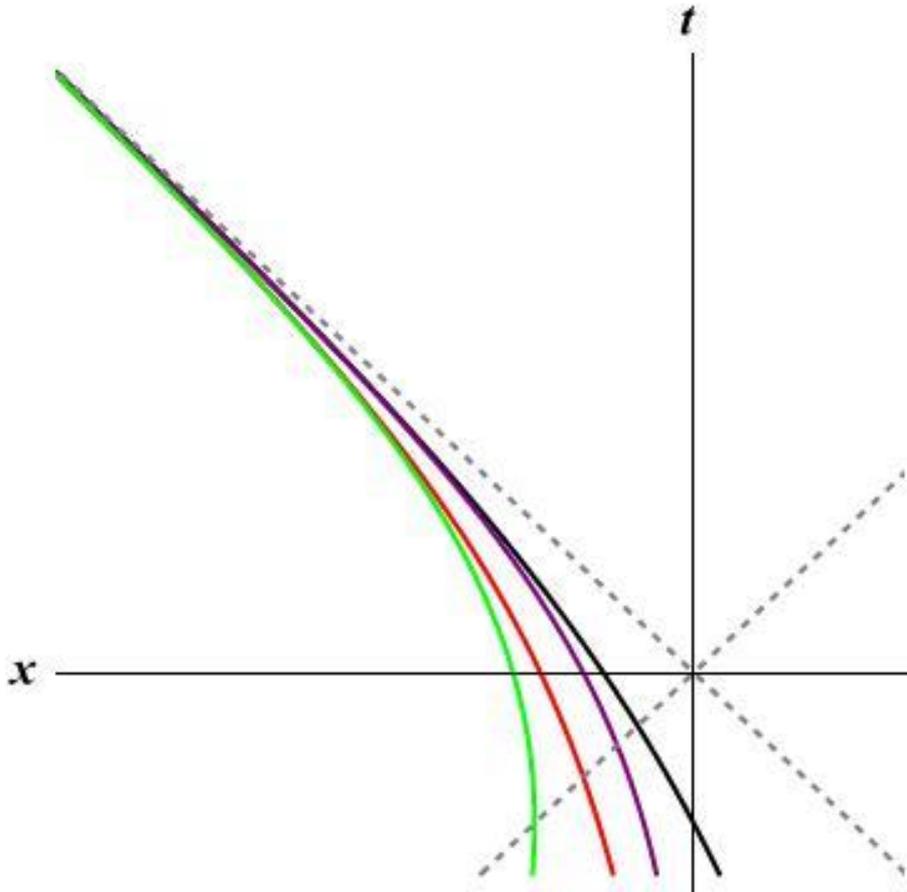
$$\epsilon = T \ln \phi = \frac{\kappa}{2\pi} \ln \frac{1+\sqrt{5}}{2} = 0.0765$$

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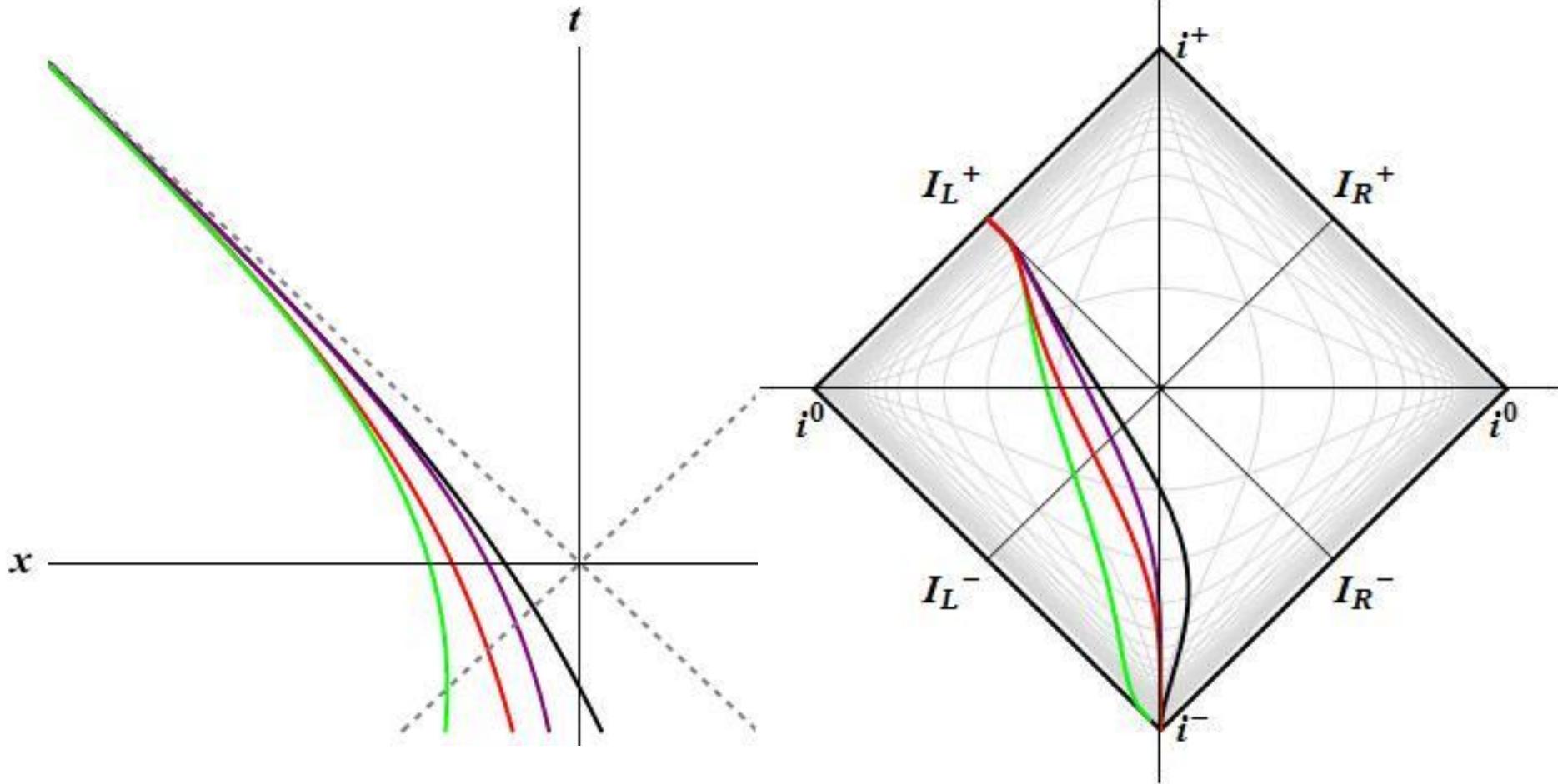
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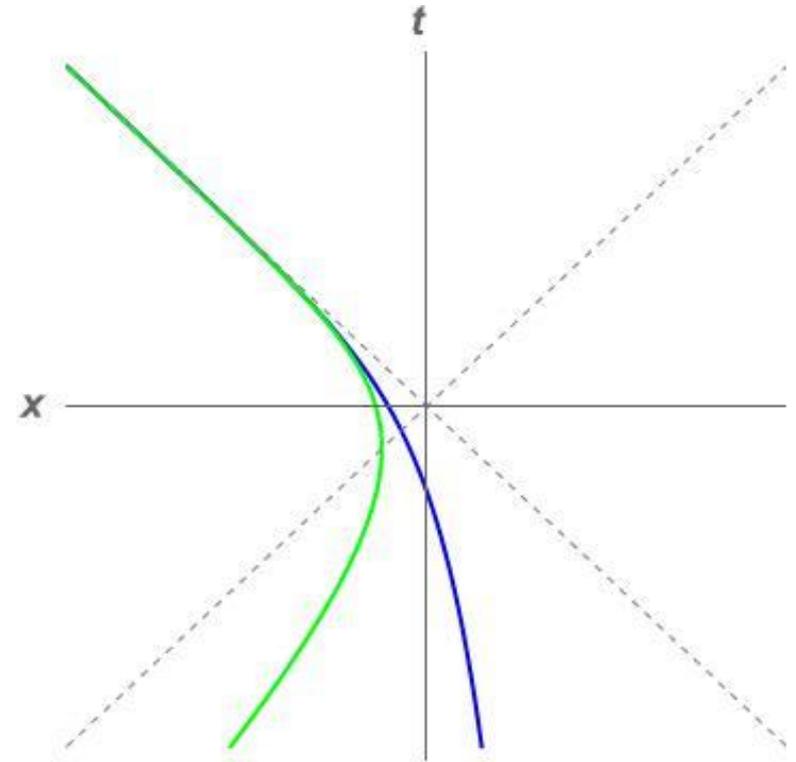
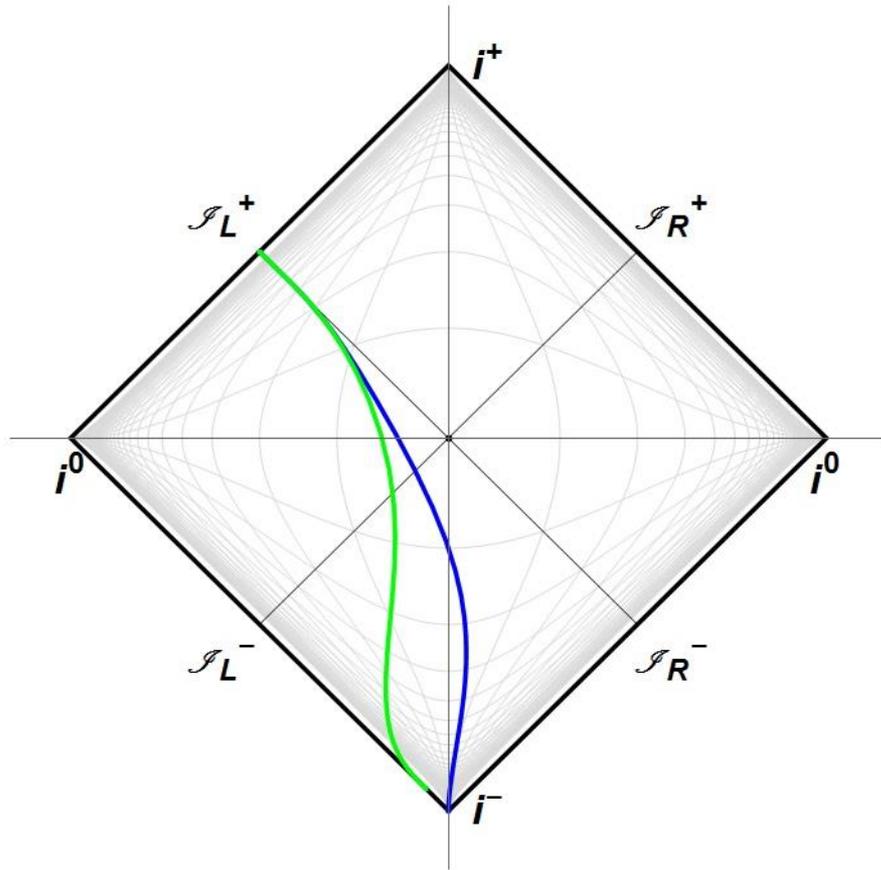
Evolution to Equilibrium



Evolution to Equilibrium



$$z(t) = -t - \frac{W[2e^{-2\kappa t}]}{2\kappa}$$



Good, M.R.R., Anderson, P.R., and Evans, C.R.

“Black Hole- Moving Mirror” , 1507.03489 & 1507.05048, 2015

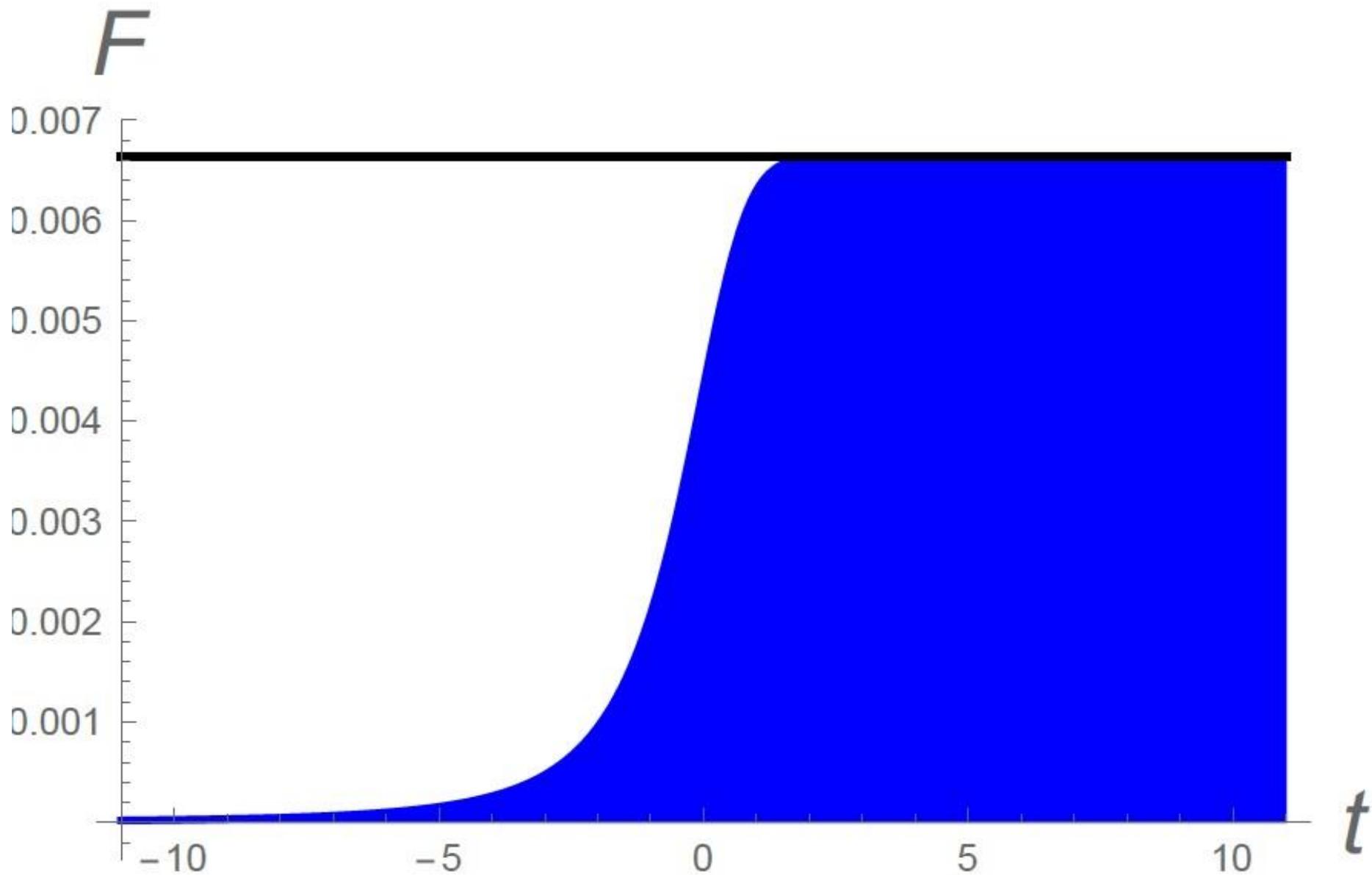
Coordinate Dependent Collapse

- Various horizon moving mirror solutions \Leftrightarrow various black hole collapse scenarios.
 - Non-stationary black hole mass function.
 - Different coordinates for geometry matching.

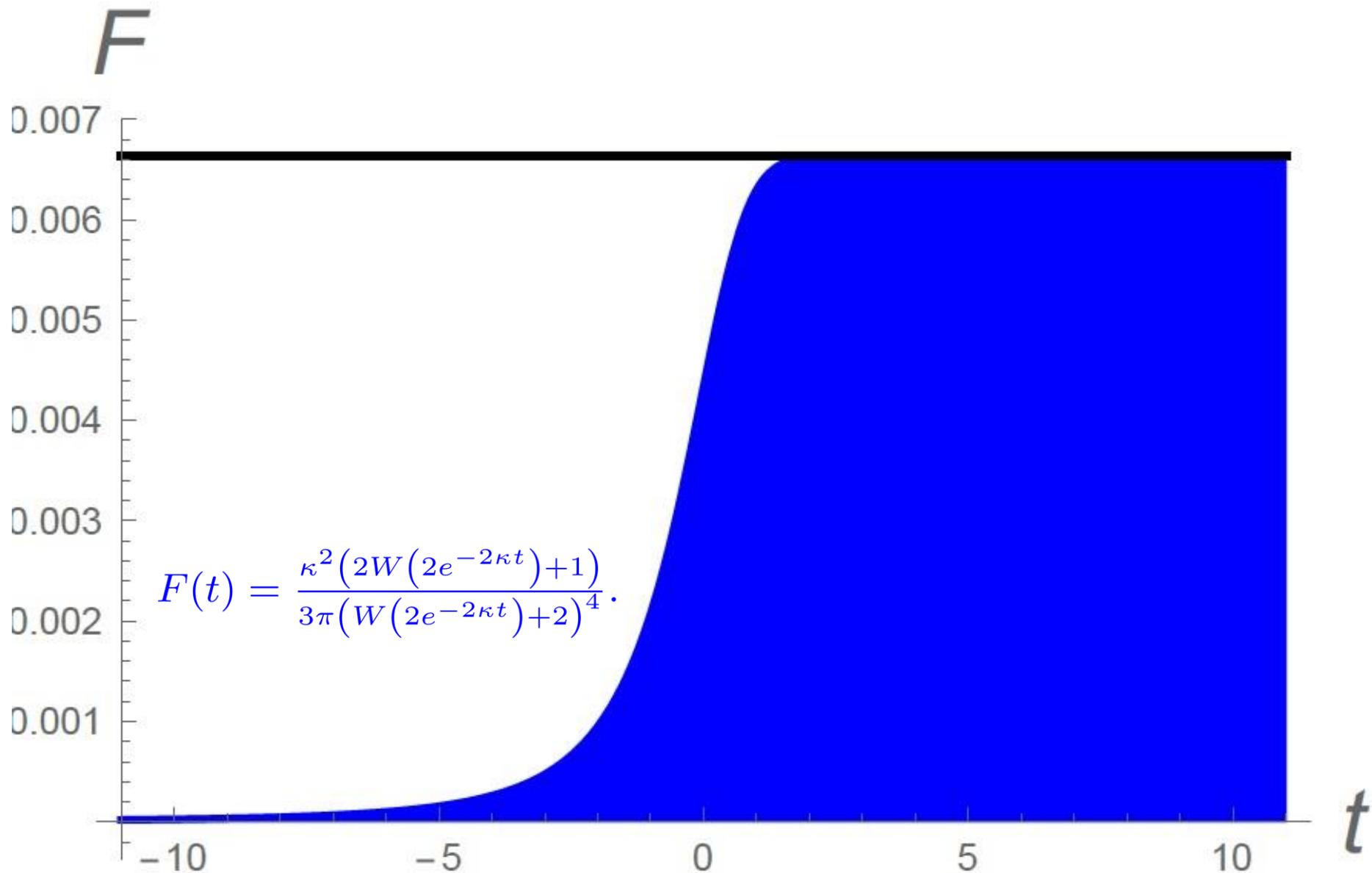
$$u_{\text{out}}(u_{\text{in}}) = u_{\text{in}} - 4M \ln \frac{|v_0 - 4M - u_{\text{in}}|}{4M}$$

$$r^* = r + 2M \ln \left| \frac{r}{2M} - 1 \right|$$

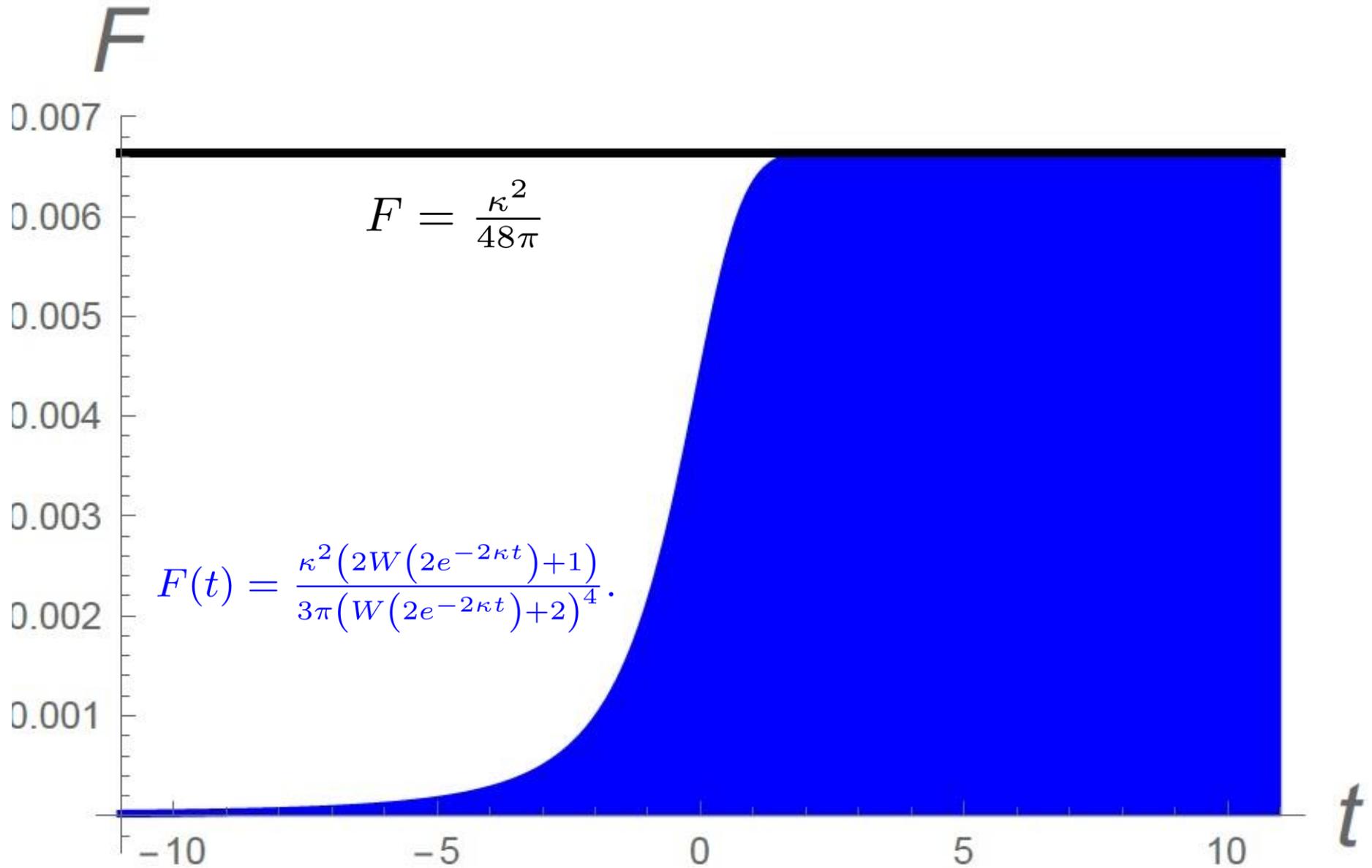
Energy Flux, $\kappa=1$.



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Non-eternally thermal Bogolubov Coefficients

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$$\beta_{\omega\omega'} = \frac{1}{4\pi\sqrt{\omega\omega'}} \frac{2\omega'}{\kappa} \Gamma\left(\frac{i\omega}{\kappa} + 1\right) \left(-\frac{i(\omega+\omega')}{\kappa}\right)^{-1-\frac{i\omega}{\kappa}}$$

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Late Times

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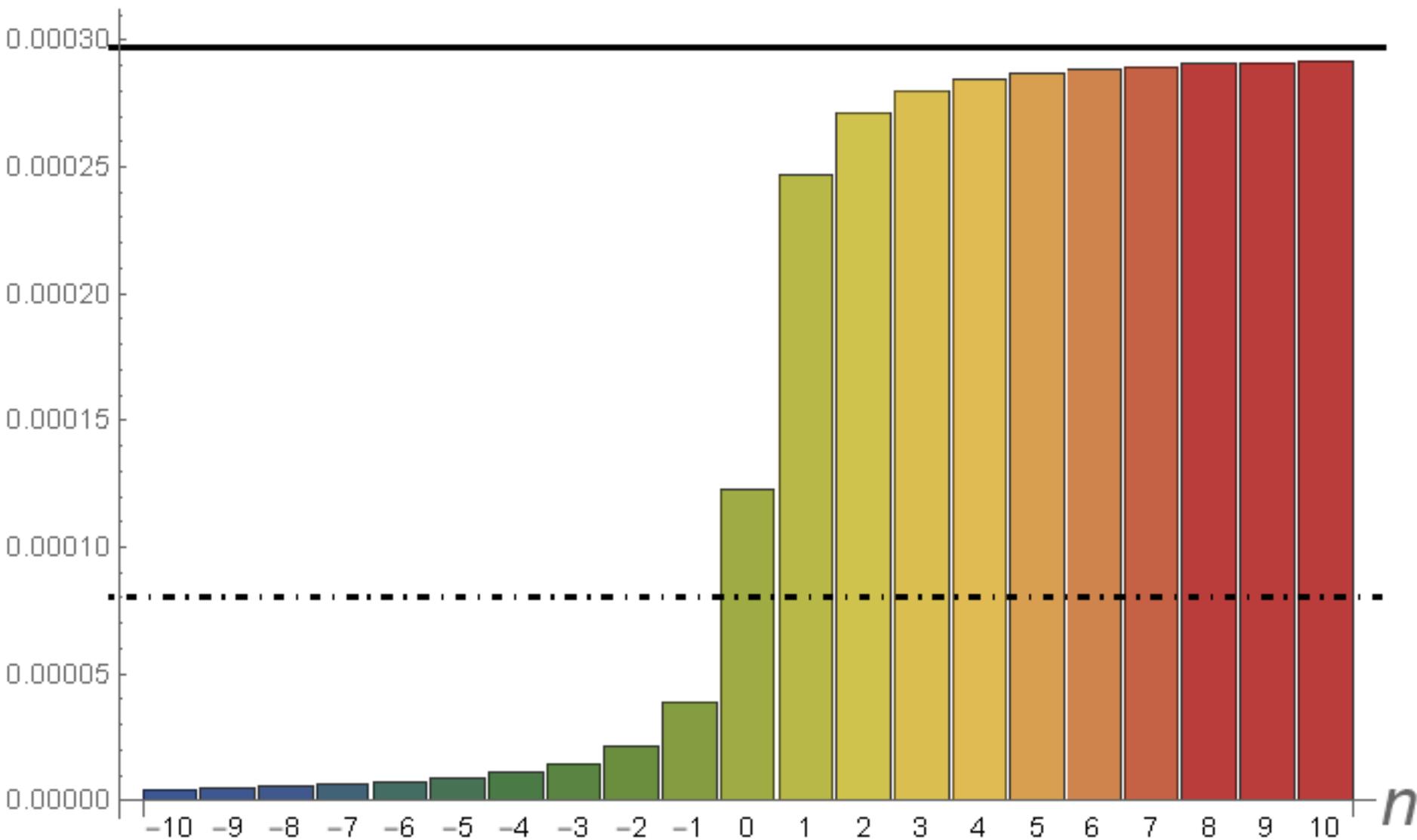
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$$T = \frac{\kappa}{2\pi} = \frac{1}{8\pi M}$$

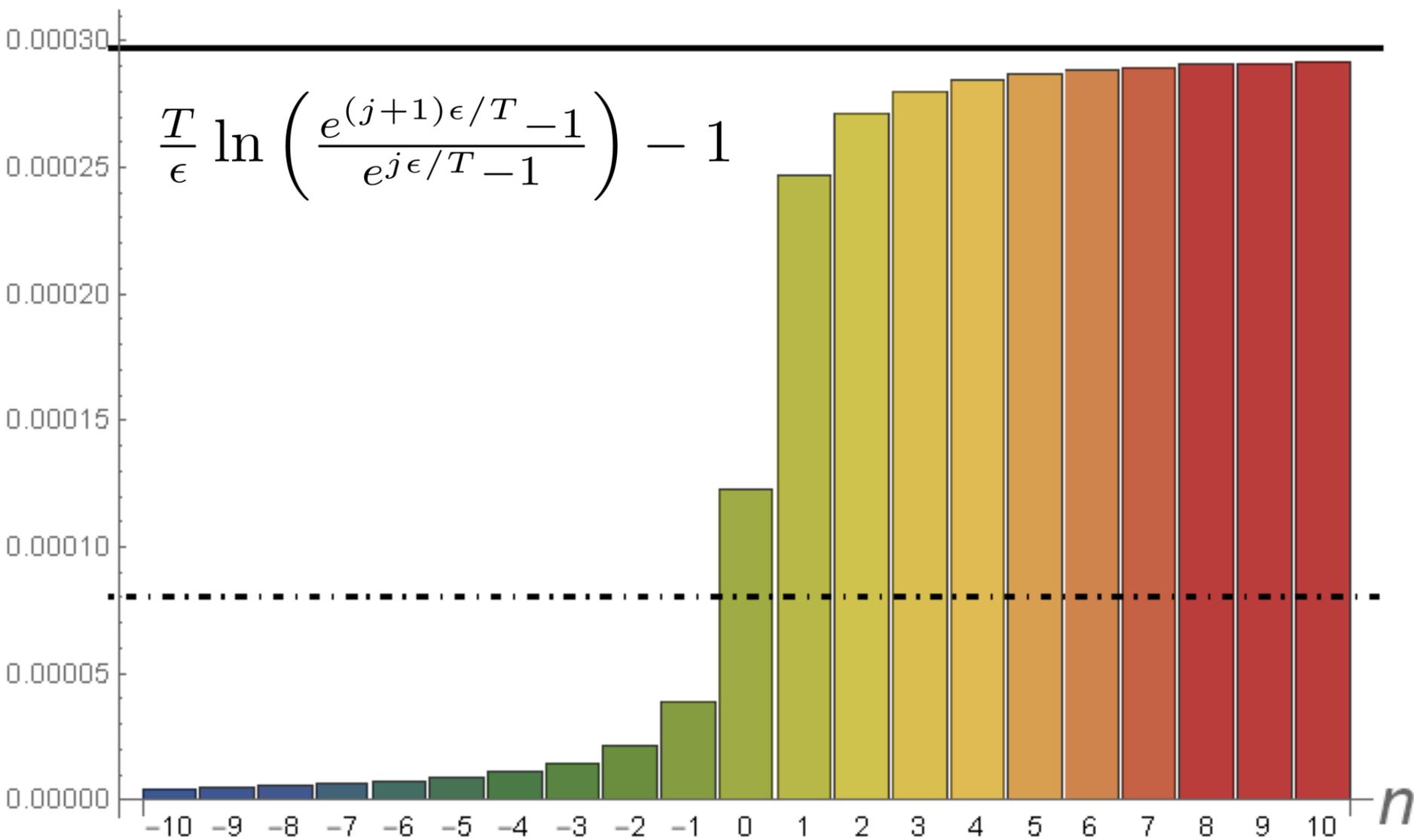
Particle Flux, $\epsilon=1, j=1, \kappa=1$.

N_{jn}



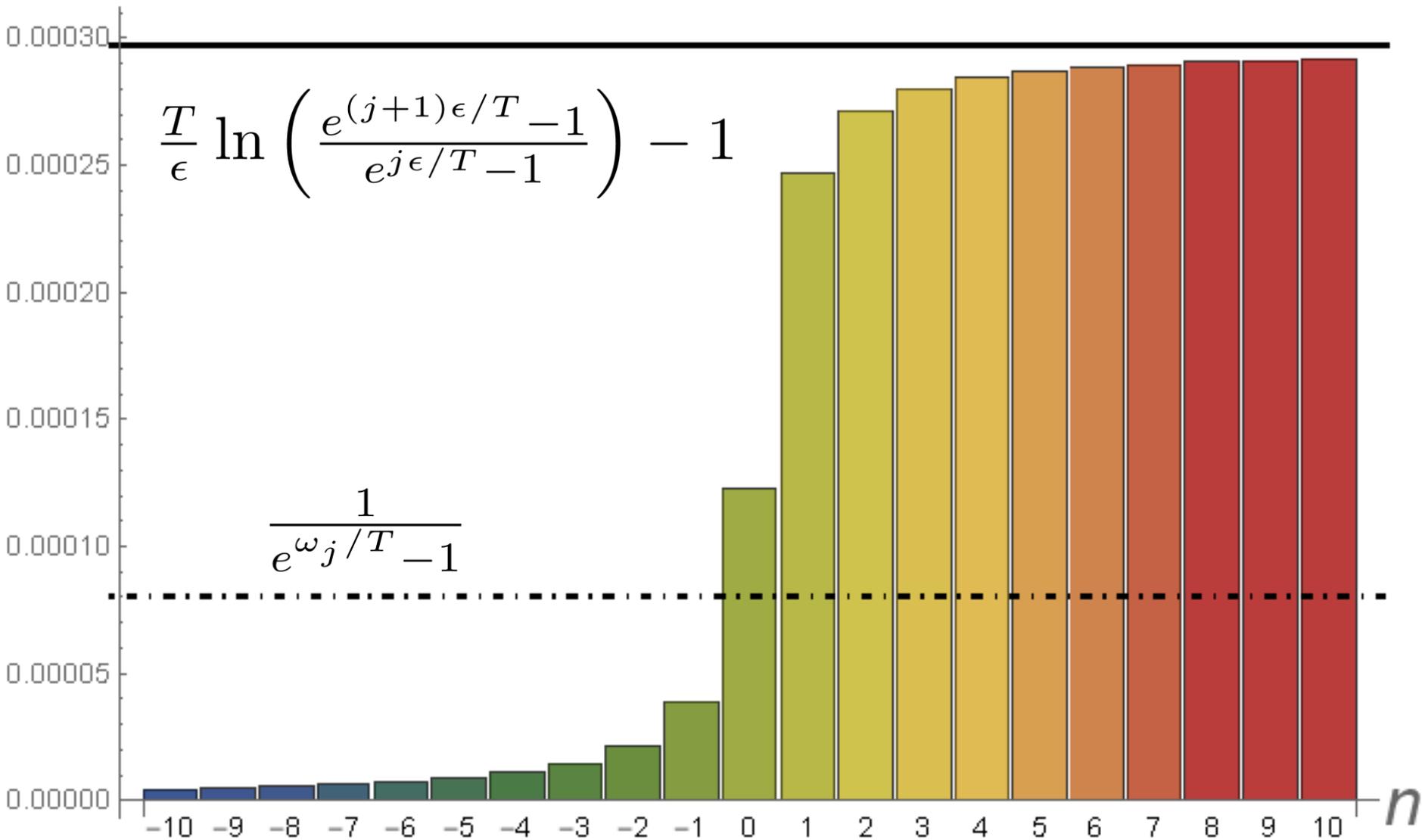
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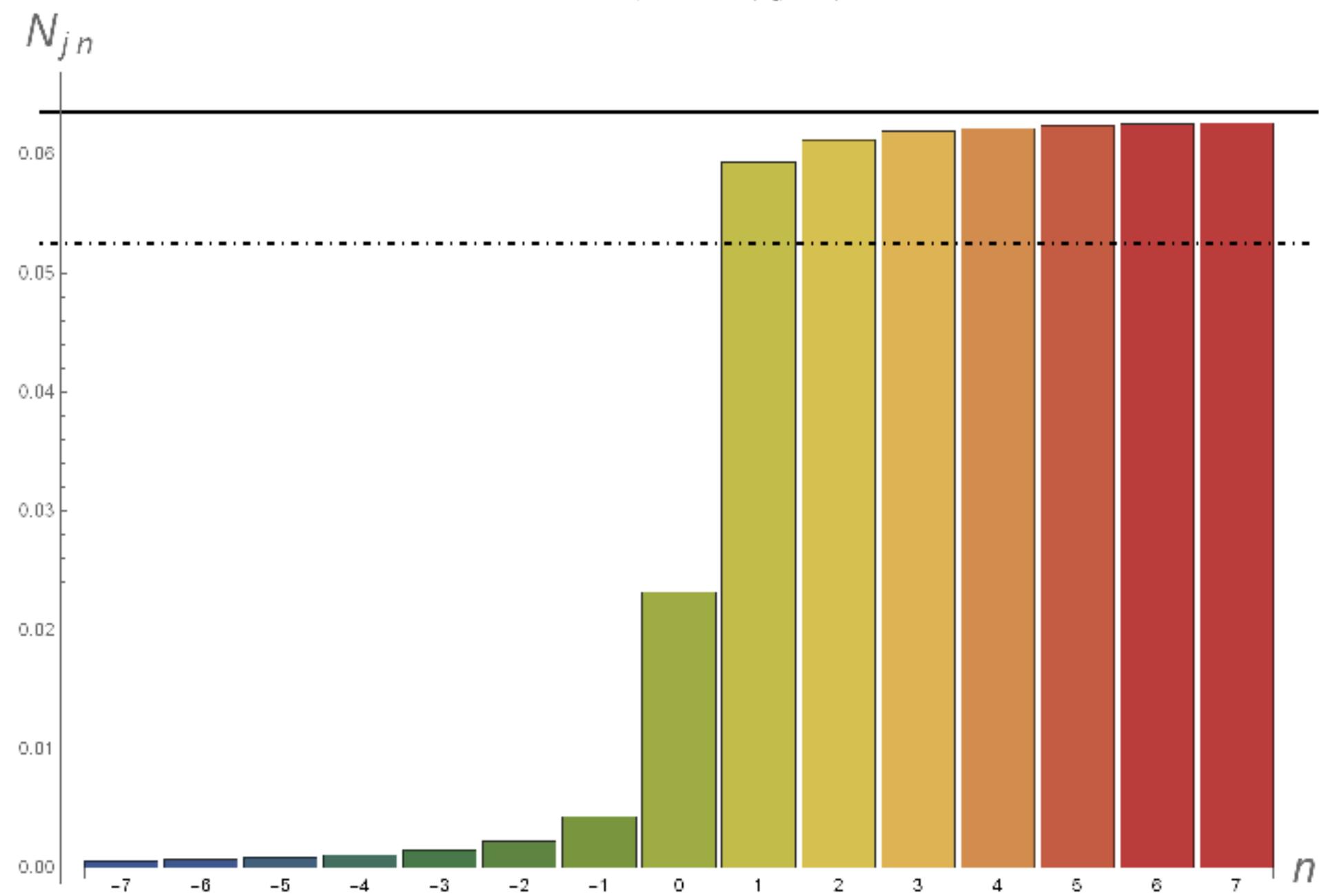


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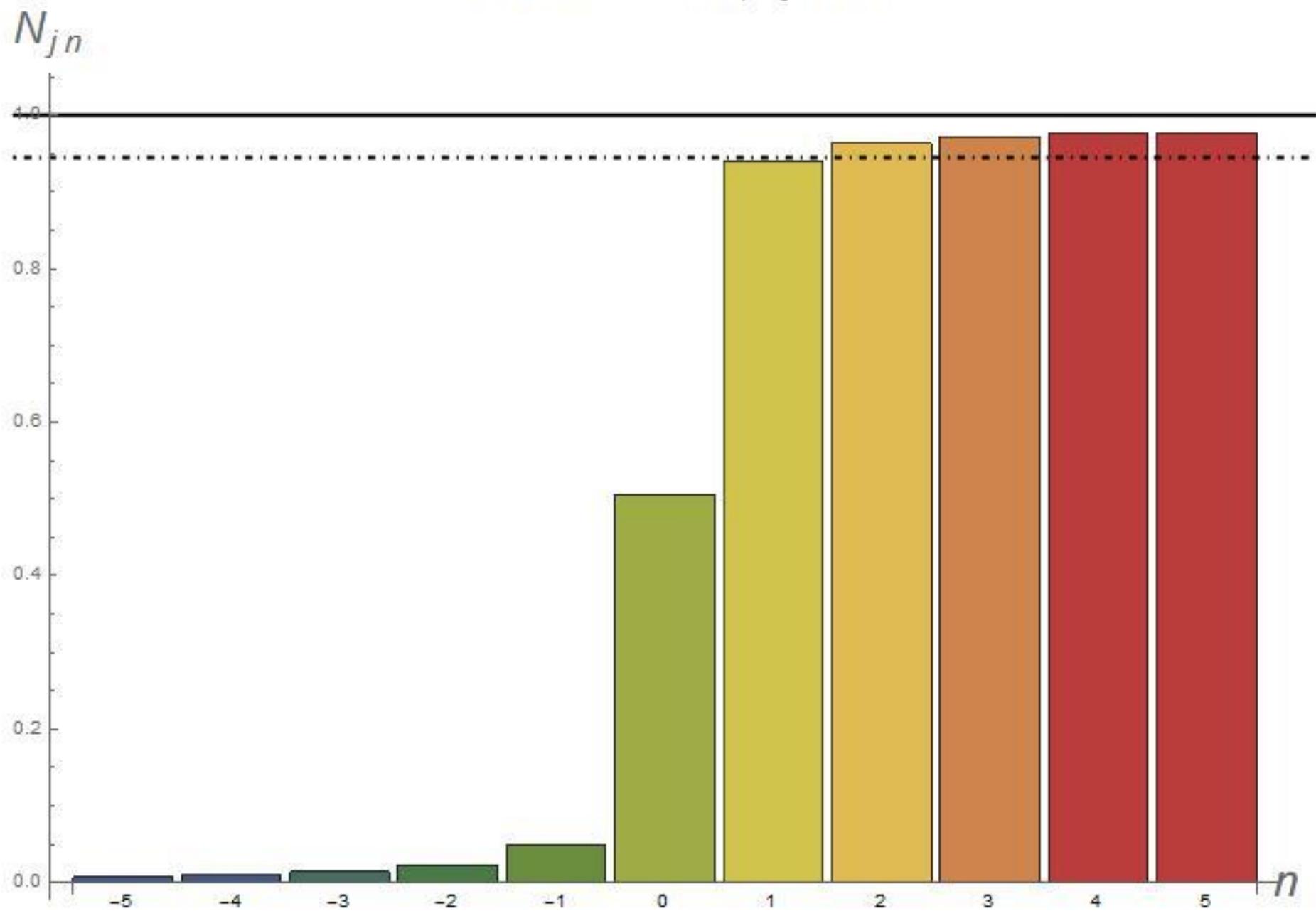
N_{jn}



Particle Flux, $\epsilon=\pi^{-1}$, $j=1$, $\kappa=1$.



Particle Flux, $2\pi\epsilon = \ln\phi$, $j=1$, $\kappa=1$.



Energy Conservation

- As it stands, the black hole radiates forever.
- However, the emission energy must be finite.

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$$E_T \leq M$$

Information Loss

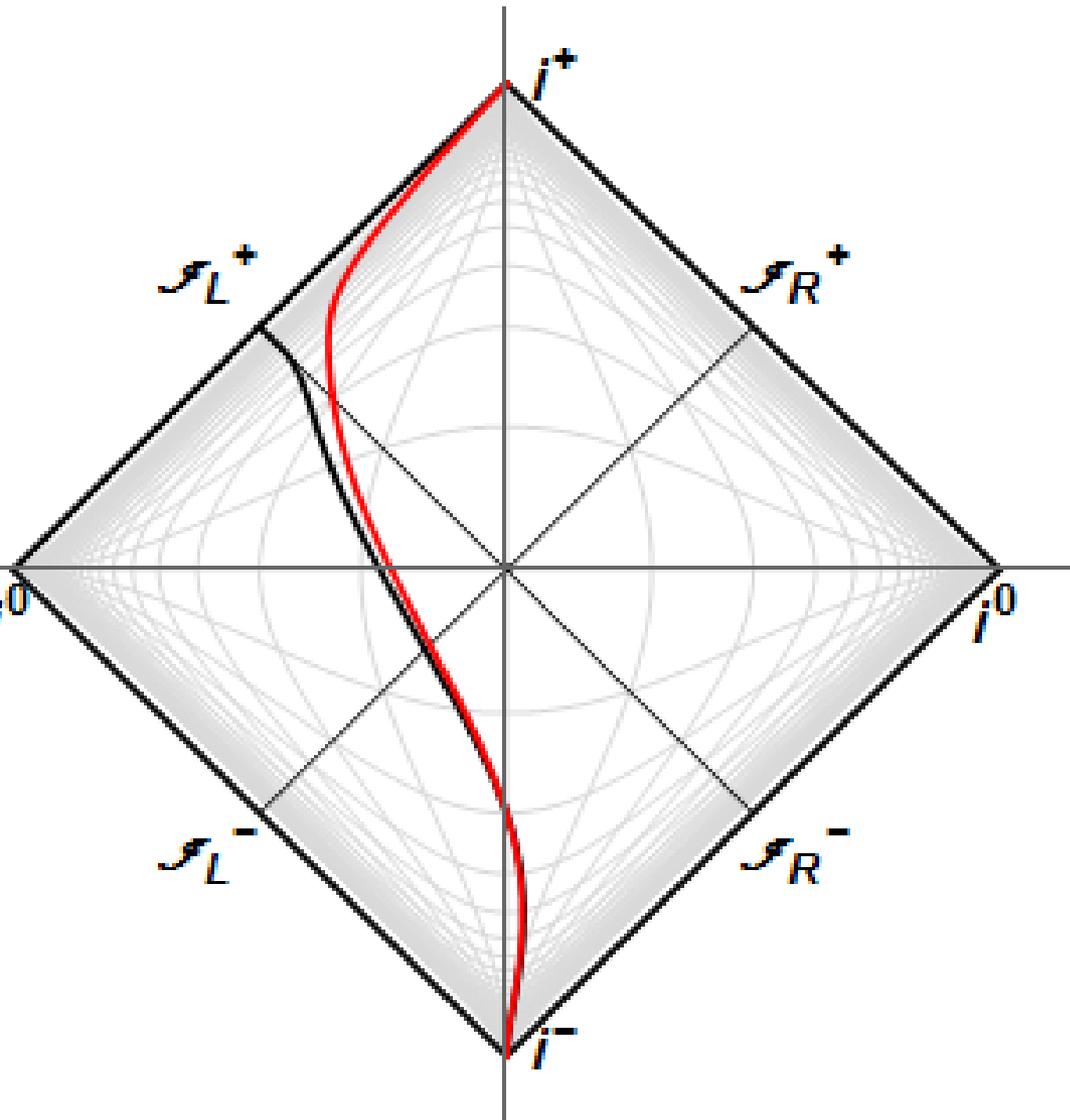
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- The mirror must be asymptotically inertial.
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Information Loss

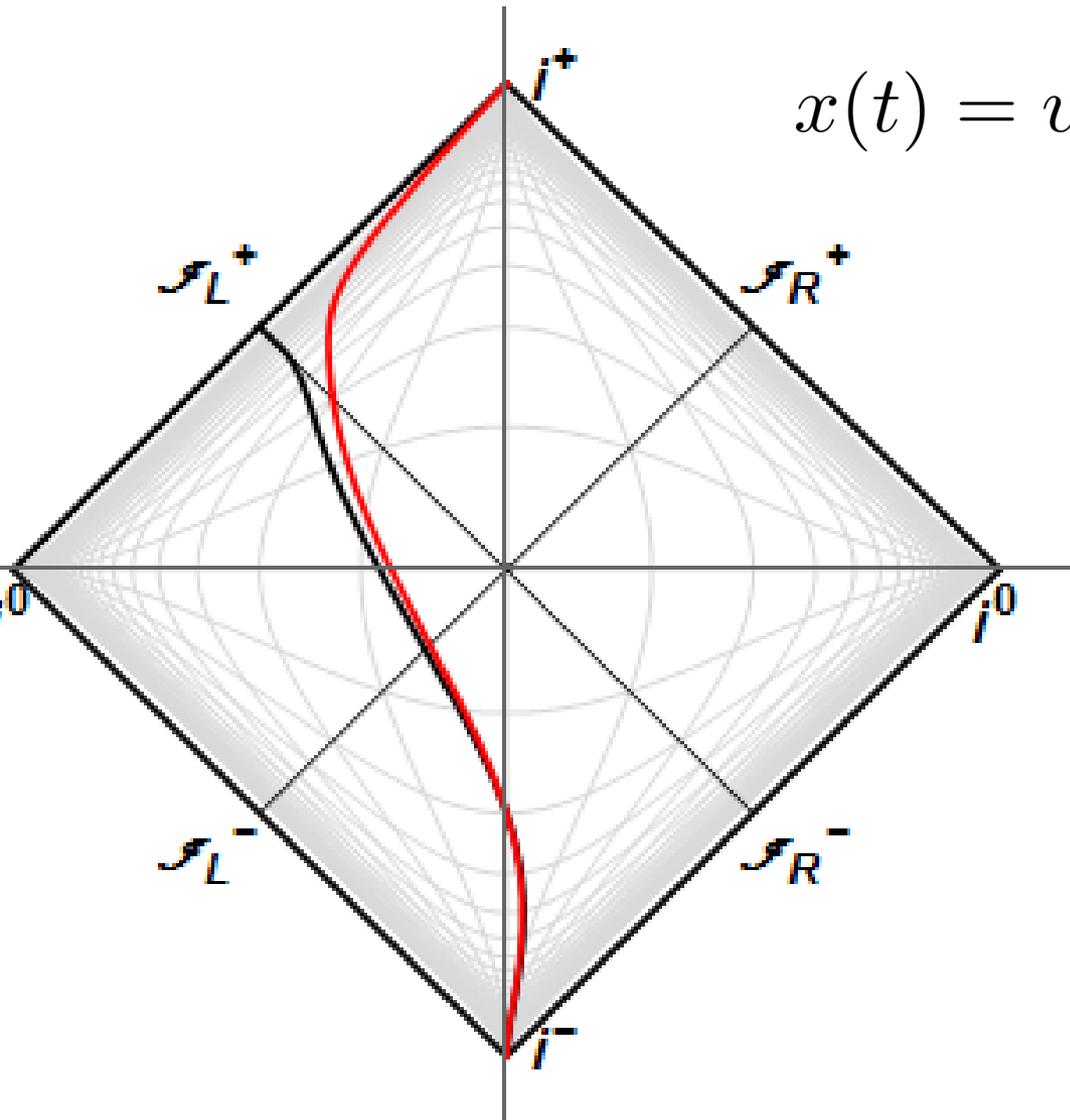
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$$12\pi F(u) = [\eta'(u)]^2 - \eta''(u)$$

Trajectory Functions

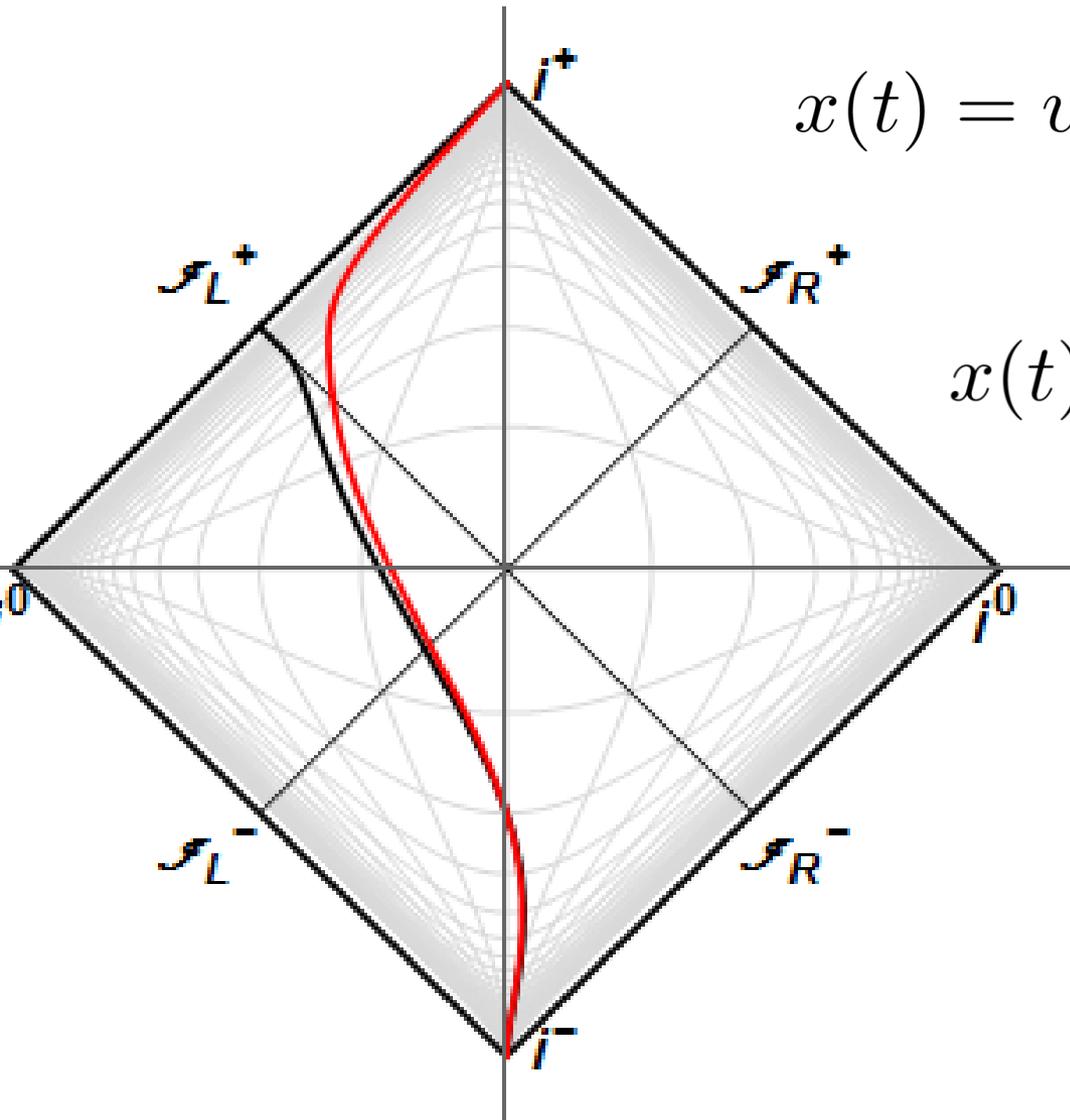


Trajectory Functions



$$x(t) = v_H - t - 2MW \left[2e^{\frac{v_H - t}{2M}} \right]$$

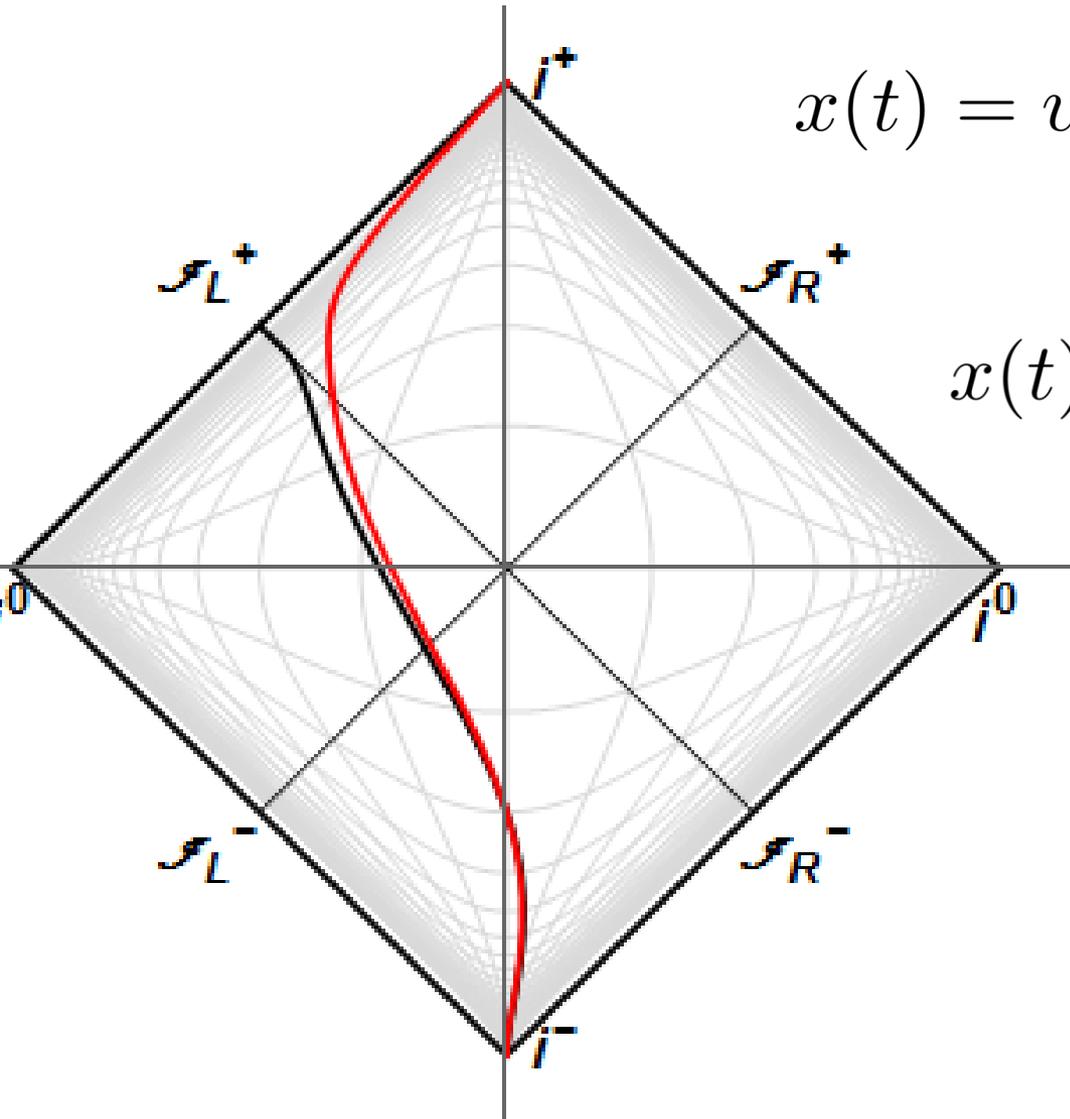
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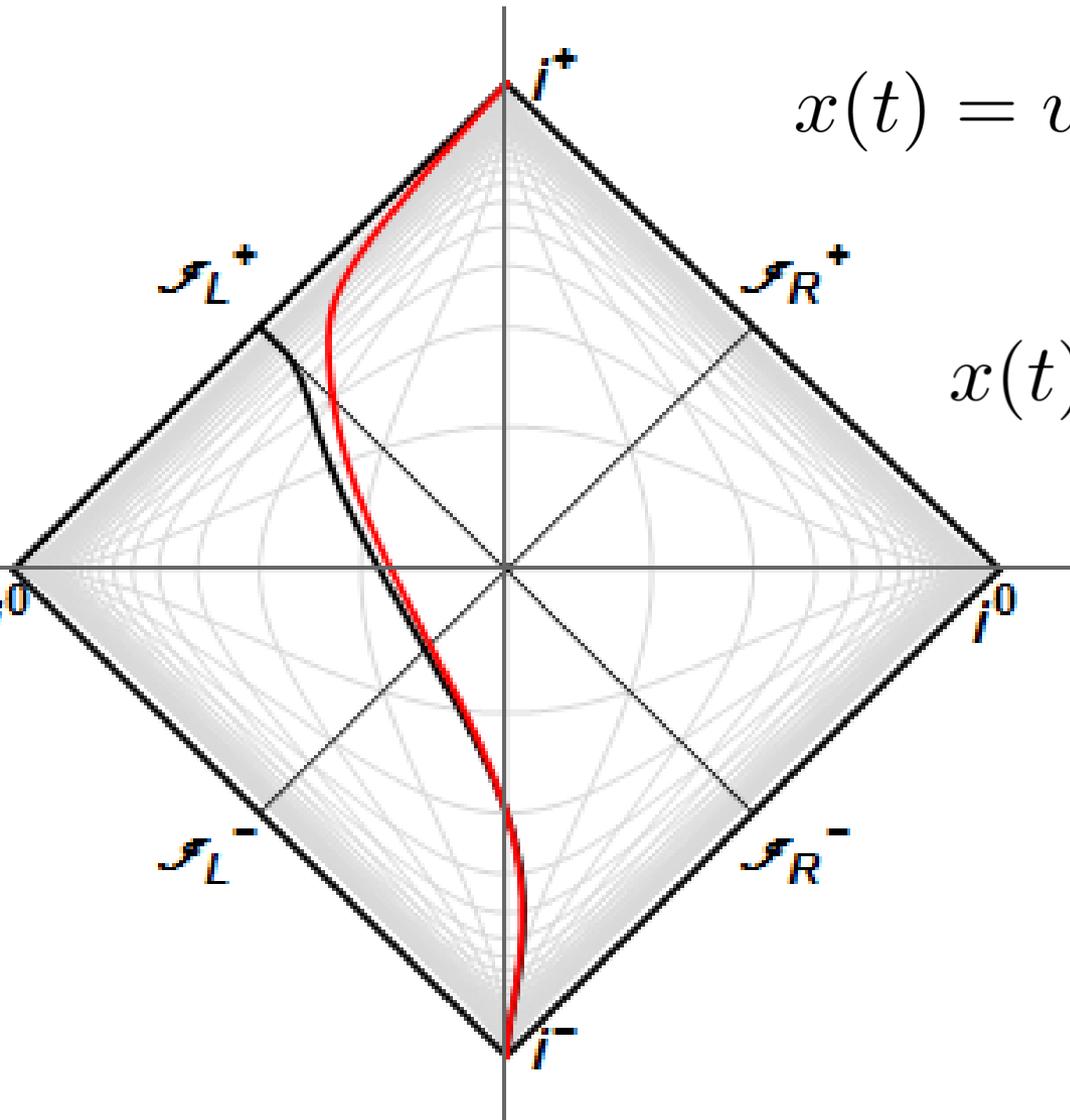


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Trajectory Functions



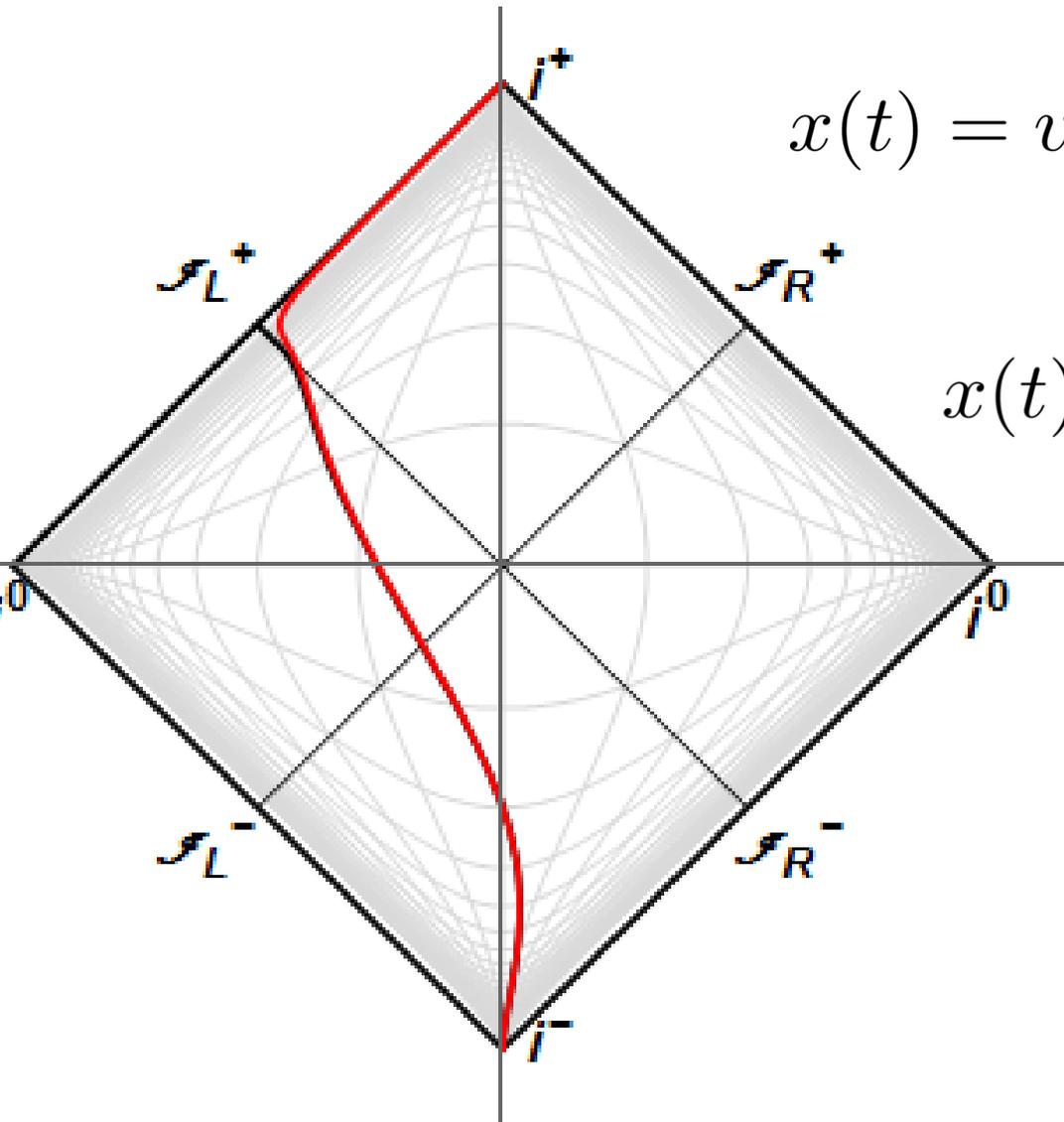
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$$\xi = 0.9$$

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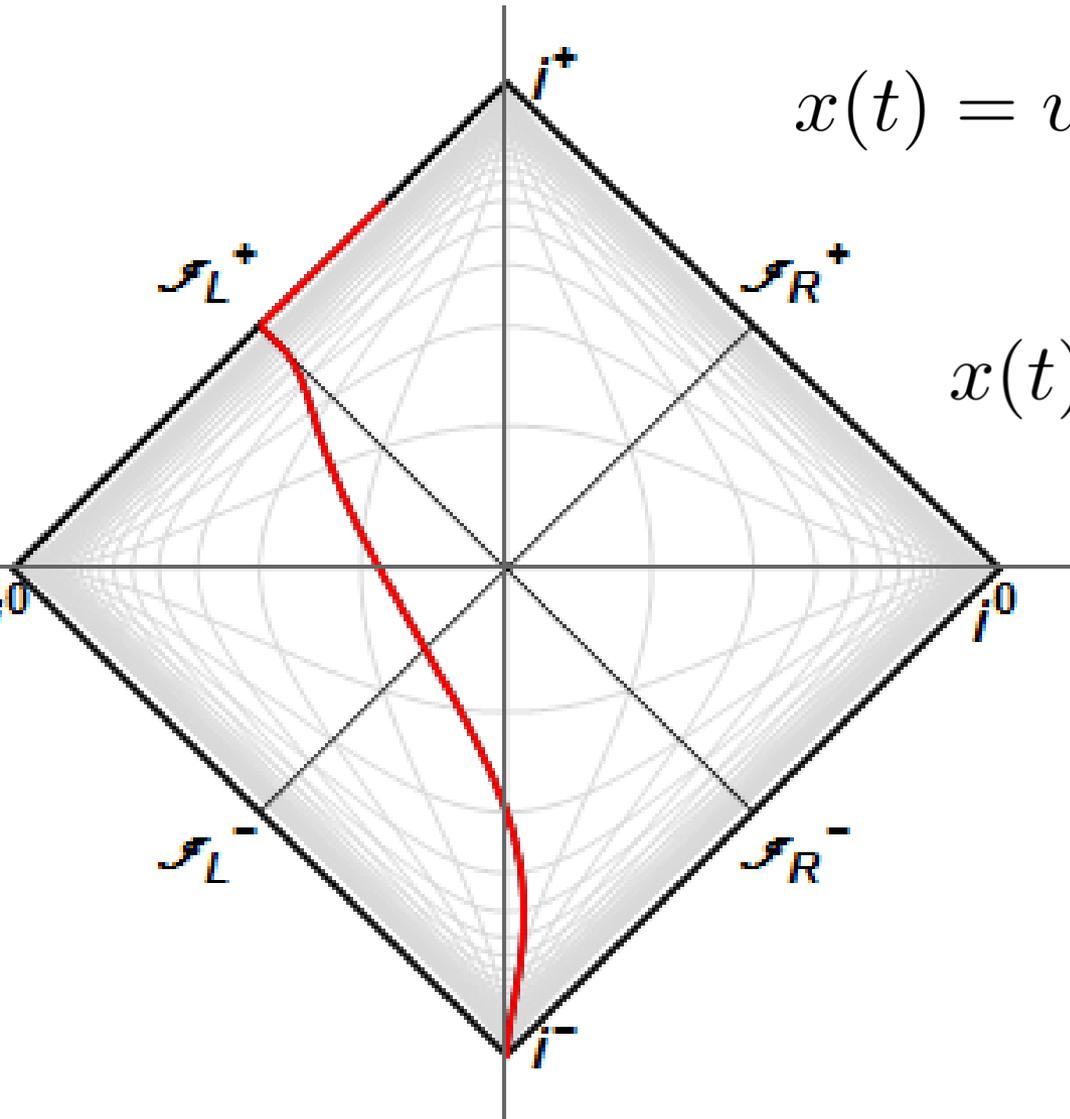
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$$\xi = 0.99$$

Trajectory Functions



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$$\xi = 0.9$$

$$\xi = 0.99$$

$$\xi = 0.999$$

Ray Tracing Functions

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$$u = v - 4M \ln \frac{v_H - v}{4M}$$

Ray Tracing Functions

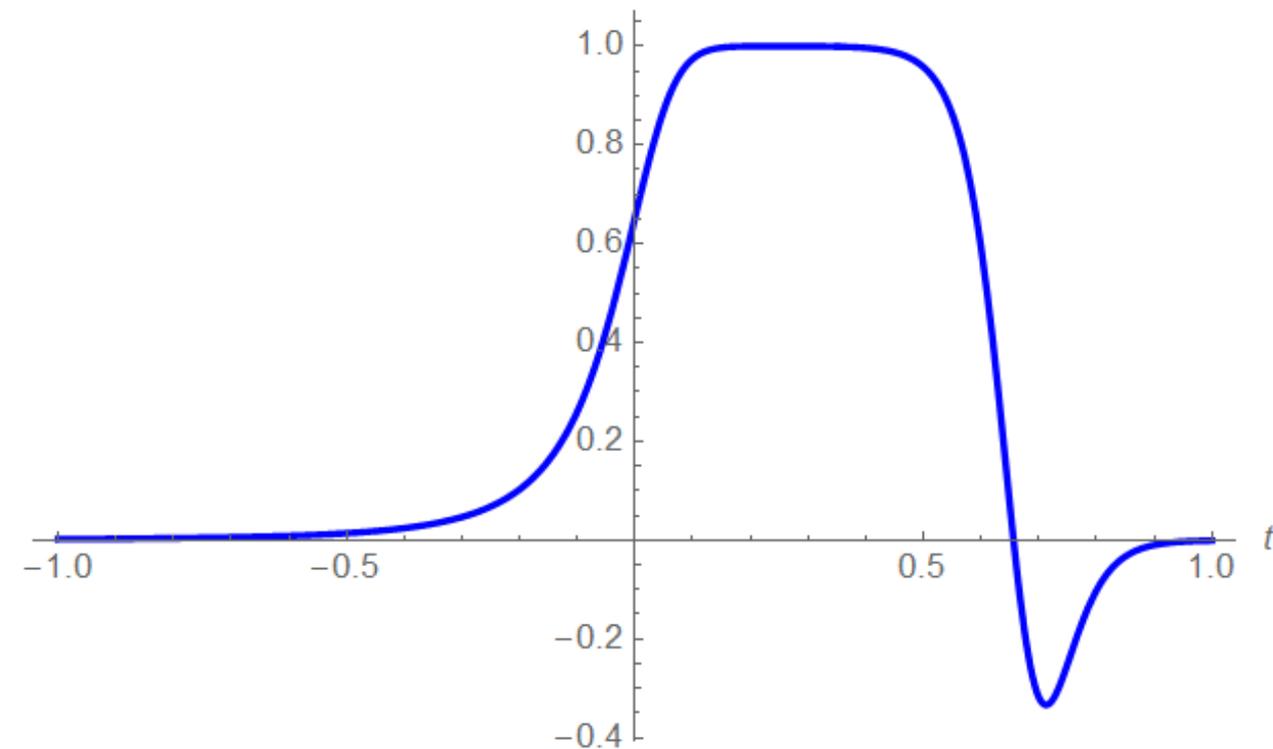
$$u = v - 4M \ln \frac{v_H - v}{4M}$$

$$u = \frac{1+\xi}{1-\xi} v + 4M\xi W \left[\frac{2e^{\frac{v_H - v}{2M(1-\xi)}}}{1-\xi} \right] - \frac{2\xi}{1-\xi} v_H$$

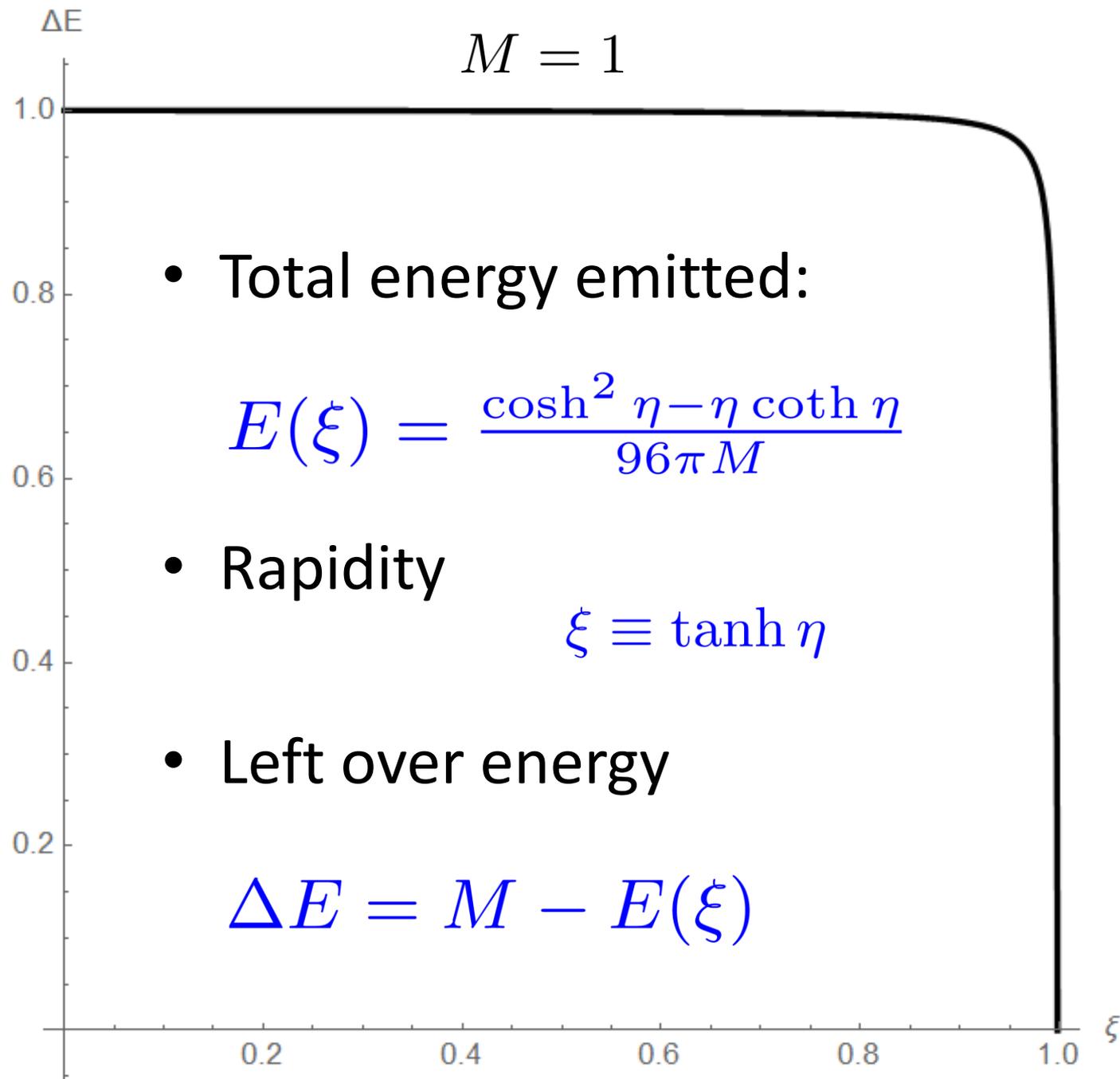
Energy Flux

$$F = \frac{\kappa^2 \xi W(2e^{-2\kappa t}) \left(\xi^2 + 2W(2e^{-2\kappa t})^2 + W(2e^{-2\kappa t}) - 1 \right)}{3\pi \left(-\xi + W(2e^{-2\kappa t}) + 1 \right)^2 \left(\xi + W(2e^{-2\kappa t}) + 1 \right)^4}$$

Energy Flux



$$\begin{aligned} \xi &= 1 - 0.1^7 \\ &= 0.99999999 \end{aligned}$$



Conclusions

- An exact one-to-one analytic correspondence exists.
- High-frequency limit gives thermal particle creation.
- Early-time particle creation is monotonic.
- Finite energy emission requires negative energy flux.
- Effects of acceleration \rightarrow Effects of gravitation.



$$e^{\epsilon/T} = \phi \quad \rightarrow \quad \langle N_{j=1} \rangle = 1$$

$$E_{st} = \int_{-\infty}^{\infty} \langle T_{uu} \rangle du$$

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$$E_{qs} = \int_0^{\infty} \omega \langle N_{\omega} \rangle d\omega$$

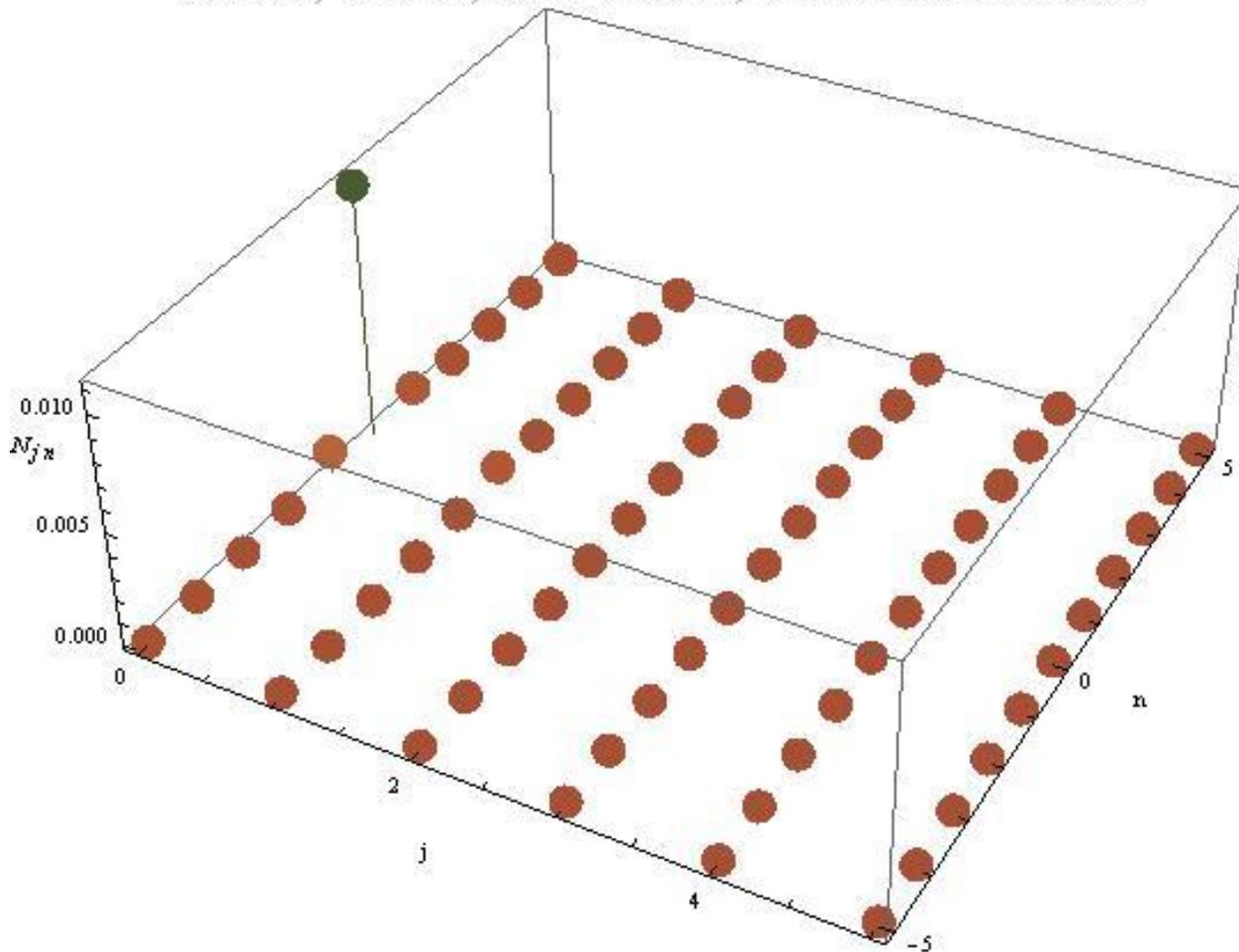
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$$E_{st} = E_{qs}$$

Spectral Resolution Using Packets

ArctX , $\epsilon=0.8$, $N = 0.0134$, 0.1% relative error



Wavepacket Spectral Resolution

$$\Delta t \Delta \omega \geq \frac{1}{2}$$

$$\Delta T \Delta \epsilon \geq 2\pi$$

insights

- Effects of acceleration will yield insights into the effects of gravitation.
- Existence of trajectories with spectral evolution.
- Moving mirrors provide a simple laboratory to construct the mathematical machinery needed to understand the time evolution of black hole evaporation.

Negative Energy Flux

Constraints on negative energy fluxes

Ford - Physical Review D, 1991

Negative Energy Seen By Accelerated Observers

Ford and Roman - Phys. Rev. D, 2013

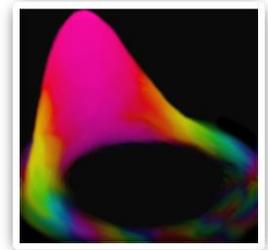
Energy flux correlations and moving mirrors

Ford and Roman - Phys.Rev. D, 2004

Finite Particle Count via Packets

$$\langle in | N_{\omega}^{out} | in \rangle = \frac{1}{e^{2\pi\omega_1/\kappa} - 1} \delta(\omega_1 - \omega_2)$$

$$\beta_{jn, \omega'} = \frac{1}{\sqrt{\epsilon}} \int_{j\epsilon}^{(j+1)\epsilon} d\omega e^{2\pi i \omega n / \epsilon} \beta_{\omega \omega'}$$



$$\langle in | N_{jn}^{out} | in \rangle = \int_0^{\infty} d\omega' |\beta_{jn, \omega'}|^2 = \frac{1}{e^{2\pi\omega_j/\kappa} - 1}$$

$$T = \frac{\kappa}{2\pi}$$

$$\langle T_{uu} \rangle = \frac{\kappa^2}{48\pi}$$

W Lambert (Product Log)

$$Y = Xe^X$$

$$3^t = 8t$$

$$X = W(Y)$$

$$-\frac{\ln 3}{8} = (-t \ln 3)e^{(-t \ln 3)}$$



$$t = -\frac{W\left(-\frac{\ln 3}{8}\right)}{\ln 3}$$

$$t \approx 0.146891$$

$$W(z) = \sum_{n=1}^{\infty} (-n)^{n-1} \frac{z^n}{n!}$$

We use this function for the transcendental inversion necessary to obtain $z(t)$.

Constant Energy Flux Trajectory

$$z = -t - \frac{1}{\kappa} W(e^{-2\kappa t})$$

$$\alpha = \gamma^3 a = \frac{-\kappa}{2\sqrt{W(e^{-2\kappa t})}}$$

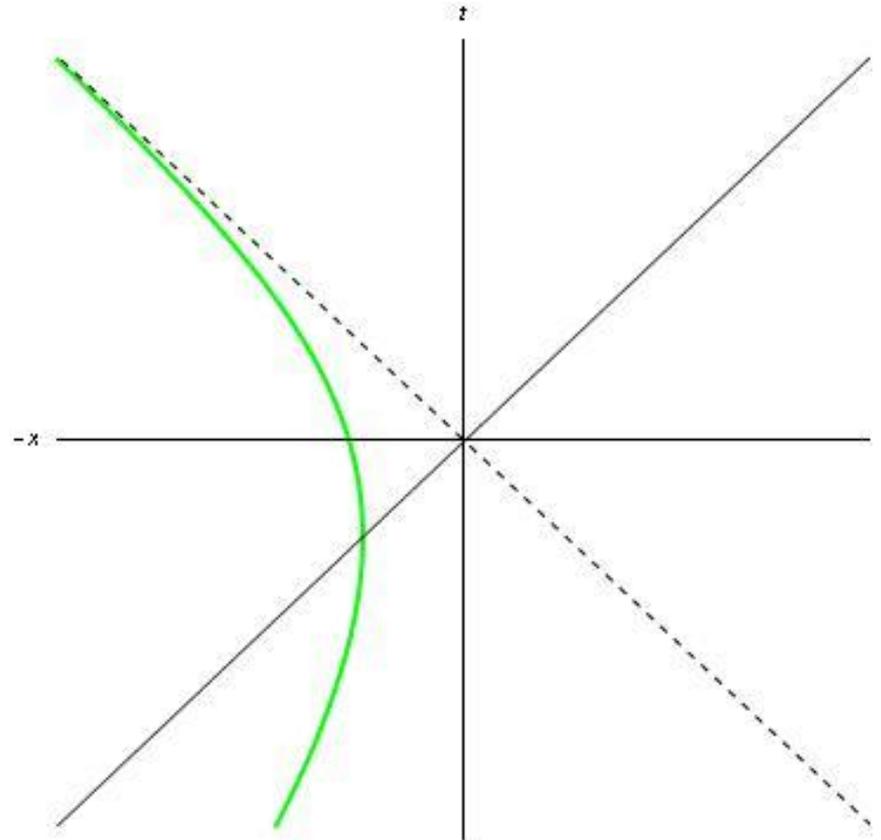
Features:

Future Horizon.

Infinite total energy production.

Infinite total particle production.

Time dependent acceleration.



$$\lim_{t \rightarrow \pm\infty} \alpha(t) = \begin{cases} -\infty \\ 0 \end{cases}$$

$$\lim_{t \rightarrow \pm\infty} v(t) = \begin{cases} -1 \\ 1 \end{cases}$$

$$x(0) = -\frac{\Omega}{\kappa}$$

$$e^{-\Omega} = \Omega$$

$$x(-\frac{1}{2\kappa}) = -\frac{1}{2\kappa}$$

$$p(u) = 2t_u - u \quad u = t - z(t)$$

$$f(v) = 2t_v - v \quad v = t + z(t)$$

Exactly Solvable Trajectory

$$z(t) = -\frac{1}{2\kappa} \sinh^{-1}(e^{\kappa t})$$

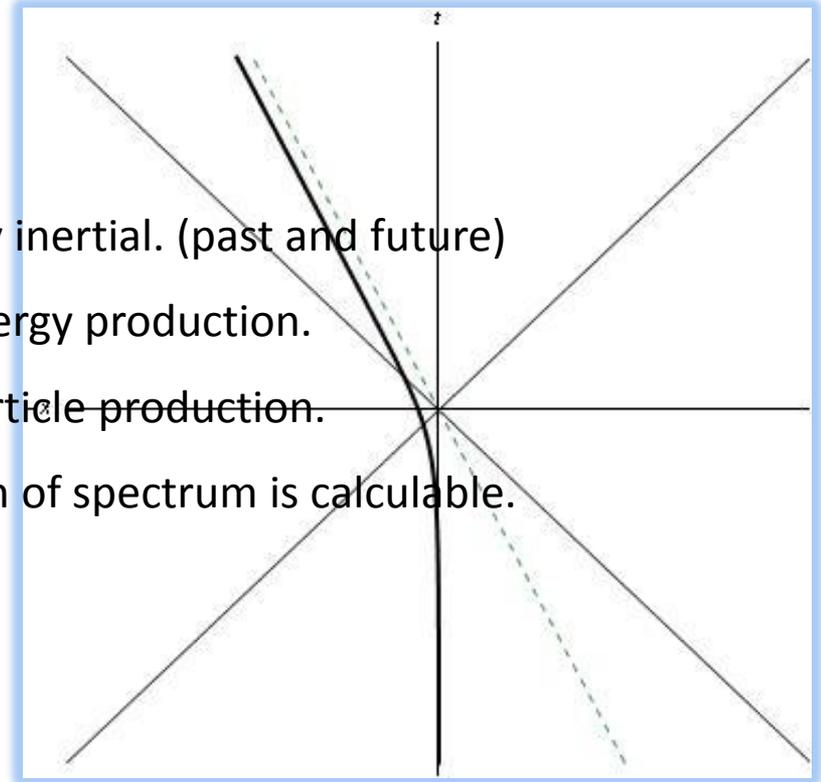
Features:

Asymptotically inertial. (past and future)

Finite total energy production.

Finite total particle production.

Time evolution of spectrum is calculable.



$$E = \int \langle T_{uu} \rangle du = \int \omega \langle N_\omega \rangle d\omega$$

$$|\beta_{\omega'\omega}|^2 = \frac{1}{16\pi\kappa} \frac{\omega\omega'(\omega + \omega')^{-1}}{\left(\frac{\omega'}{4\kappa} + \frac{3\omega}{4\kappa}\right)\left(\frac{\omega}{4\kappa} + \frac{3\omega'}{4\kappa}\right)} \frac{\sinh\left(\frac{\pi\omega}{4\kappa} + \frac{3\pi\omega'}{4\kappa}\right)}{\sinh\left(\frac{\pi\omega}{\kappa} + \frac{\pi\omega'}{\kappa}\right) \sinh\left(\frac{\pi\omega'}{4\kappa} + \frac{3\pi\omega}{4\kappa}\right)}$$