

Bjorken flow in Holographic First Order Phase Transition

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Recent Trends in String Theory and
Related Topics, IPM, April 7, 2018

In collaboration with:
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Outline

Motivation and Introduction

Gravity setup dual to non-CFT

Linearized Gravity (QNMs)

Dynamics of 1st order phase transition

Boost invariant expansion

Summary and Future directions

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- The nature has many examples of strongly coupled systems: QGP, high-T superconductors, big bang
- Traditional tools (perturbation theory, l-QCD) break down
- The AdS/CFT has built a bridge between problems from the Quantum world and methods from General Relativity.

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- Linearized Einstein EoMs for gauge invariant perturbations identifies the dispersion relation $\omega(k) = ?$

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2- **Full time evolution** of a given initial configuration

- Solving the **time dependent** Einstein EoMs
- reading the data corresponds to the **dual FT** from **near boundary** expansion (like 1-pt functions)

Gravity Setup dual to non-CFT

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→ Assuming AdS/CFT works

→ Model the **gravity**+**matter** with a potential $V(\varphi)$

→ We may choose $V(\varphi)$ such that it reproduces the physics of interest

(like **QCD** EoS, 1st or 2nd order **phase transitions**)

Gravity Setup dual to non-CFT

The simplest setup dual to $\mathcal{L}_{\text{CFT}} + \Lambda^{d-\Delta} \langle O_\phi \rangle$

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$$\rightarrow S = \frac{1}{16\pi G_{d+1}} \int d^{d+1}x \sqrt{|g|} \left[R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

Gubser, Nellore '08 & Guryov, Kiritsis, Mazzanti, Nitti '09

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→ The potential has a general form:

$$V(\phi) = -d(d-1) \cosh(\gamma\phi) + b_2 \phi^2 + b_4 \phi^4$$

$$\sim -d(d-1) + \frac{m^2}{2} \phi^2 + \mathcal{O}(\phi^4), \quad m^2 = \Delta(\Delta - d)$$

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→ We are interested in FT's in flat background with fixed sources (Λ at different temperatures (corresponding to φ_H))

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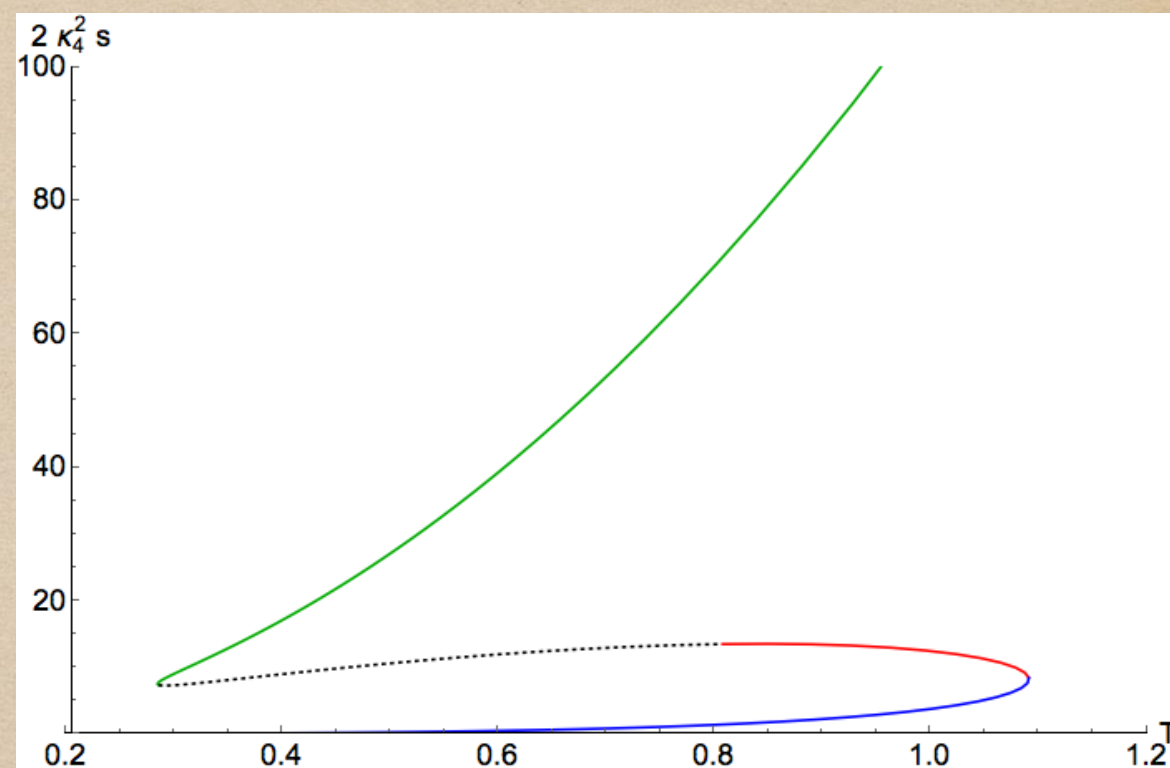
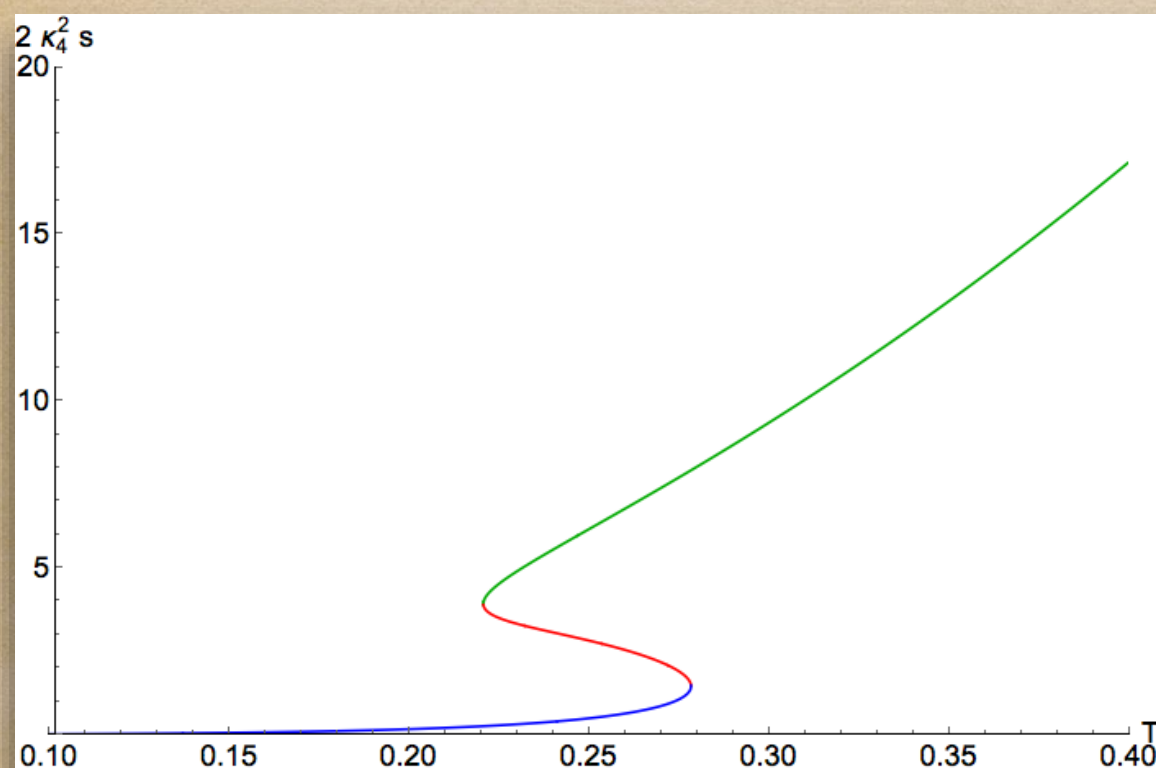
→ Equation of State: $s(T)$, $T_{ij}(T)$, $C_s^2 = \frac{s}{T} (ds/dT)^{-1}$

Gravity Setup dual to non-CFT

Examples: 3+1-gravity/CFT $_{\mp}(\Delta=2)$

$$V_1 = -6 \cosh(\phi/\sqrt{3}) - 0.2 \phi^4,$$

$$V_2 = -6 \cosh(\phi/\sqrt{3}) - 0.3 \phi^4$$



Stable regimes:

Green: Large BH

Blue: Small BH

Unstable regimes:

Red: thermodynamic instability

Black-dotted: Dynamical instability

Linearized Gravity
Quasi Normal Mods (QNM)

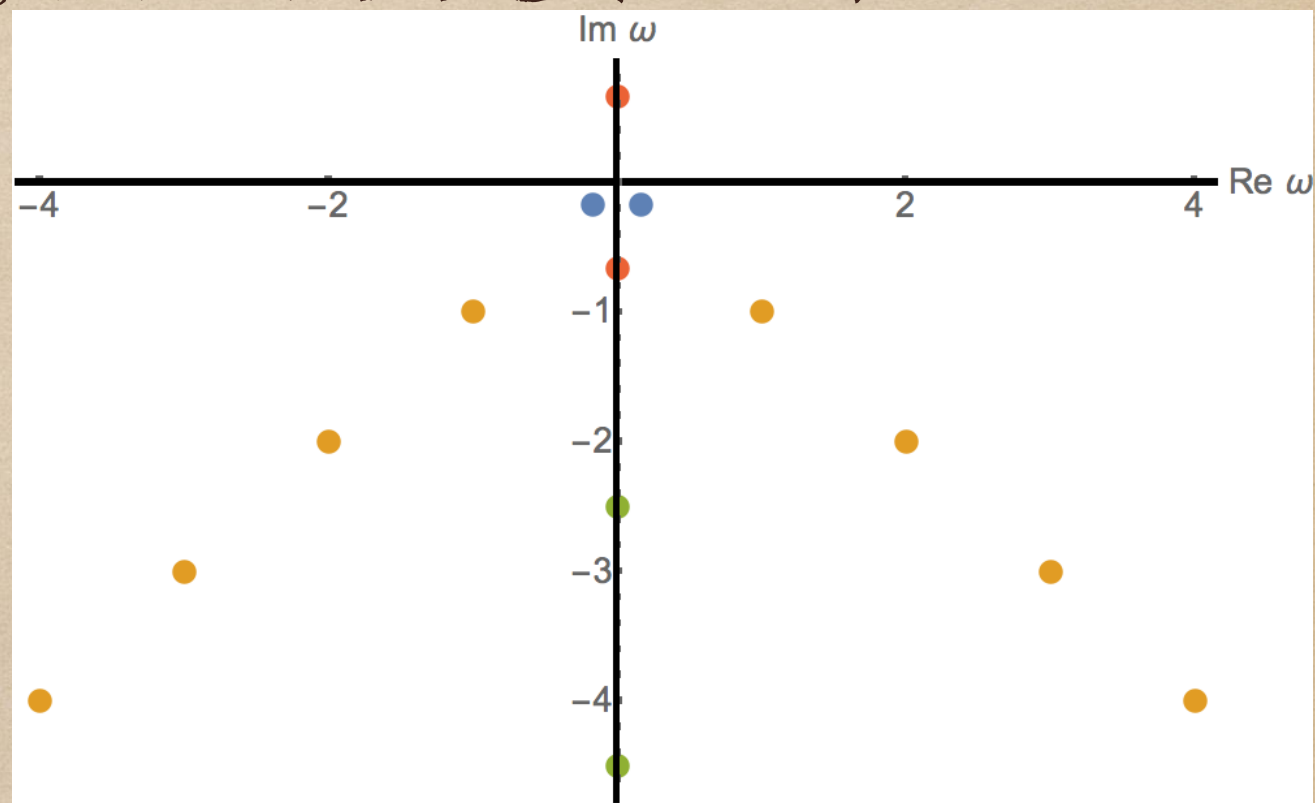
Linearized Gravity (QNMs)

QNM equations: linearized Einstein equations ($\delta\Phi_i \sim e^{-i\omega t + ikx} \Psi_i(r)$)
$$\mathbb{E} \circ M(g_{\mu\nu} + \epsilon h_{\mu\nu}, \varphi + \epsilon \psi) = \mathbb{E} \circ M(g_{\mu\nu}, \varphi) + \epsilon \text{QNM}_{\text{eqs}}$$

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 $E_{OM}(g_{\mu\nu} + \epsilon h_{\mu\nu}, \varphi + \epsilon \psi) = E_{OM}(g_{\mu\nu}, \varphi) + \epsilon QNM_{eqs}$

Typical structure of the QNM's in AA_{dS} ($k = \text{finite}$)



Red: Unstable mode

Green: Diffusive modes

Yellow: Non-hydrodynamic modes

Blue: Hydrodynamic modes, $\omega_{(k \rightarrow 0)} \rightarrow 0$

Linearized Gravity (QNMs)

New features in non-CFT cases:

→ **Sign for bubble formation** in 1st order phase transitions!

Janík, Jankowski, HS, PRL '16

→ **Dynamical instability**: A thermodynamically stable black hole **MAY** suffers an unstable non-hydro mode!

Janík, Jankowski, HS, PRL '16
Gursoy, Jansen, Wilke, PRD '16

→ **Crossing** between hydro and nonhydro modes!

(limits applicability of hydrodynamics)

Janík, Jankowski, HS, JHEP '16, ...

→ **Joining modes**: Hydro modes and nonhydro modes might be indistinguishable!

Janík, Jankowski, HS, JHEP '16

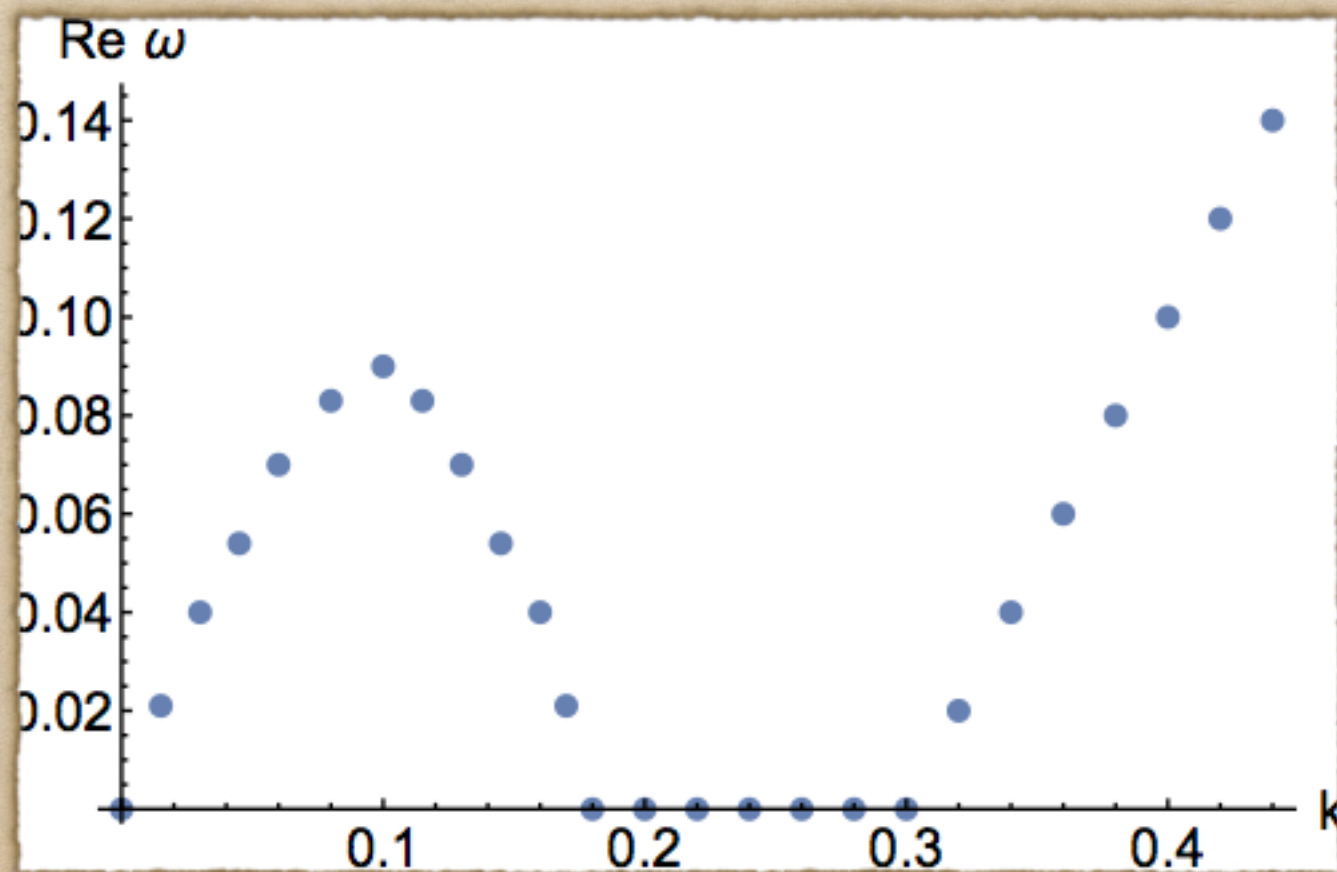
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($\text{Re}(\omega)=0$ for $k_{\min} < k < k_{\max}$)

Janik, Jankowski, HS, PRL '16

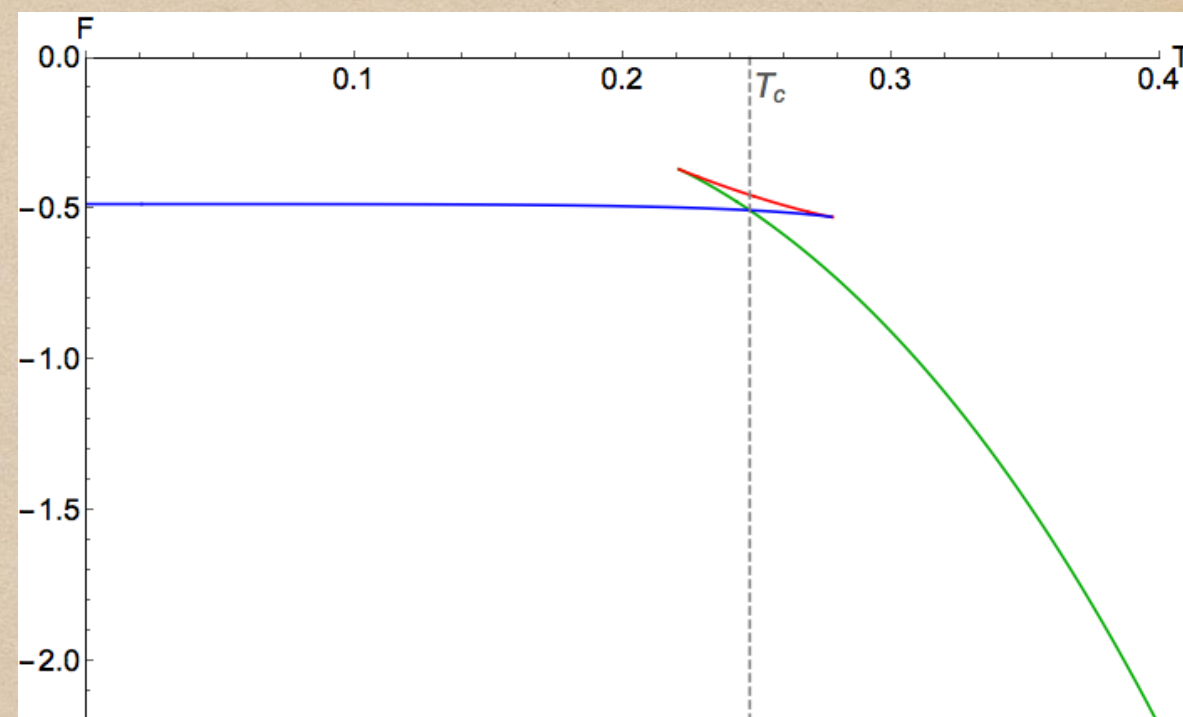
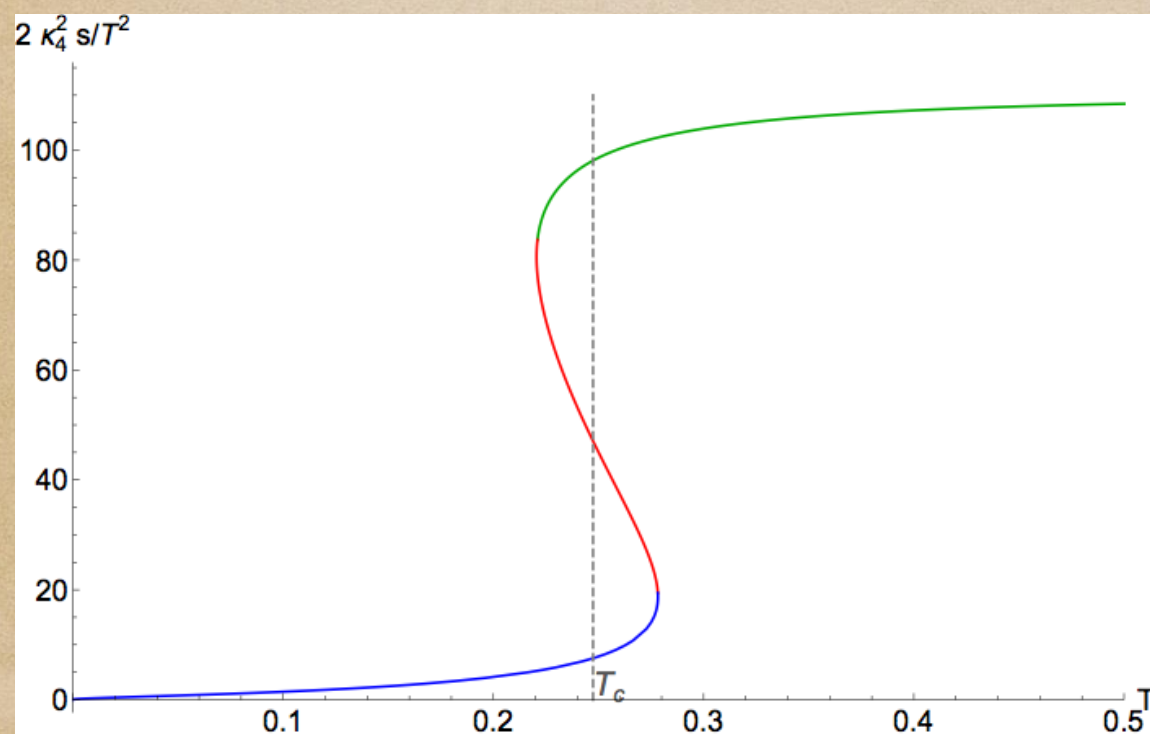


Dynamics of 1st order
phase transition

Dynamics of 1st order phase transition

Potential in 4D: $V(\phi) = -6 \cosh(\phi/\sqrt{3}) - 0.2 \phi^4$, $\Delta = 2$

Janik, Jankowski, HS, PRL '17



Initial configurations: perturbed black holes
in Spinodal region

Dynamics of 1st order phase transition

Some technical points:

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Metric Ansatz (in Eddington-Finkelstein): $x \sim x + i 2\pi$

$$ds^2 = -A(u, x, t) dt^2 - \frac{2dtdu}{u^2} - 2B(u, x, t) dt dx + \Sigma(u, x, t)^2 [H(u, x, t) dx^2 + H(u, x, t)^{-1} dy^2]$$

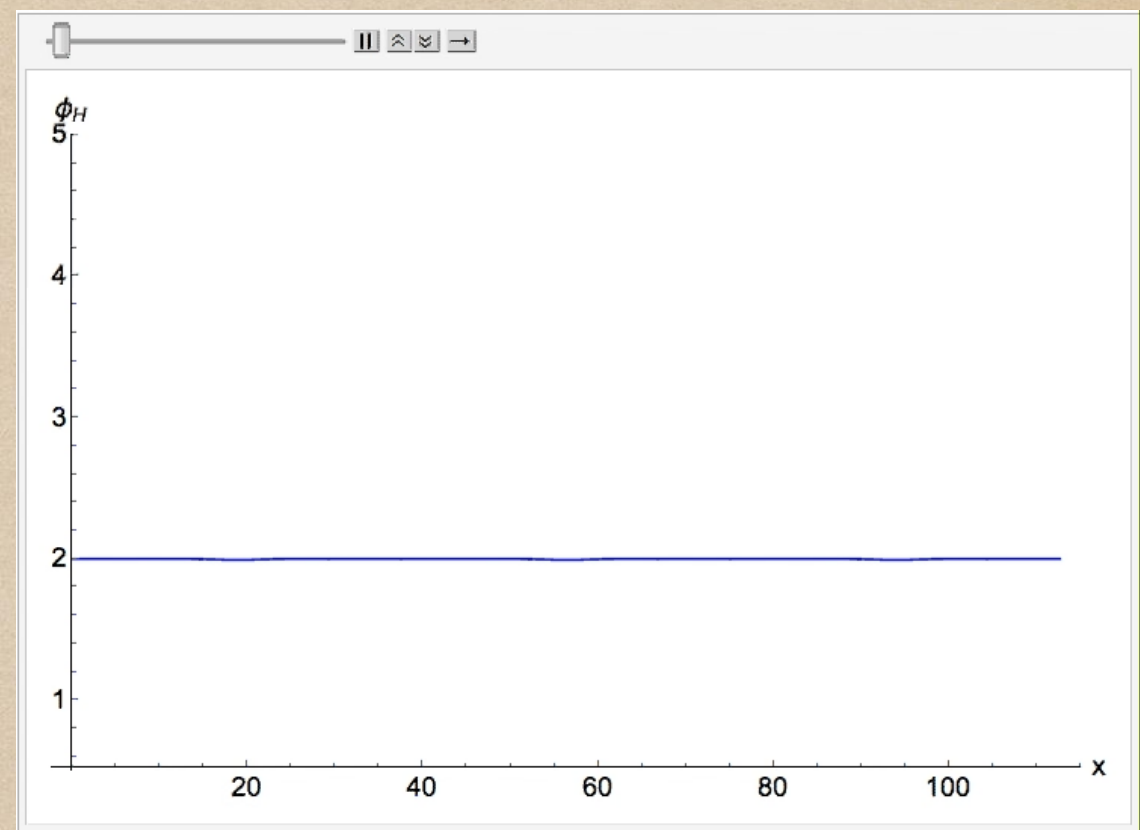
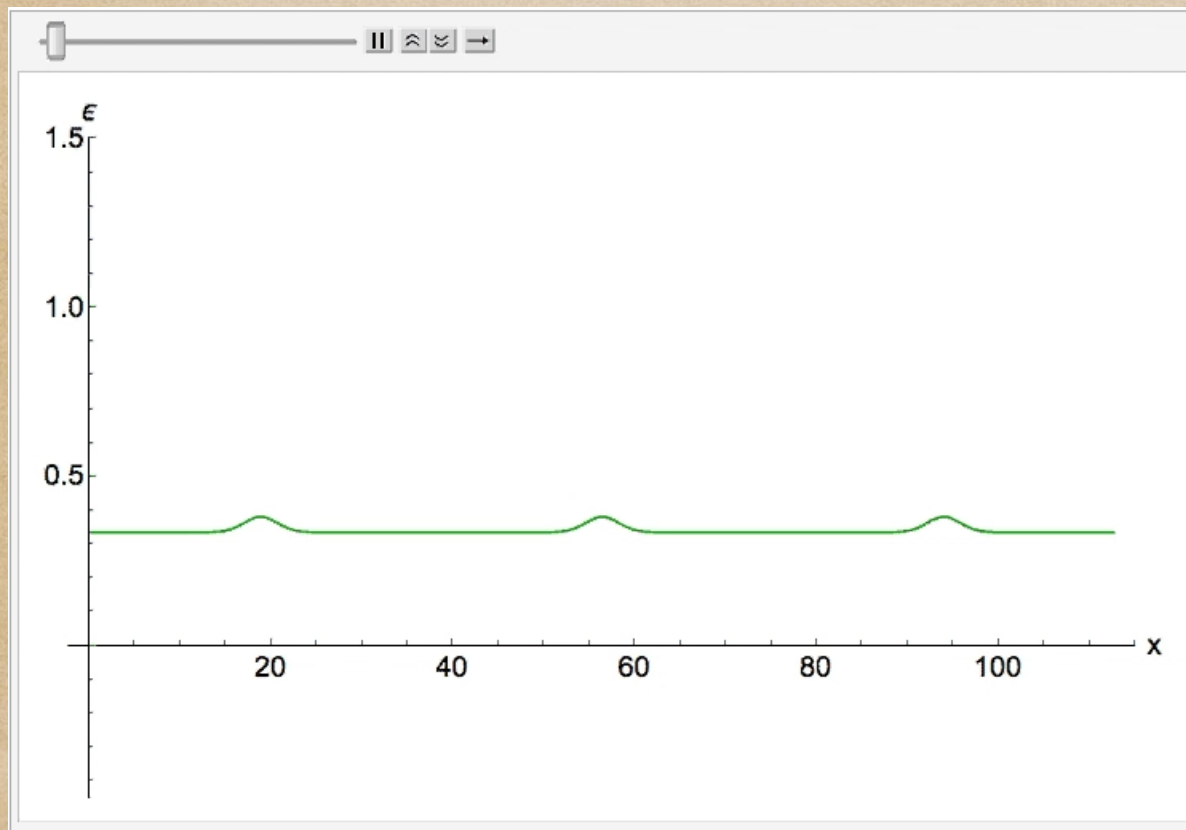
Comments:

- Time integration: Runge-Kutta and Adams-Bashforth methods
- Proper boundary conditions at AdS and at the apparent horizon
- Holographic renormalization: $\langle T_{ij} \rangle$, $\langle O_\phi \rangle$
- Ward identities: $\langle T^i_i \rangle = \langle O_\phi \rangle$, $\nabla^i \langle T_{ij} \rangle = 0$

Dynamics of 1st order phase transition

Results: initial config.:

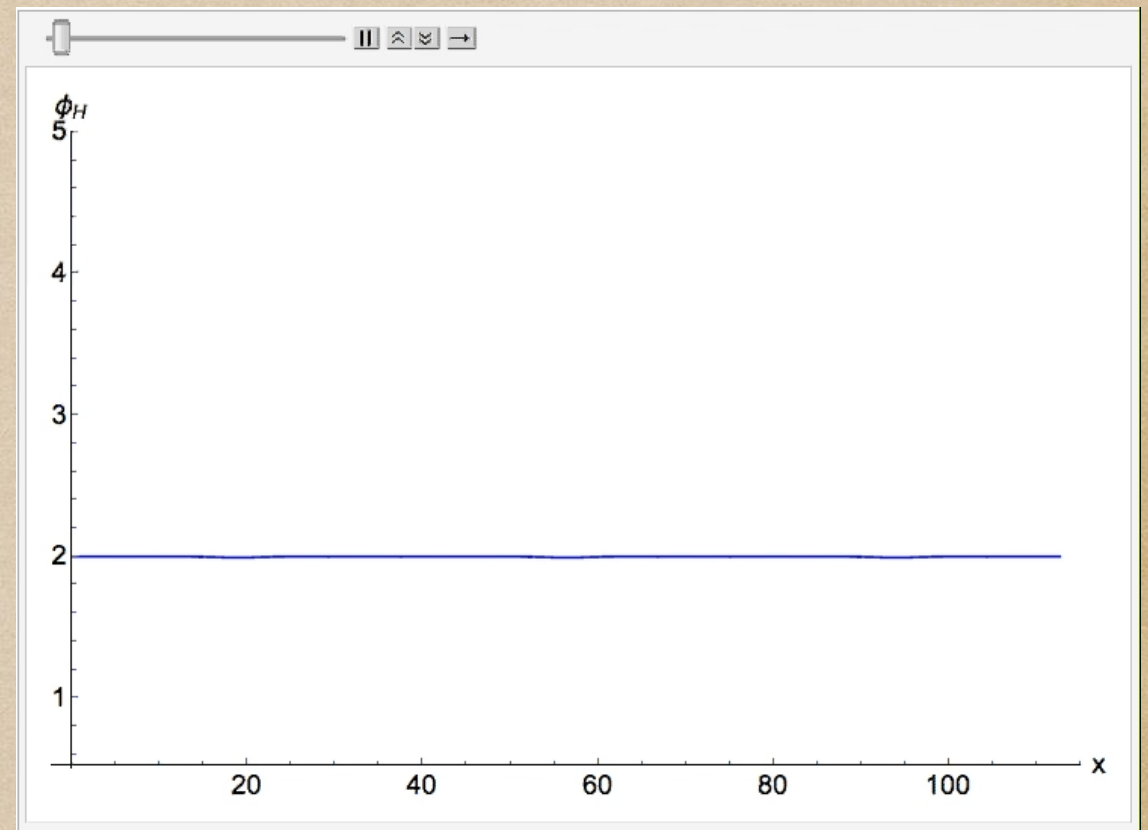
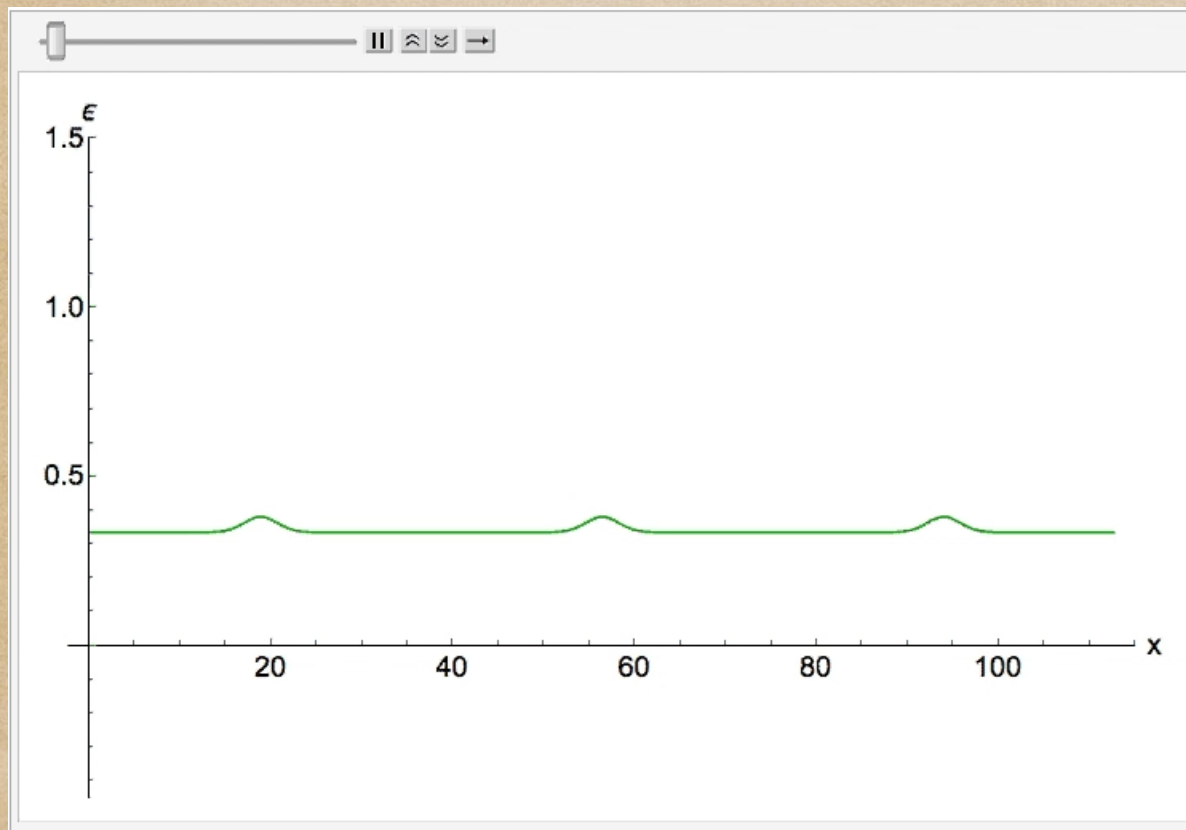
$$\varphi_H = 2 \quad \delta\Sigma \propto \exp\left(-\frac{\cos(kx)^2}{w}\right)$$



Dynamics of 1st order phase transition

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The final state: Static Black holes with inhomogenous horizons
With homogenous sources

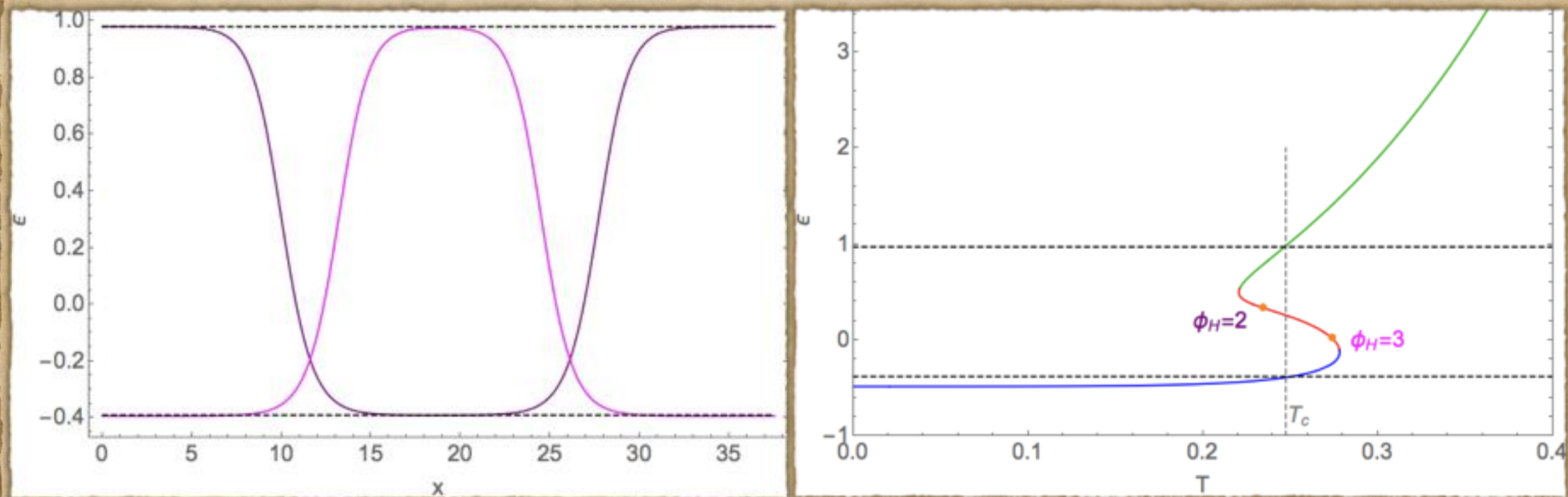
Dynamics of 1st order phase transition

Results: New static black hole: inhomogeneous horizon
and temperature T_c

Días, Santos, Way, JHEP '17

Dynamics of 1st order phase transition

Results: New static black hole: inhomogeneous horizon and temperature T_c



Green lines: Min and Max points on EoS at T_c

Red and Blue line: $\epsilon(x)$ for the final state of perturbations

Hydrodynamics is applicable with 2nd-order transport Coefficients

Boost invariant expansion

Expanding plasma: Boost-invariant expansion as a simple model

The physics of $(d+1)$ -dim FT
depends only on τ and
not rapidity y . $i=1, \dots, d-1$

Bjorken, PRD, '83

$$ds^2 = -dt^2 + d\tilde{y}^2 + dx_i^2$$

$$t = \tau \cosh y, \quad \tilde{y} = \tau \sinh y$$

$$ds^2 = -d\tau^2 + \tau^2 dy^2 + dx_i^2$$

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Dual gravity Ansatz in $3+1$ d:

$$ds^2 = -A(u, \tau) d\tau^2 - \frac{2d\tau du}{u^2} + \Sigma(u, \tau)^2 [H(u, \tau) dx^2 + \tau^2 H(u, \tau)^{-1} dy^2]$$

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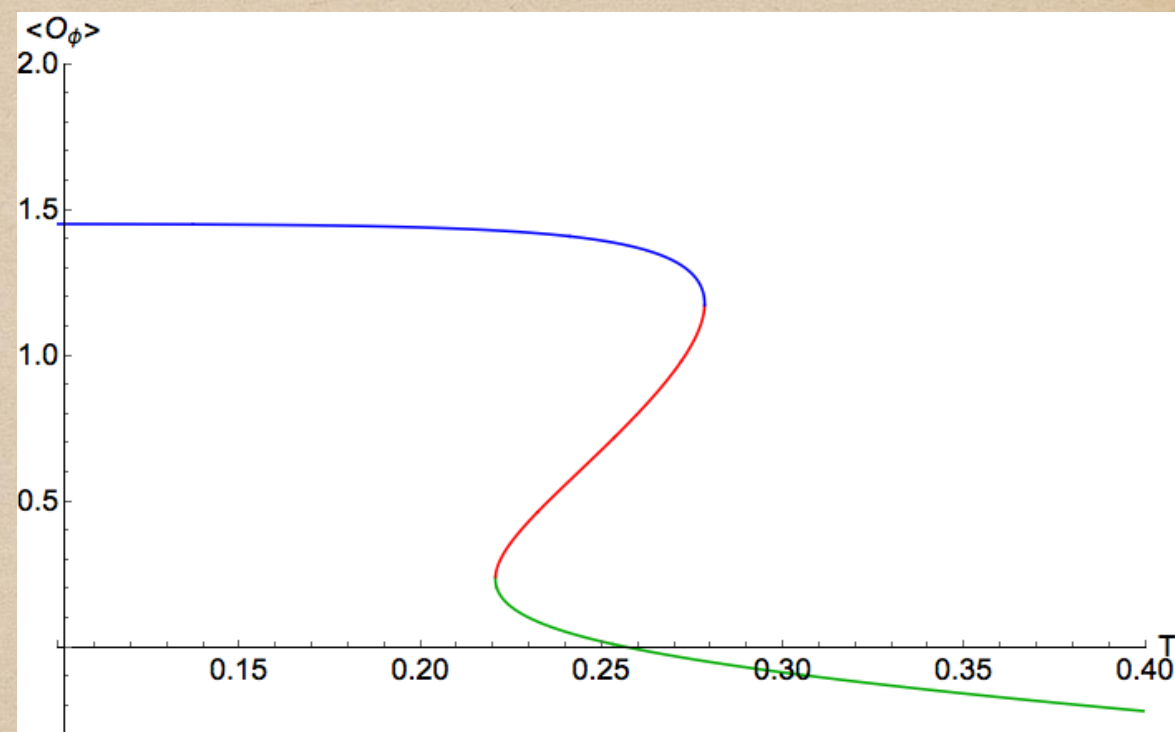
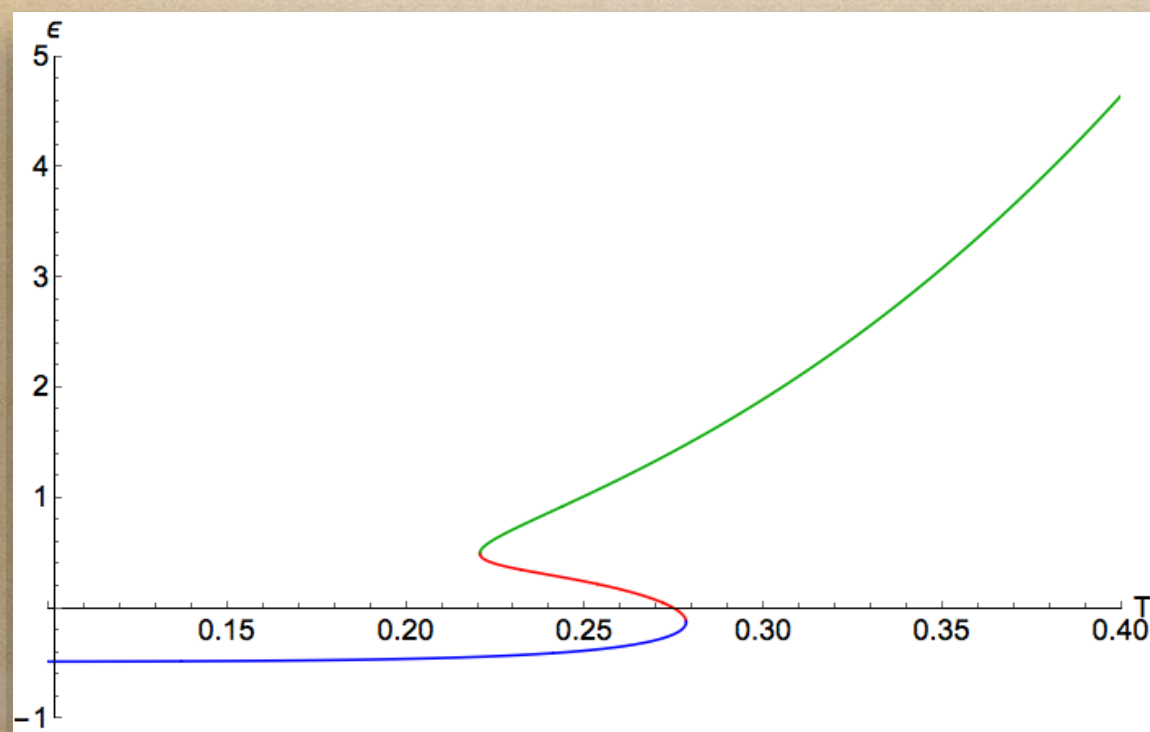
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Janik, Jankowski, HS, "soon"

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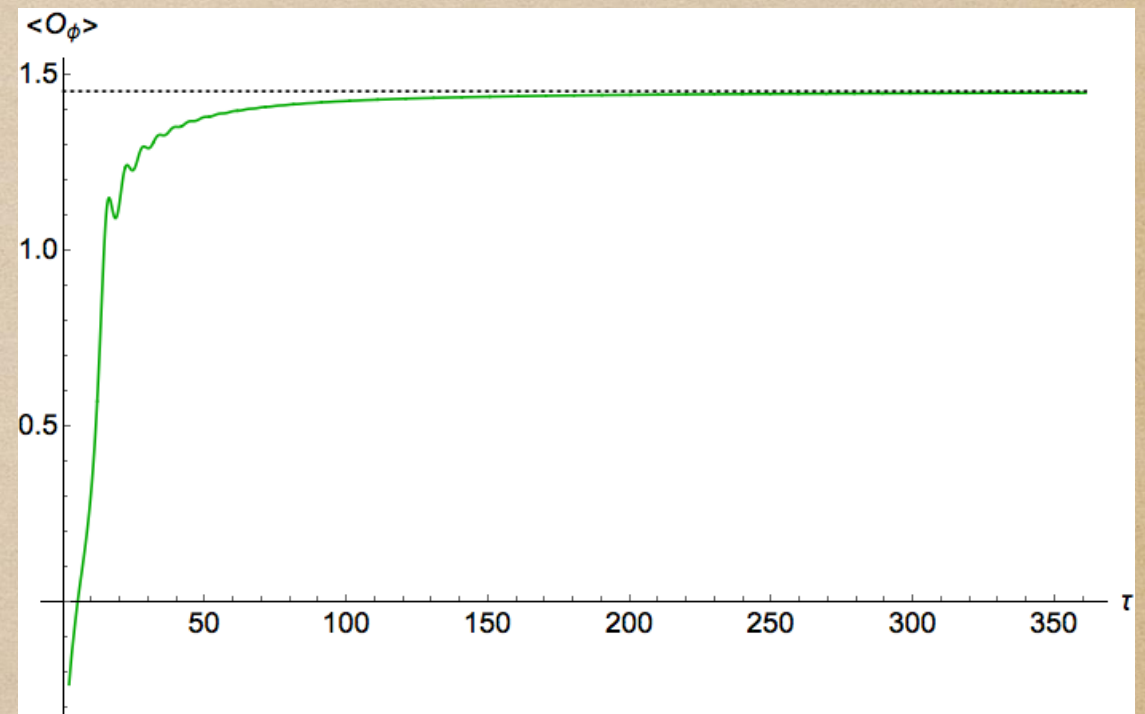
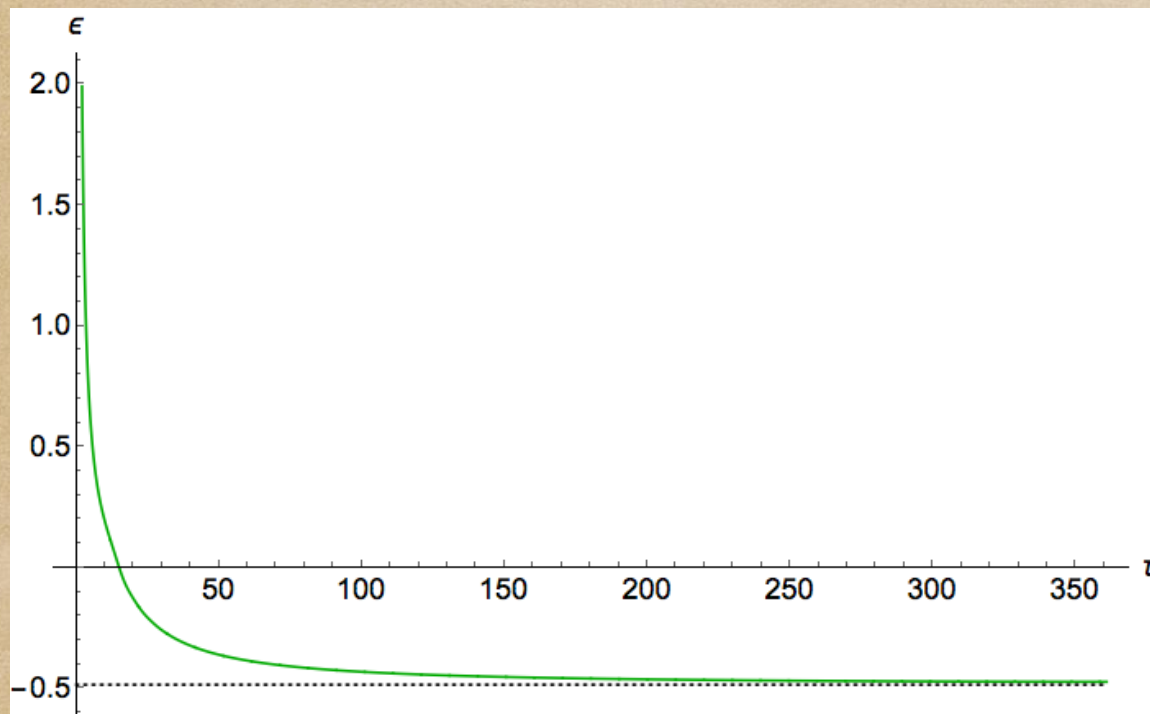


Initial configuration has "enough" energy

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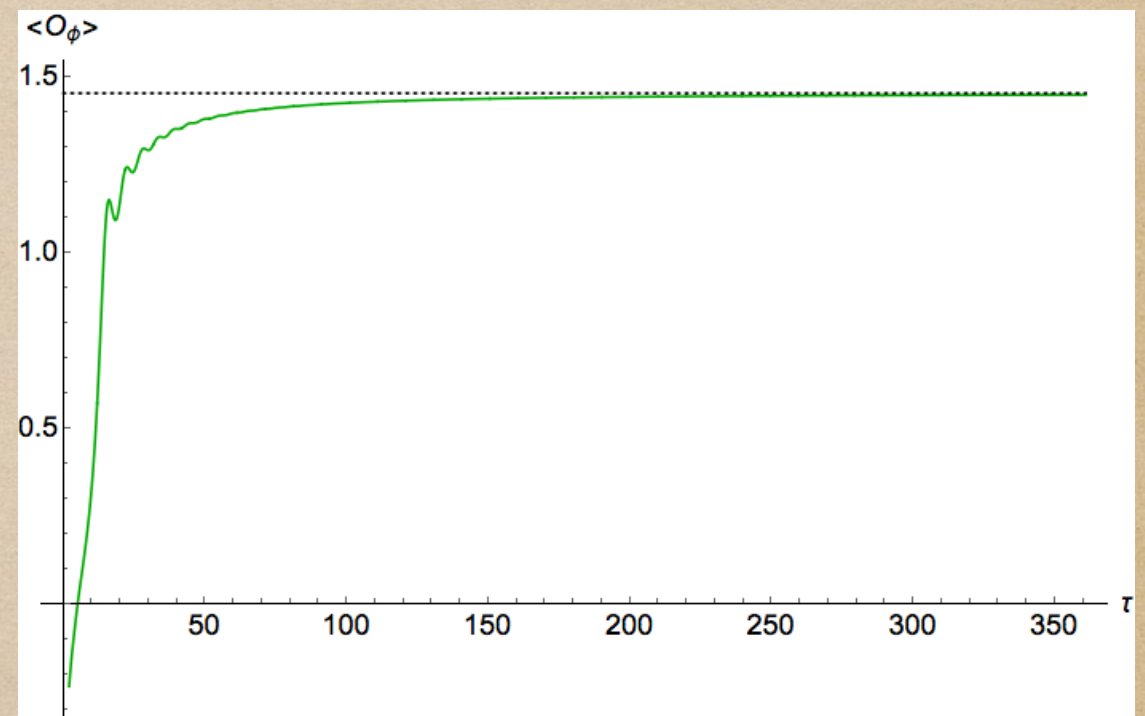
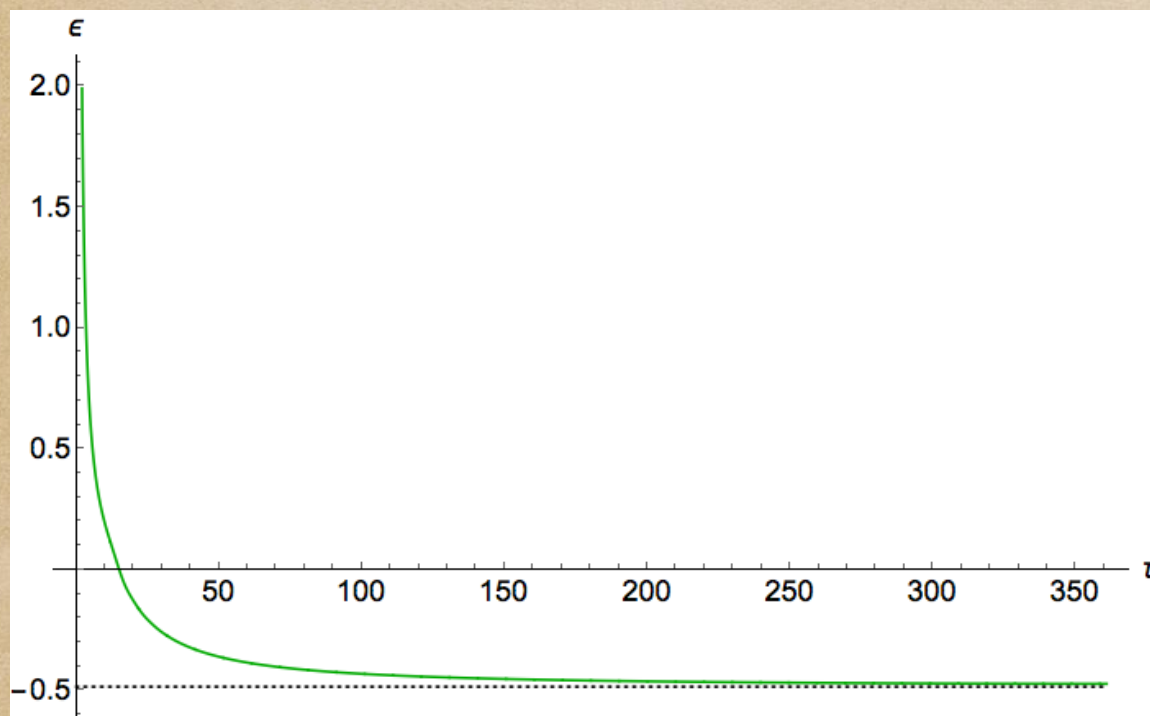
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→ The energy is decreasing

→ The expansion passes the Spinodal region "trivially"

Expanding plasma: Boost-invariant expansion as a simple model

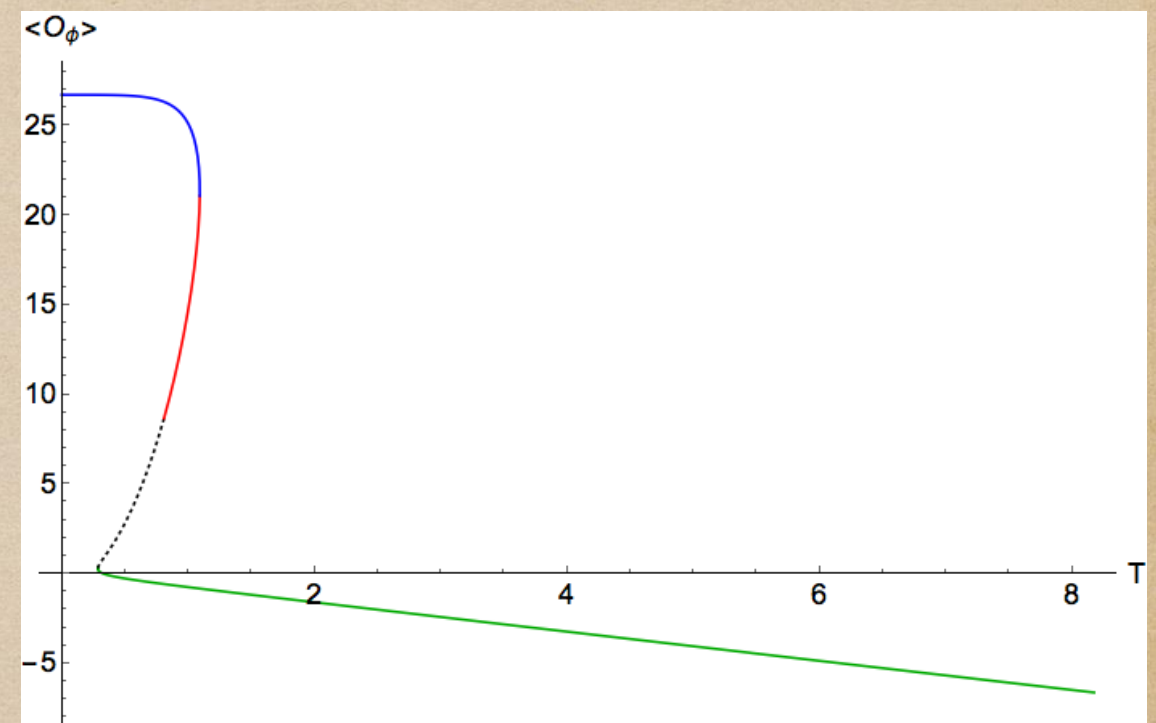
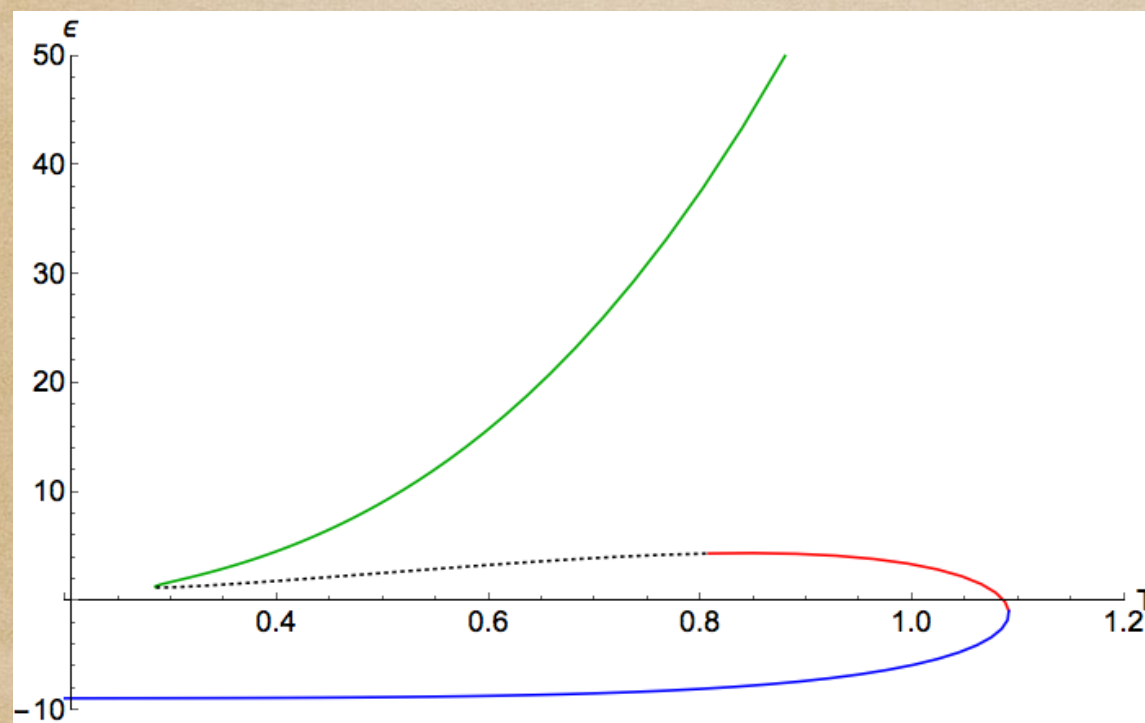
2nd Potential in 4D: $V_2(\phi) = -6 \cosh(\phi/\sqrt{3}) - 0.3 \phi^4, \quad \Delta = 2$

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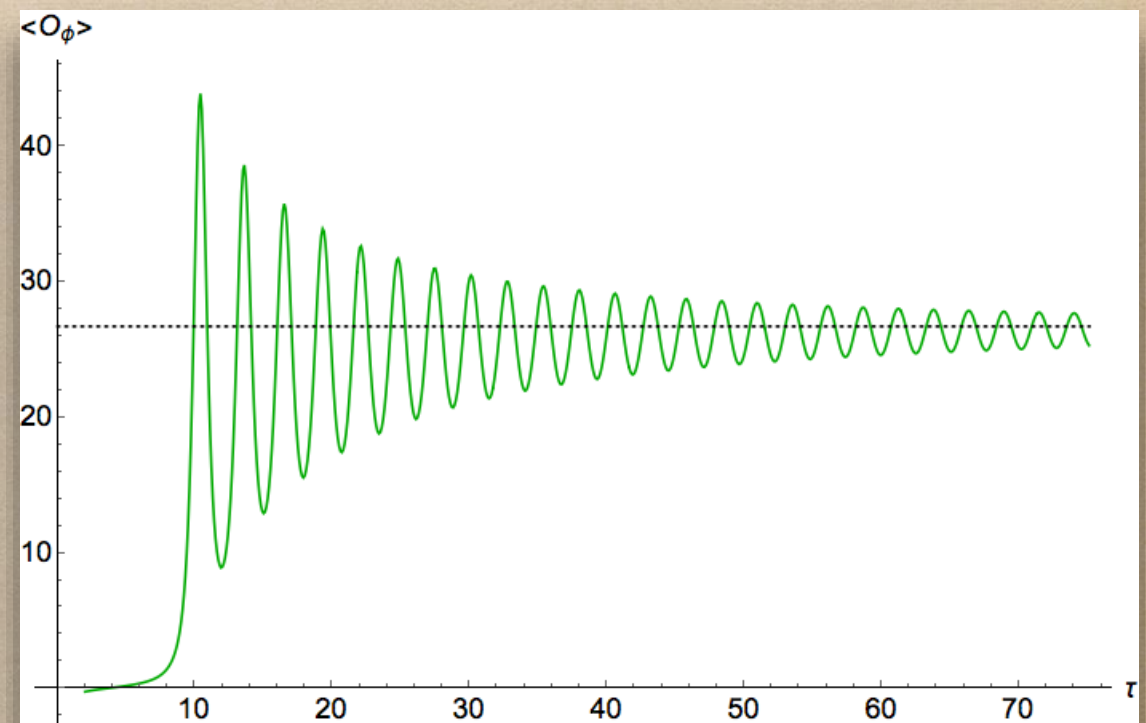
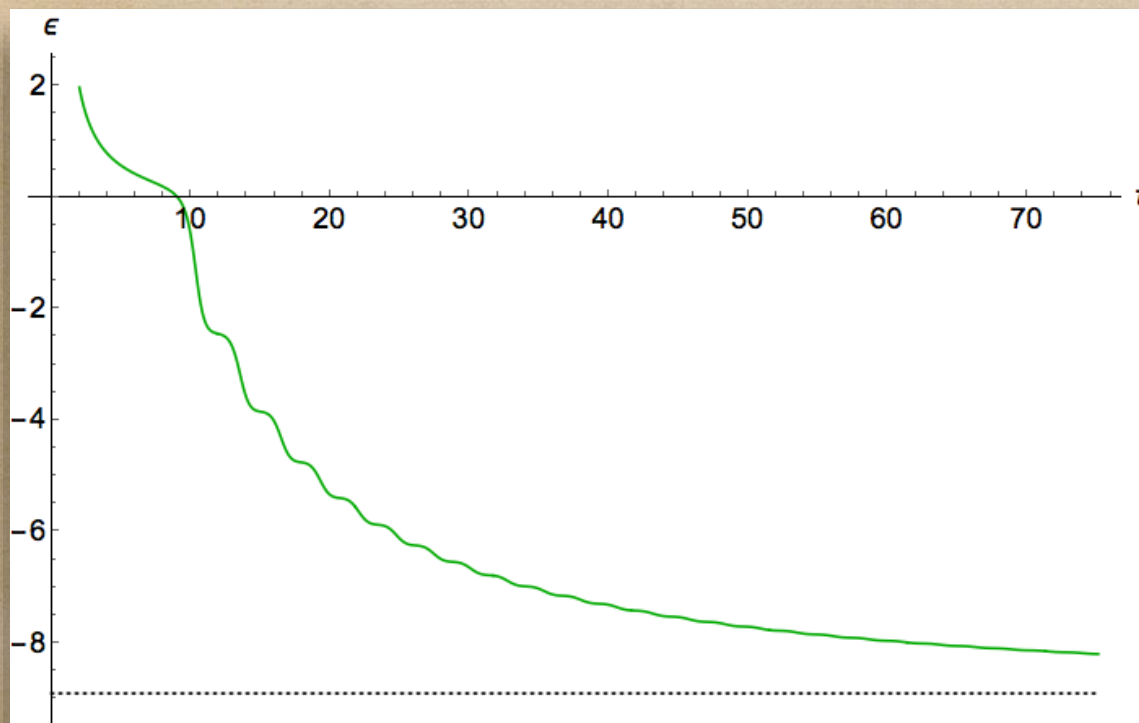


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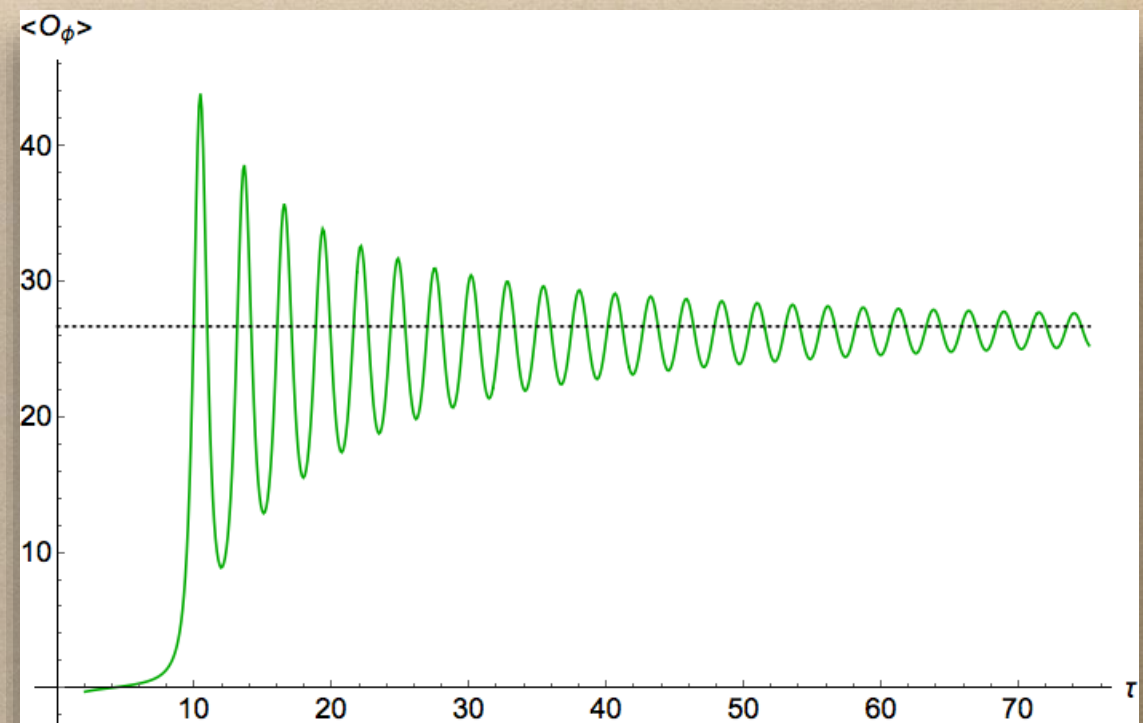
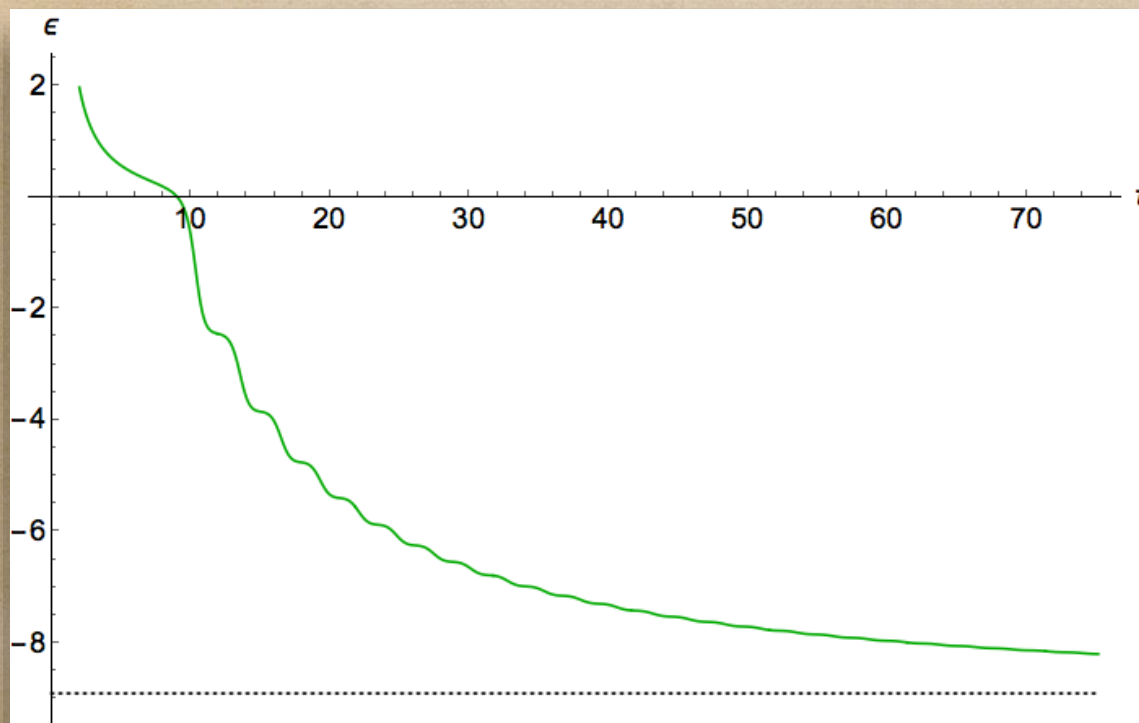
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- The energy is decreasing
- There is a reminiscent of dynamical instability until the late time in expansion!!!

Summary

- Higher QNMs introduce new instabilities in BHs
- Bubble formation is studied (in strongly coupled regime, with universal properties)
- New black holes with homogenous sources but inhomogenous horizon
- Plasma expansion (boost invariant flow) is studied and needs more,...

Summary and future directions

- Higher QNMs introduce new instabilities in BHs
- Bubble formation is studied (in strongly coupled regime, with universal properties)
- New black holes with homogenous sources but inhomogenous horizon
- Plasma expansion (boost invariant flow) is studied and needs more,...
- Different dimensions (in particular 3D gravity + matter)
- Investigating anisotropic plasma
- Collision of bubbles in 1st order phase transition
- Using probes: strings, Wilson line/loop, EE, ...
- New black holes: Domain wall solutions to GR

Thanks for your attention

Dynamics of 1st order phase transition

Some technical points:

Metric Ansatz (in Eddington-Finkelstein): $x \sim x + i 2\pi$

$$ds^2 = -A(u, x, t) dt^2 - \frac{2dtdu}{u^2} - 2B(u, x, t) dt dx + \Sigma(u, x, t)^2 [H(u, x, t) dx^2 + H(u, x, t)^{-1} dy^2]$$

EoMs: Using the new time derivative:

$$d_+ := \partial_t - \frac{u^2 A}{2} \partial_u$$

$$\partial_u \phi = S_\phi(\Sigma, H), \quad \partial_{u^2} B + p_1 \partial_u B + p_2 B = S_B(\Sigma, H, \phi)$$

$$\partial_u (d_+ \Sigma) + p_3 d_+ \Sigma = S_{d_+ \Sigma}(\Sigma, H, \phi, B)$$

$$\partial_u (d_+ H) + p_4 d_+ H = S_{d_+ H}(\Sigma, H, \phi, B, d_+ \Sigma)$$

$$\partial_u (d_+ \phi) + p_5 d_+ \phi = S_{d_+ \phi}(\Sigma, H, \phi, B, d_+ \Sigma, d_+ H)$$

$$\partial_{u^2} A + p_6 \partial_u A = S(\Sigma, H, \phi, B, d_+ \Sigma, d_+ H, d_+ \phi)$$

Solve the equations at $t=0$ and then one step forward