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Non-Gaussianity

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Outline

- Inflation
- Observational cosmology
- Theoretical cosmology
- Non-Gaussianity and its charcteristics



Standard Model of Cosmology



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Problems of standard model

- Flatness problem
- Horizon problem
- Monopole problem

Solving the problems: Inflation

Accelerating universe can solve the problems.

A.Guth 1980

- Solution to:
 - Flatness problem
 - Horizon problem
 - Monopoles problem

Dynamics of Inflation



- Matter content of universe at inflationary phase is inflaton, a scalar field
- Accelerating is possible:

$$\rho_{\phi} = \frac{1}{2}\dot{\phi}^{2} + V(\phi) \quad , \quad P_{\phi} = \frac{1}{2}\dot{\phi}^{2} - V(\phi)$$

conditions for inflation:

Slow-Roll conditions:

$$\epsilon = \frac{m_{\rm pl}^2}{16\pi} \left(\frac{V_{\phi}}{V}\right)^2 , \ \eta = \frac{m_{\rm pl}^2 V_{\phi\phi}}{8\pi V}$$

 $\epsilon \ll 1$ and $|\eta| \ll 1$

• In order to inflation occurs and lasts sufficiently long time.



B.~A.~Bassett, S.~Tsujikawa and D.~Wands, %``Inflation dynamics and reheating," Rev.\ Mod.\ Phys.\ \ {\bf 78}, 537 (2006) [astro-ph/0507632].

CMB: The observations

CMB: Radio waves remained from early universe. We observe them when they reach to our horizon.





What we learn from CMB data?



• Power spectrum:

The relation between CMB anisotropy and quantum fluctuations of early universe

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What we learn from CMB data?

Non-Gaussianity:

Measures the correlation between different scales or different parts of the universe.

 Gaussian perturbations have zero non-Gaussianity

How can we build fluctuations on CMB?

• Quantum fluctuations exist naturally

• Can they be the initial seeds of CMB fluctuations?

Inflation can do that

How can we model the CMB? Perturbation theory in cosmology

Perturbation of metric around FRW background:

$$ds^{2} = -(1+2A)dt^{2} + 2a(\partial_{i}B - S_{i})dx^{i}dt +a^{2}\left[(1-2\psi)\delta_{ij} + 2\partial_{ij}E + 2\partial_{(i}F_{j)} + h_{ij}\right]dx^{i}dx^{j}$$

+ Perturbation of matter:

 $\phi = \phi_0(t) + \delta\phi(t, x)$

A good quantity!

Curvature perturbation:
$$\mathcal{R} \equiv \psi + \frac{H}{\dot{\phi}} \delta \phi$$

Power spectrum/non-Gaussianity and CMB

Power spectrum of curvature perturbation:

$$\langle \mathcal{R}_{k_1}\mathcal{R}_{k_2} \rangle \longrightarrow \mathcal{P}_{\mathcal{R}} \equiv \frac{4\pi k^3}{(2\pi)^3} |\mathcal{R}^2|$$

Non-Gaussianity:

$$<\mathcal{R}_{k_1}\mathcal{R}_{k_2}\mathcal{R}_{k_3}>=-6/5f_{NL}(<\mathcal{R}_{k_1}\mathcal{R}_{k_2}>+2perm.)$$

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To gather...

- Inflation
- Observation of temperature fluctuations
- Cosmological perturbation theory and inflation
- Power spectrum and non-Gaussianity





Why non-Gaussianity is important?

- Includes information about statistics of fluctuations
- Can constrain or rule out models of inflation
- Is both detectable and calculable (in principle)

Characteristics of Non-Gaussianity



• These characteristics are detectable distinguishable

Shape of Non-Gaussianity

Shape is the scale dependence of non-Gaussianity



http://www.umich.edu/~mctp/SciPrgPgs/eve nts/2011/CosmoNongaussianity/index.html (by Matarrese)

Size of Non-Gaussianity

| http://www.umich.edu/ | -mctp/SciPrgP | | |
|-------------------------------------|-----------------|-------------------------------|-----------------------|
| gs/events/2011/CosmoNo ndex.html | ongaussianity/i | WMAP 7-yrs | (Komatsu et al. 2010) |
| | Local | - 10 < f _{NL} < 74 | |
| | Equilateral | - 214 < f _{NL} < 266 | (95% c.l) |
| | Orthogonal | -410 < f _{NL} <6 | |
| | | | |

 $\langle \mathcal{R}_{k_1}\mathcal{R}_{k_2}\mathcal{R}_{k_3} \rangle = -6/5f_{NL}(\langle \mathcal{R}_{k_1}\mathcal{R}_{k_2} \rangle + 2perm.)$

Planck will have more to say about non-Gaussianity

Is the non-Gaussianity large?

Temperature $(f_{NL} = 10^4)$



http://ipht.cea.fr/Pisp/pontavignon2008/pr ogramme.htm (by Wandelt & Yadav)

How to calcualte the non-Gaussianity?

- Maldacena's Method
- δN formalism

Maldacena's method

Quadratic action(The propagator):

$$S_2 = \int \mathrm{d}\tau \,\mathrm{d}^3 x \,a^2 \left[\frac{\Sigma}{H^2} (\mathcal{R}')^2 - \varepsilon (\partial \mathcal{R})^2\right]$$

%``Primordial non-Gaussianities in single field inflation," JCAP\ {\bf 0506}, 003 (2005) [astro-ph/0503692].

Maldacena's method

• Cubic action(The interaction):

$$S_{3} = \int \mathrm{d}t \,\mathrm{d}^{3}x \,\left[-\frac{2}{3}a^{3}\varepsilon \left(u + \frac{\varepsilon}{\varepsilon_{X}}\frac{s}{3} \right) \frac{\dot{\mathcal{R}}^{3}}{H} + 4a^{5}\varepsilon (2u+\varepsilon)H\dot{\mathcal{R}}^{2}\partial^{-2}\dot{\mathcal{R}} - 4a^{3}\varepsilon\dot{\mathcal{R}}\partial^{2}\mathcal{R}\partial^{-2}\dot{\mathcal{R}} \right],$$

• coupling constants!

%``Primordial non-Gaussianities in single field inflation," JCAP\ {\bf 0506}, 003 (2005) [astro-ph/0503692].

δN Formalism



D.~Wands, K.~A.~Malik, D.~H.~Lyth and A.~R.~Liddle, %``A New approach to the evolution of cosmological perturbations on large scales,"

Phys.\ Rev.\ D\ {\bf 62}, 043527 (2000) [astro-ph/0003278].

Some Examples

• Typical slow-roll single field inflation:

The size of non-Gaussianity is of order of slow-roll parameters and **non-detectable**.

• Multiple field inflation:

Entropic perturbations can induce large non-Gaussianity

Some Examples

• Dirac-Born-Infield (DBI) inflation: M.~Alishahiha, E.~Silverstein and D.~Tong, %``DBI in the sky," Non-standard kinetic term

Phys.\ Rev.\ D\ {\bf 70}, 123505 (2004) [hep-th/0404084].

$$\mathcal{L}_{\text{eff}} = -\frac{1}{g_s} \sqrt{-g} \left(f(\phi)^{-1} \sqrt{1 + f(\phi)g^{\mu\nu}\partial_{\mu}\phi\partial_{\nu}\phi} + V(\phi) \right)$$

$$\frac{d^2 v_k}{d\tau^2} + \left(c_s^2 k^2 - \frac{\nu^2 - \frac{1}{4}}{\tau^2}\right) v_k = 0$$
 sponsible for large non-

Gaussianity

$$\begin{split} S_3 &= \frac{1}{2} \int \mathrm{d}t \, \mathrm{d}^3x \, a^3 \bigg[-\frac{4}{3} \varepsilon \left(\frac{1}{3c_{\mathrm{s}}^2} \frac{\varepsilon}{\varepsilon_X} s + u \right) \frac{\dot{\mathcal{R}}^3}{Hc_{\mathrm{s}}^2} + \frac{2\varepsilon}{c_{\mathrm{s}}^2} \left(3u + \frac{\varepsilon}{c_{\mathrm{s}}^2} \right) \mathcal{R} \dot{\mathcal{R}}^2 \\ &+ \frac{2\varepsilon}{a^2 c_{\mathrm{s}}^2} (\varepsilon - 2s - u c_{\mathrm{s}}^2) \mathcal{R} (\partial \mathcal{R})^2 \\ &- \frac{4}{a^2} \frac{\varepsilon}{c_{\mathrm{s}}^2} \dot{\mathcal{R}} \mathcal{R}_{,i} \chi_{,i} + \mathcal{R}^2 \dot{\mathcal{R}} \frac{\varepsilon}{c_{\mathrm{s}}^2} \frac{\mathrm{d}}{\mathrm{d}t} \left(\frac{\eta}{c_{\mathrm{s}}^2} \right) - \frac{\varepsilon^3}{c_{\mathrm{s}}^4} \mathcal{R} \dot{\mathcal{R}}^2 + \frac{1}{a^2} \varepsilon \mathcal{R} \chi_{,ij} \chi_{,ij} \bigg] \end{split}$$

%``Primordial non-Gaussianities in single field inflation," JCAP\ {\bf 0506}, 003 (2005) [astro-ph/0503692].

Some Examples

• Ghost inflation:

N.~Arkani-Hamed, P.~Creminelli, S.~Mukohyama and M.~Zaldarriaga, %``Ghost inflation," JCAP\ {\bf 0404}, 001 (2004) [hep-th/0312100].

The non-standard dispersion relation results large non-Gaussinaity

$$S = \int d^4x \; \frac{1}{2} \dot{\pi}^2 - \frac{\alpha^2}{2M^2} (\nabla^2 \pi)^2$$

Conclusion

- Non-Gaussianity is the three point correlation function of temperature fluctuations on CMB.
- Different models predict different non-Gaussianity
- The detection of size, shape and sign of Non-Gaussianity in forthcoming observations constrains models of inflation.

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