Timelike Entanglement Entropy

Ali Mollabashi



Based on: Phys.Rev.Lett. 130 (2023) 3, 031601, [2210.09457] JHEP 05 (2023) 052, [2302.11695]

Collaborators: Kazuki Doi, Jonathan Harper, Tadashi Takayanagi, Yusuke Taki

IPM, IRCHEP (1402)

Ali Mollabashi

Timelike EE

Warm up: why TEE, ...

• Timelike entanglement "entropy" is a new *complex-valued* measure of information

Warm up: why TEE, ...

- Timelike entanglement "entropy" is a new *complex-valued* measure of information
- An early motivation: Concerns from dS/CFT; Is it possible to understand about the entanglement structure of theories dual to dS geometries via weak rotation of well-known results in AdS?

Warm up: why TEE, ...

- Timelike entanglement "entropy" is a new *complex-valued* measure of information
- An early motivation: Concerns from dS/CFT; Is it possible to understand about the entanglement structure of theories dual to dS geometries via weak rotation of well-known results in AdS?
- The content of this talk:
 - ▶ A reminder: Definition of pseudo entanglement
 - Definition of timelike entanglement in QFT
 - ▶ TEE is an example of pseudo entanglement
 - ▶ Holographic prescription to calculate TEE (v.1)
 - A comment about: TEE in AdS \Leftrightarrow EE in dS/CFT

• Entanglement entropy is defined for a single state $|\psi\rangle$

- Entanglement entropy is defined for a single state $|\psi\rangle$
- Pseudo Entropy is defined for two states $|\psi\rangle_1$ and $|\psi\rangle_2$

$$\rho = \frac{|\psi\rangle \langle \psi|}{\langle \psi|\psi\rangle} \longrightarrow \tau^{\psi_1|\psi_2} = \frac{|\psi_1\rangle \langle \psi_2|}{\langle \psi_2|\psi_1\rangle}$$

- Entanglement entropy is defined for a single state $|\psi\rangle$
- Pseudo Entropy is defined for two states $|\psi\rangle_1$ and $|\psi\rangle_2$

$$\rho = \frac{|\psi\rangle \langle \psi|}{\langle \psi|\psi\rangle} \longrightarrow \tau^{\psi_1|\psi_2} = \frac{|\psi_1\rangle \langle \psi_2|}{\langle \psi_2|\psi_1\rangle}$$

Pseudo entropy [Nakata-Takayanagi-Taki-Tamaoka-Wei '20]:

$$S(\tau_A) = -\operatorname{Tr}_A \left[\tau_A \log \tau_A \right] \quad , \quad \tau_A = \operatorname{Tr}_B \left[\tau_{AB}^{\psi_1 | \psi_2} \right]$$

- Entanglement entropy is defined for a single state $|\psi\rangle$
- Pseudo Entropy is defined for two states $|\psi\rangle_1$ and $|\psi\rangle_2$

$$\rho = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle} \longrightarrow \tau^{\psi_1|\psi_2} = \frac{|\psi_1\rangle\langle\psi_2|}{\langle\psi_2|\psi_1\rangle}$$

Pseudo entropy [Nakata-Takayanagi-Taki-Tamaoka-Wei '20]:

$$S(\tau_A) = -\operatorname{Tr}_A \left[\tau_A \log \tau_A \right] , \quad \tau_A = \operatorname{Tr}_B \left[\tau_{AB}^{\psi_1 | \psi_2} \right]$$

• $\tau_{AB}^{\psi_1|\psi_2}$ naturally arises after post-selection measurements

- Entanglement entropy is defined for a single state $|\psi\rangle$
- Pseudo Entropy is defined for two states $|\psi\rangle_1$ and $|\psi\rangle_2$

$$\rho = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle} \longrightarrow \tau^{\psi_1|\psi_2} = \frac{|\psi_1\rangle\langle\psi_2|}{\langle\psi_2|\psi_1\rangle}$$

Pseudo entropy [Nakata-Takayanagi-Taki-Tamaoka-Wei '20]:

$$S(\tau_A) = -\operatorname{Tr}_A \left[\tau_A \log \tau_A \right] , \quad \tau_A = \operatorname{Tr}_B \left[\tau_{AB}^{\psi_1 | \psi_2} \right]$$

- $\tau_{AB}^{\psi_1|\psi_2}$ naturally arises after post-selection measurements
- ▶ Weak value [Aharonov-Albert-Vaidman '88]

$$\frac{\langle \psi_1 | O_A | \psi_2 \rangle}{\langle \psi_1 | \psi_2 \rangle} = \operatorname{Tr} \left[O_A \tau_A^{2|1} \right]$$

Ali Mollabashi

- Entanglement entropy is defined for a single state $|\psi\rangle$
- Pseudo Entropy is defined for two states $|\psi\rangle_1$ and $|\psi\rangle_2$

$$\rho = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle} \longrightarrow \tau^{\psi_1|\psi_2} = \frac{|\psi_1\rangle\langle\psi_2|}{\langle\psi_2|\psi_1\rangle}$$

Pseudo entropy [Nakata-Takayanagi-Taki-Tamaoka-Wei '20]:

$$S(\tau_A) = -\operatorname{Tr}_A \left[\tau_A \log \tau_A \right] , \quad \tau_A = \operatorname{Tr}_B \left[\tau_{AB}^{\psi_1 | \psi_2} \right]$$

- $\tau_{AB}^{\psi_1|\psi_2}$ naturally arises after post-selection measurements
- ▶ Weak value [Aharonov-Albert-Vaidman '88]

$$\frac{\langle \psi_1 | O_A | \psi_2 \rangle}{\langle \psi_1 | \psi_2 \rangle} = \operatorname{Tr} \left[O_A \tau_A^{2|1} \right]$$

Start with $|\psi_1\rangle$, perform a measurement O_A , discard all outcomes except those which the final state is $|\psi_2\rangle$

Ali Mollabashi

Timelike EE







[Nakata-Takayanagi-Taki-Tamaoka-Wei '20]

• It is much easier to work with ρ^n rather than $\log \rho$

- It is much easier to work with ρ^n rather than $\log \rho$
- ▶ Replica method [Callan-Wilczek '94]

$$S_A^{(n)} = \frac{1}{1-n} \log \operatorname{Tr} [\rho_A^n] \quad , \quad \lim_{n \to 1} S_A^{(n)} = S_A$$

- It is much easier to work with ρ^n rather than $\log \rho$
- ▶ Replica method [Callan-Wilczek '94]

$$S_A^{(n)} = \frac{1}{1-n} \log \operatorname{Tr} [\rho_A^n] \quad , \quad \lim_{n \to 1} S_A^{(n)} = S_A$$

▶ In Euclidean formalism we have to calculate

$$\operatorname{Tr}\left[\rho_{A}^{n}\right] \propto \int_{\mathcal{R}_{n}} \mathcal{D}\phi \, e^{-S_{E}\left[\phi\right]} = \mathcal{Z}_{n}$$

- It is much easier to work with ρ^n rather than $\log \rho$
- Replica method [Callan-Wilczek '94]

$$S_A^{(n)} = \frac{1}{1-n} \log \operatorname{Tr} [\rho_A^n] \quad , \quad \lim_{n \to 1} S_A^{(n)} = S_A$$

▶ In Euclidean formalism we have to calculate

$$\operatorname{Tr}\left[\rho_{A}^{n}\right] \propto \int_{\mathcal{R}_{n}} \mathcal{D}\phi \, e^{-S_{E}\left[\phi\right]} = \mathcal{Z}_{n}$$

▶ The key point [Calabrese-Cardy '04]

$$\mathbb{C}$$
 in presence of $\mathcal{O}_{(n)} \Leftrightarrow \mathcal{R}_n$

- It is much easier to work with ρ^n rather than $\log \rho$
- Replica method [Callan-Wilczek '94]

$$S_A^{(n)} = \frac{1}{1-n} \log \operatorname{Tr} [\rho_A^n] \quad , \quad \lim_{n \to 1} S_A^{(n)} = S_A$$

▶ In Euclidean formalism we have to calculate

$$\operatorname{Tr}\left[\rho_{A}^{n}\right] \propto \int_{\mathcal{R}_{n}} \mathcal{D}\phi \, e^{-S_{E}\left[\phi\right]} = \mathcal{Z}_{n}$$

▶ The key point [Calabrese-Cardy '04]

$$\mathbb{C}$$
 in presence of $\mathcal{O}_{(n)} \Leftrightarrow \mathcal{R}_n$

• For A: a single interval

$$\mathcal{Z}_n \propto \langle \sigma_n \bar{\sigma}_n \rangle_{\mathbb{C}} \quad , \quad \Delta_n = \frac{c}{12} \left(n - \frac{1}{n} \right)$$

- It is much easier to work with ρ^n rather than $\log \rho$
- ▶ Replica method [Callan-Wilczek '94]

$$S_A^{(n)} = \frac{1}{1-n} \log \operatorname{Tr} [\rho_A^n] \quad , \quad \lim_{n \to 1} S_A^{(n)} = S_A$$

▶ In Euclidean formalism we have to calculate

$$\operatorname{Tr}\left[\rho_{A}^{n}\right] \propto \int_{\mathcal{R}_{n}} \mathcal{D}\phi \, e^{-S_{E}\left[\phi\right]} = \mathcal{Z}_{n}$$

▶ The key point [Calabrese-Cardy '04]

$$\mathbb{C}$$
 in presence of $\mathcal{O}_{(n)} \Leftrightarrow \mathcal{R}_n$

• For A: a single interval

$$\mathcal{Z}_n \propto \langle \sigma_n \bar{\sigma}_n \rangle_{\mathbb{C}}$$
, $\Delta_n = \frac{c}{12} \left(n - \frac{1}{n} \right)$
This leads to

$$\operatorname{Tr}\left[\rho_{A}^{n}\right] = \frac{\mathcal{Z}_{n}}{\mathcal{Z}_{1}^{n}} \propto \left(\frac{L_{A}}{\epsilon}\right)^{-\overline{6}\left(n-\frac{1}{n}\right)}$$

Ali Mollabashi

Timelike EE

• Reminder: For spacelike regions

$$S_A^{(n)} = \frac{1}{1-n} \log \langle \sigma_n(P) \bar{\sigma}_n(Q) \rangle$$
$$= \frac{1}{1-n} \log \left[\left(\frac{\epsilon}{\sqrt{(x_P - x_Q)^2 - (t_P - t_Q)^2}} \right)^{2\Delta_n} \right]$$

• Reminder: For spacelike regions

$$S_A^{(n)} = \frac{1}{1-n} \log \langle \sigma_n(P) \bar{\sigma}_n(Q) \rangle$$
$$= \frac{1}{1-n} \log \left[\left(\frac{\epsilon}{\sqrt{(x_P - x_Q)^2 - (t_P - t_Q)^2}} \right)^{2\Delta_n} \right]$$

▶ EE is given by

$$S_A = S_A^{(1)} = \frac{c}{3} \log \left[\frac{\sqrt{(\Delta x)^2 - (\Delta t)^2}}{\epsilon} \right]$$

• Reminder: For spacelike regions

$$S_A^{(n)} = \frac{1}{1-n} \log \langle \sigma_n(P) \bar{\sigma}_n(Q) \rangle$$
$$= \frac{1}{1-n} \log \left[\left(\frac{\epsilon}{\sqrt{(x_P - x_Q)^2 - (t_P - t_Q)^2}} \right)^{2\Delta_n} \right]$$

▶ EE is given by

$$S_A = S_A^{(1)} = \frac{c}{3} \log \left[\frac{\sqrt{(\Delta x)^2 - (\Delta t)^2}}{\epsilon} \right]$$

• Continue to $(\Delta x)^2 - (\Delta t)^2 < 0$, for pure timelike region T_0

$$S_A^{(\mathrm{T})} = \frac{c}{3}\log\frac{T_0}{\epsilon} + \frac{c\pi}{6}i$$

Ali Mollabashi

Timelike EE

• Finite Size: CFT on a circle (R), For a pure timelike region T_0

$$S_A^{(\mathrm{T})} = \frac{c}{3} \log \left[\frac{R}{\pi \epsilon} \sin \frac{\pi T_0}{R} \right] + \frac{i \pi c}{6}$$

• Finite Size: CFT on a circle (R), For a pure timelike region T_0

$$S_A^{(\mathrm{T})} = \frac{c}{3} \log \left[\frac{R}{\pi \epsilon} \sin \frac{\pi T_0}{R} \right] + \frac{i\pi c}{6}$$

• Finite Temperature: CFT at temperature $(1/\beta)$, For a pure timelike region T_0

$$S_A^{(\mathrm{T})} = \frac{c}{3} \log \left[\frac{\beta}{\pi\epsilon} \sinh \frac{\pi T_0}{\beta}\right] + \frac{i\pi c}{6}$$

• Consider a free scalar theory with $x \sim x + R$

▶ Consider a free scalar theory with $x \sim x + R$

$$S = \frac{1}{2} \int dt dx \left[(\partial_t \phi)^2 - (\partial_x \phi)^2 - m^2 \phi^2 \right]$$

the partition function is given by $Z_{\phi} = \int D\phi \, e^{iS}$

▶ Consider a free scalar theory with $x \sim x + R$

$$S = \frac{1}{2} \int dt dx \left[(\partial_t \phi)^2 - (\partial_x \phi)^2 - m^2 \phi^2 \right]$$

the partition function is given by Z_{ϕ} = $\int D\phi \, e^{iS}$

▶ Consider t: spatial direction, T = -ix: real-time



▶ The "Hamiltonian" reads

$$H = \frac{-i}{2} \int dt \left[\pi^2 + (\partial_t \phi)^2 - m^2 \phi^2 \right]$$

where $\pi = -\partial_x \phi$ is the canonical momentum such that

$$[\phi(t),\pi(t')] = i\delta(t-t')$$

▶ The "Hamiltonian" reads

$$H = \frac{-i}{2} \int dt \left[\pi^2 + (\partial_t \phi)^2 - m^2 \phi^2 \right]$$

where $\pi = -\partial_x \phi$ is the canonical momentum such that

$$[\phi(t),\pi(t')] = i\delta(t-t')$$

• In this formulation $(\tilde{H} = iH)$

$$Z_{\phi} = \mathrm{Tr}\left[e^{-RH}\right]$$

▶ The "Hamiltonian" reads

$$H = \frac{-i}{2} \int dt \left[\pi^2 + (\partial_t \phi)^2 - m^2 \phi^2 \right]$$

where $\pi = -\partial_x \phi$ is the canonical momentum such that

$$[\phi(t),\pi(t')] = i\delta(t-t')$$

▶ In this formulation $(\tilde{H} = iH)$

$$Z_{\phi} = \operatorname{Tr}\left[e^{-RH}\right] = \operatorname{Tr}\left[e^{iR\tilde{H}}\right]$$

▶ The "Hamiltonian" reads

$$H = \frac{-i}{2} \int dt \left[\pi^2 + (\partial_t \phi)^2 - m^2 \phi^2 \right]$$

where $\pi = -\partial_x \phi$ is the canonical momentum such that

$$[\phi(t),\pi(t')] = i\delta(t-t')$$

• In this formulation $(\tilde{H} = iH)$

$$Z_{\phi} = \operatorname{Tr}\left[e^{-RH}\right] = \operatorname{Tr}\left[e^{iR\tilde{H}}\right]$$

▶ Prescription: With $\beta_S \rightarrow -iR$ and $m \rightarrow -im$, we can find TEE from EE

Ali Mollabashi

Timelike EE

▶ Finite temperature EE to finite size TEE

$$S_A = \frac{c}{3} \log \left[\frac{\beta_S}{\pi \tilde{\epsilon}} \sinh \frac{\pi X_0}{\beta_S} \right] \xrightarrow[\tilde{\epsilon} \to -i\epsilon]{\tilde{\epsilon} \to -i\epsilon} S_A^{(T)} = \frac{c}{3} \log \left[\frac{R}{\pi \epsilon} \sin \frac{\pi T_0}{R} \right] + \frac{i\pi c}{6}$$

▶ Finite temperature EE to finite size TEE

$$S_A = \frac{c}{3} \log \left[\frac{\beta_S}{\pi \tilde{\epsilon}} \sinh \frac{\pi X_0}{\beta_S} \right] \xrightarrow[\tilde{\epsilon} \to -i\epsilon]{\tilde{\epsilon} \to -i\epsilon} S_A^{(T)} = \frac{c}{3} \log \left[\frac{R}{\pi \epsilon} \sin \frac{\pi T_0}{R} \right] + \frac{i\pi c}{6}$$

▶ Take $t \sim t - i\beta$ and $R \rightarrow \infty$ to define finite temperature TEE
Timelike EE in QFT II: Wick Rotation of Coordinates

▶ Finite temperature EE to finite size TEE

$$S_A = \frac{c}{3} \log \left[\frac{\beta_S}{\pi \tilde{\epsilon}} \sinh \frac{\pi X_0}{\beta_S} \right] \xrightarrow[\tilde{\epsilon} \to -i\epsilon]{\tilde{\epsilon} \to -i\epsilon} S_A^{(T)} = \frac{c}{3} \log \left[\frac{R}{\pi \epsilon} \sin \frac{\pi T_0}{R} \right] + \frac{i\pi c}{6}$$

- ▶ Take $t \sim t i\beta$ and $R \rightarrow \infty$ to define finite temperature TEE
- ▶ Finite size EE to finite temperature TEE

$$S_A = \frac{c}{3} \log \left[\frac{R_S}{\pi \tilde{\epsilon}} \sin \left(\frac{\pi X_0}{R_S} \right) \right] \xrightarrow[\tilde{\epsilon} \to -i\epsilon]{\tilde{\epsilon} \to -i\epsilon} S_A^{(\mathrm{T})} = \frac{c}{3} \log \left[\frac{R}{\pi \epsilon} \sinh \frac{\pi T_0}{R} \right] + \frac{i\pi c}{6}$$

TEE as Pseudo Entanglement Entropy

▶ The reduced density matrix corresponding to TEE was NOT hermitian (remember $H = i\tilde{H}$)

TEE as Pseudo Entanglement Entropy

- ▶ The reduced density matrix corresponding to TEE was NOT hermitian (remember $H = i\tilde{H}$)
- ▶ More explicitly consider the following purification

$$|\Psi\rangle = \frac{1}{\sqrt{Z(\delta)}} \sum_{n} e^{+i(R+i\delta)E_n/2} |n\rangle_1 |n\rangle_2$$

$$|\Psi^*\rangle = \frac{1}{\sqrt{Z(\delta)}} \sum_{n} e^{-i(R-i\delta)E_n/2} |n\rangle_1 |n\rangle_2$$

TEE as Pseudo Entanglement Entropy

- ▶ The reduced density matrix corresponding to TEE was NOT hermitian (remember $H = i\tilde{H}$)
- More explicitly consider the following purification

$$|\Psi\rangle = \frac{1}{\sqrt{Z(\delta)}} \sum_{n} e^{+i(R+i\delta)E_n/2} |n\rangle_1 |n\rangle_2$$

$$|\Psi^*\rangle = \frac{1}{\sqrt{Z(\delta)}} \sum_{n} e^{-i(R-i\delta)E_n/2} |n\rangle_1 |n\rangle_2$$

Tracing over one copy

$$\operatorname{Tr}_2|\Psi\rangle\langle\Psi^*| = e^{i(R+i\delta)\tilde{H}}$$

TEE is naturally expressed in terms of $\mathbf{pseudo}\ \mathbf{entropy}$

Ali Mollabashi

• Spectrum of $\tau_A \longrightarrow \text{TEE}$

- Spectrum of $\tau_A \longrightarrow \text{TEE}$
- How to read the spectrum of τ_A ?

$$\langle \mathcal{O}_A \rangle = \operatorname{Tr}_A \left[\tau_A \mathcal{O}_A \right]$$

- Spectrum of $\tau_A \longrightarrow \text{TEE}$
- How to read the spectrum of τ_A ?

$$\langle \mathcal{O}_A \rangle = \operatorname{Tr}_A \left[\tau_A \mathcal{O}_A \right]$$

▶ In free theories Wick's theorem \longrightarrow A unique way to real the spectrum of τ_A

- Spectrum of $\tau_A \longrightarrow \text{TEE}$
- How to read the spectrum of τ_A ?

$$\langle \mathcal{O}_A \rangle = \operatorname{Tr}_A \left[\tau_A \mathcal{O}_A \right]$$

- ▶ In free theories Wick's theorem \longrightarrow A unique way to real the spectrum of τ_A
- ▶ All we need are the two-point functions

- Spectrum of $\tau_A \longrightarrow \text{TEE}$
- How to read the spectrum of τ_A ?

$$\langle \mathcal{O}_A \rangle = \operatorname{Tr}_A \left[\tau_A \mathcal{O}_A \right]$$

- ▶ In free theories Wick's theorem \longrightarrow A unique way to real the spectrum of τ_A
- ▶ All we need are the two-point functions
- ▶ Similar to the standard formulation of KG theory

$$H = -i \int d^{d-1}k_y \, dk \, \Omega_{\mathbf{k}} \, a_{\mathbf{k}}^{\dagger} a_{\mathbf{k}}$$

• In 2*d* case $(\Omega_k = \sqrt{k^2 - m^2})$ the structure of $\operatorname{Tr}_A[\tau_A \mathcal{O}_A]$

$$\sum_{n} e^{iR\Omega_k n} f(n)$$

Ali Mollabashi

- Spectrum of $\tau_A \longrightarrow \text{TEE}$
- How to read the spectrum of τ_A ?

$$\langle \mathcal{O}_A \rangle = \operatorname{Tr}_A \left[\tau_A \mathcal{O}_A \right]$$

- ▶ In free theories Wick's theorem \longrightarrow A unique way to real the spectrum of τ_A
- ▶ All we need are the two-point functions
- ▶ Similar to the standard formulation of KG theory

$$H = -i \int d^{d-1}k_y \, dk \, \Omega_{\mathbf{k}} \, a_{\mathbf{k}}^{\dagger} a_{\mathbf{k}}$$

• In 2*d* case $(\Omega_k = \sqrt{k^2 - m^2})$ the structure of $\operatorname{Tr}_A[\tau_A \mathcal{O}_A]$

$$\sum_{n} e^{iR\Omega_k n} f(n) \longrightarrow \sum_{n} e^{i(R+i\delta)\Omega_k n} f(n)$$

Ali Mollabashi

- Spectrum of $\tau_A \longrightarrow \text{TEE}$
- How to read the spectrum of τ_A ?

$$\langle \mathcal{O}_A \rangle = \operatorname{Tr}_A \left[\tau_A \mathcal{O}_A \right]$$

- \blacktriangleright In free theories Wick's theorem \longrightarrow A unique way to real the spectrum of τ_A
- ▶ All we need are the two-point functions
- ▶ Similar to the standard formulation of KG theory

$$H = -i \int d^{d-1}k_y \, dk \, \Omega_{\mathbf{k}} \, a_{\mathbf{k}}^{\dagger} a_{\mathbf{k}}$$

• In 2*d* case $(\Omega_k = \sqrt{k^2 - m^2})$ the structure of $\operatorname{Tr}_A[\tau_A \mathcal{O}_A]$

$$\sum_{n} e^{iR\Omega_k n} f(n) \longrightarrow \sum_{n} e^{i(R+i\delta)\Omega_k n} f(n) = \sum_{n} e^{-(R+i\delta)\Omega_{ik} n} f(n)$$

Ali Mollabashi

• The relevant correlation functions (similarly for $\Pi_{tt'}$)

$$\Phi_{tt'} \equiv \operatorname{Tr}\left[e^{i(R+i\delta)\tilde{H}}\phi(t)\phi(t')\right] = \int \frac{dk}{2\pi} \frac{i}{2\Omega_{ik}} \operatorname{coth}\left(\frac{(R+i\delta)\Omega_{ik}}{2}\right) e^{ik(t-t')}$$

• The relevant correlation functions (similarly for $\Pi_{tt'}$)

$$\Phi_{tt'} \equiv \operatorname{Tr}\left[e^{i(R+i\delta)\tilde{H}}\phi(t)\phi(t')\right] = \int \frac{dk}{2\pi} \frac{i}{2\Omega_{ik}} \operatorname{coth}\left(\frac{(R+i\delta)\Omega_{ik}}{2}\right) e^{ik(t-t')}$$



• The relevant correlation functions (similarly for $\Pi_{tt'}$)

$$\Phi_{tt'} \equiv \operatorname{Tr}\left[e^{i(R+i\delta)\tilde{H}}\phi(t)\phi(t')\right] = \int \frac{dk}{2\pi} \frac{i}{2\Omega_{ik}} \operatorname{coth}\left(\frac{(R+i\delta)\Omega_{ik}}{2}\right) e^{ik(t-t')}$$



Strong numerical evidence for the imaginary part

• The relevant correlation functions (similarly for $\Pi_{tt'}$)

$$\Phi_{tt'} \equiv \operatorname{Tr}\left[e^{i(R+i\delta)\tilde{H}}\phi(t)\phi(t')\right] = \int \frac{dk}{2\pi} \frac{i}{2\Omega_{ik}} \operatorname{coth}\left(\frac{(R+i\delta)\Omega_{ik}}{2}\right) e^{ik(t-t')}$$



- ▶ Strong numerical evidence for the imaginary part
- ▶ Numerical results also agree in the finite temperature case

Holographic Timelike EE in Pure AdS_3

▶ Poincare AdS₃

$$ds^2 = \frac{dz^2 - dt^2 + dx^2}{z^2},$$

with continuation $X_0 \rightarrow iT_0$

$$t = \sqrt{z^2 + (T_0/2)^2}$$

Holographic Timelike EE in Pure AdS_3

▶ Poincare AdS₃

$$ds^2 = \frac{dz^2 - dt^2 + dx^2}{z^2},$$

with continuation $X_0 \rightarrow iT_0$

$$t = \sqrt{z^2 + (T_0/2)^2}$$



Holographic Timelike EE in Pure AdS_3

▶ Poincare AdS₃

$$ds^2 = \frac{dz^2 - dt^2 + dx^2}{z^2},$$

with continuation $X_0 \rightarrow iT_0$

$$t = \sqrt{z^2 + (T_0/2)^2}$$



Timelike part comes from the homology condition

$$S_A^{(\mathrm{T})} = \frac{c}{3} \log\left(\frac{T_0}{\epsilon}\right) + \frac{c}{6}\pi i$$

Ali Mollabashi

• Can a union of timelike and spacelike geodesics result from an extremizing prescription?

- Can a union of timelike and spacelike geodesics result from an extremizing prescription?
- AdS_2 slicing of AdS_3 shows an affirmative answer

$$ds^{2} = d\eta^{2} + \cosh^{2}\eta \left(-\cosh^{2}rdt^{2} + dr^{2}\right)$$

- Can a union of timelike and spacelike geodesics result from an extremizing prescription?
- \blacktriangleright AdS₂ slicing of AdS₃ shows an affirmative answer

$$ds^{2} = d\eta^{2} + \cosh^{2}\eta \left(-\cosh^{2}rdt^{2} + dr^{2}\right)$$

we analytically continue $\eta < 0$ by $\eta = i \tilde{\eta}$

$$ds^{2} = -d\tilde{\eta}^{2} + \cos^{2}\tilde{\eta}\left(-\cosh^{2}rdt^{2} + dr^{2}\right), \quad \left(0 < \tilde{\eta} < \pi/2\right)$$

- Can a union of timelike and spacelike geodesics result from an extremizing prescription?
- AdS_2 slicing of AdS_3 shows an affirmative answer

$$ds^{2} = d\eta^{2} + \cosh^{2}\eta \left(-\cosh^{2}rdt^{2} + dr^{2}\right)$$

we analytically continue $\eta < 0$ by $\eta = i \tilde{\eta}$

$$ds^2 = -d\tilde{\eta}^2 + \cos^2\tilde{\eta} \left(-\cosh^2 r dt^2 + dr^2 \right), \quad \left(0 < \tilde{\eta} < \pi/2 \right)$$



$$ds^2 = -(d\tilde{\eta}^2 + \cos^2\tilde{\eta}dt^2), \quad (0 < \tilde{\eta} < \pi/2)$$

• Due to symmetry Γ_A is expected to be on

$$ds^2 = -(d\tilde{\eta}^2 + \cos^2\tilde{\eta}dt^2), \quad (0 < \tilde{\eta} < \pi/2)$$

▶ For compactified t this is a $-ds_{S^2}^2$ with imaginary length

$$ds^2 = -(d\tilde{\eta}^2 + \cos^2\tilde{\eta}dt^2), \quad (0 < \tilde{\eta} < \pi/2)$$

- ▶ For compactified t this is a $-ds_{S^2}^2$ with imaginary length
- Extremizing w.r.t the points on $\eta = 0$ leads to the geodesic along the boundary of S^2 (with length $i\pi$)

$$ds^{2} = -(d\tilde{\eta}^{2} + \cos^{2}\tilde{\eta}dt^{2}), \quad (0 < \tilde{\eta} < \pi/2)$$

- ▶ For compactified t this is a $-ds_{S^2}^2$ with imaginary length
- Extremizing w.r.t the points on $\eta = 0$ leads to the geodesic along the boundary of S^2 (with length $i\pi$)
- \Rightarrow the real part is the sum of geodesics with fixed endpoints

$$ds^2 = -(d\tilde{\eta}^2 + \cos^2\tilde{\eta}dt^2), \quad (0 < \tilde{\eta} < \pi/2)$$

- ▶ For compactified t this is a $-ds_{S^2}^2$ with imaginary length
- Extremizing w.r.t the points on $\eta = 0$ leads to the geodesic along the boundary of S^2 (with length $i\pi$)
- $\blacktriangleright \Rightarrow$ the real part is the sum of geodesics with fixed endpoints
- Prescription summary:
 - 1. Construct candidates of Γ_A from a union of spacelike and timelike geodesics such that $\partial \Gamma_A = \partial A$
 - 2. (Considering Wick rotated geometry for the timelike geodesics), require all variations of the joining points to be stationary

$$ds^{2} = -4\frac{dudv}{(1+uv)^{2}} + \frac{(1-uv)^{2}}{(1+uv)^{2}}r_{+}^{2}d\phi^{2}$$

$$ds^{2} = -4\frac{dudv}{(1+uv)^{2}} + \frac{(1-uv)^{2}}{(1+uv)^{2}}r_{+}^{2}d\phi^{2}$$

Solving the geodesic equation u''(1+uv) + 2u'(u-u'v) = 0• How extremization works?

$$\begin{cases} (a_1, -\frac{1}{a_1}) & \text{subregion bdy} \\ (a_2, -\frac{1}{a_2}) & \text{subregion bdy} \\ (s, \frac{1}{s}) & \text{future singularity} \\ (-q, -\frac{1}{q}) & \text{past singularity} \end{cases}$$

$$ds^{2} = -4\frac{dudv}{(1+uv)^{2}} + \frac{(1-uv)^{2}}{(1+uv)^{2}}r_{+}^{2}d\phi^{2}$$

Solving the geodesic equation u''(1+uv) + 2u'(u-u'v) = 0• How extremization works?

$$\begin{cases} (a_1, -\frac{1}{a_1}) & \text{subregion bdy} \\ (a_2, -\frac{1}{a_2}) & \text{subregion bdy} \\ (s, \frac{1}{s}) & \text{future singularity} \\ (-q, -\frac{1}{q}) & \text{past singularity} \end{cases}$$

• Existence of timelike geodesic \Rightarrow s = q

Ali Mollabashi

$$ds^{2} = -4\frac{dudv}{(1+uv)^{2}} + \frac{(1-uv)^{2}}{(1+uv)^{2}}r_{+}^{2}d\phi^{2}$$

Solving the geodesic equation u''(1+uv) + 2u'(u-u'v) = 0• How extremization works?

$$\begin{cases} (a_1, -\frac{1}{a_1}) & \text{subregion bdy} \\ (a_2, -\frac{1}{a_2}) & \text{subregion bdy} \\ (s, \frac{1}{s}) & \text{future singularity} \\ (-q, -\frac{1}{q}) & \text{past singularity} \end{cases}$$

- Existence of timelike geodesic \Rightarrow s = q
- Extremizing w.r.t. $s \Rightarrow s^2 = a_1 a_2$

Ali Mollabashi

Putting three pieces together

$$S_A^{(\mathrm{T})} = \frac{c}{3} \log \left(\frac{\beta}{\pi \delta} \sinh \left(\frac{\pi}{\beta} (T_2 - T_1) \right) \right) + \frac{c}{6} i \pi$$



• Can we fully understand the entanglement structure of TFD states without TEE?

- Can we fully understand the entanglement structure of TFD states without TEE?
- How to interpret: σ_n in CFT_R and $\overline{\sigma}_n$ in CFT_L ?

- Can we fully understand the entanglement structure of TFD states without TEE?
- How to interpret: σ_n in CFT_R and $\bar{\sigma}_n$ in CFT_L ?



- Can we fully understand the entanglement structure of TFD states without TEE?
- How to interpret: σ_n in CFT_R and $\overline{\sigma}_n$ in CFT_L ?



• There is no spacelike A that σ_n and $\bar{\sigma}_n$ are at ∂A The Hilbert space is constructed from fields on disconnected circles **cannot** be interpreted as EE!
Necessity of TEE in TFD States

- Can we fully understand the entanglement structure of TFD states without TEE?
- How to interpret: σ_n in CFT_R and $\overline{\sigma}_n$ in CFT_L ?



- There is no spacelike A that σ_n and $\bar{\sigma}_n$ are at ∂A The Hilbert space is constructed from fields on disconnected circles **cannot** be interpreted as EE!
- ▶ This configuration should be interpreted as TEE!

• Is timelike EE related to entropy in field theories dual to de Sitter spacetime?

- Is timelike EE related to entropy in field theories dual to de Sitter spacetime?
- ► dS / CFT is conjectured for the theory leaving on S^d part of dS space $ds^2 = R_{\rm dS}^2 (-d\tau^2 + \cosh^2 \tau \, d\Omega_d^2)$

and a Euclidean CFT_d

 The dictionary analogous to the GKPW [Gubser-Klebanov-Polyakov, Witten '98]

$$Z_{\rm CFT}[\phi_0] = \Psi_{\rm dS}[\phi_0] = \int_{\phi|_{\tau=\tau_{\infty}}=\phi_0} \mathcal{D}\phi \, e^{iI_{\rm dS}[\phi]} \Psi_{\rm in}$$

- Is timelike EE related to entropy in field theories dual to de Sitter spacetime?
- ► dS / CFT is conjectured for the theory leaving on S^d part of dS space $ds^2 = R_{\rm dS}^2 (-d\tau^2 + \cosh^2 \tau \, d\Omega_d^2)$

and a Euclidean CFT_d

 The dictionary analogous to the GKPW [Gubser-Klebanov-Polyakov, Witten '98]

$$Z_{\rm CFT}[\phi_0] = \Psi_{\rm dS}[\phi_0] = \int_{\phi|_{\tau=\tau_\infty}=\phi_0} \mathcal{D}\phi \, e^{iI_{\rm dS}[\phi]} \Psi_{\rm in}$$

 Conjecture: CFT is non-unitary with c_{dS} = ic_{AdS} (2d) [Maldacena '02]

- Is timelike EE related to entropy in field theories dual to de Sitter spacetime?
- ► dS / CFT is conjectured for the theory leaving on S^d part of dS space $ds^2 = R_{\rm dS}^2 (-d\tau^2 + \cosh^2 \tau \, d\Omega_d^2)$

and a Euclidean CFT_d

 The dictionary analogous to the GKPW [Gubser-Klebanov-Polyakov, Witten '98]

$$Z_{\rm CFT}[\phi_0] = \Psi_{\rm dS}[\phi_0] = \int_{\phi|_{\tau=\tau_{\infty}}=\phi_0} \mathcal{D}\phi \, e^{iI_{\rm dS}[\phi]} \Psi_{\rm in}$$

- Conjecture: CFT is non-unitary with c_{dS} = ic_{AdS} (2d) [Maldacena '02]
- non-unitarity $\longrightarrow \langle \psi | \neq | \psi \rangle^{\dagger}$ $\rho = | \psi \rangle \langle \psi |$ is a transition matrix

Holographic Pseudo Entropy in dS_3/CFT_2

• In global dS_3

$$ds^{2} = R_{\rm dS}^{2}(-d\tau^{2} + \cosh^{2}\tau(dt_{\rm E}^{2} + \cos^{2}t_{\rm E}d\theta^{2})), \qquad (\tau > 0)$$

with the Euclidean part corresponding to the initial state

$$ds^{2} = R_{\rm dS}^{2} (d\tau_{\rm E}^{2} + \cos^{2} \tau_{\rm E} (dt_{\rm E}^{2} + \cos^{2} t_{\rm E} d\theta^{2})), \quad (0 < \tau_{\rm E} < \pi).$$

we consider A at τ_{∞} on $t_E = 0$

Holographic Pseudo Entropy in dS_3/CFT_2

► In global dS₃

$$ds^{2} = R_{\rm dS}^{2}(-d\tau^{2} + \cosh^{2}\tau(dt_{\rm E}^{2} + \cos^{2}t_{\rm E}d\theta^{2})), \qquad (\tau > 0)$$

with the Euclidean part corresponding to the initial state

$$ds^{2} = R_{\rm dS}^{2} (d\tau_{\rm E}^{2} + \cos^{2} \tau_{\rm E} (dt_{\rm E}^{2} + \cos^{2} t_{\rm E} d\theta^{2})), \quad (0 < \tau_{\rm E} < \pi).$$

we consider A at τ_{∞} on $t_E = 0$



Holographic Pseudo Entropy in dS_3/CFT_2

► In global dS₃

$$ds^{2} = R_{\rm dS}^{2} (-d\tau^{2} + \cosh^{2}\tau (dt_{\rm E}^{2} + \cos^{2}t_{\rm E}d\theta^{2})), \qquad (\tau > 0)$$

with the Euclidean part corresponding to the initial state

$$ds^{2} = R_{\rm dS}^{2} (d\tau_{\rm E}^{2} + \cos^{2}\tau_{\rm E} (dt_{\rm E}^{2} + \cos^{2}t_{\rm E}d\theta^{2})), \quad (0 < \tau_{\rm E} < \pi).$$

we consider A at τ_{∞} on $t_E = 0$



Extrimization of timelike and spacelike parts gives

$$S_A^{(\mathrm{P})} = -i\frac{c_{\mathrm{dS}}}{3}\log\left(\frac{2\sin\frac{\phi_0}{2}}{\epsilon_{\mathrm{dS}}}\right) + \frac{\pi c_{\mathrm{dS}}}{6}, \quad (\epsilon_{\mathrm{dS}} \equiv 2e^{-\tau_{\infty}})$$

Ali Mollabashi

Timelike EE

▶ Starting from Euclidean AdS

$$ds^{2} = R_{\text{AdS}}^{2} \frac{dz^{2} + dt_{\text{E}}^{2} + dx^{2}}{z^{2}} \quad , \quad S_{A} = \frac{c_{\text{AdS}}}{3} \log\left(\frac{T_{0}}{\epsilon_{\text{AdS}}}\right)$$

Starting from Euclidean AdS

$$ds^{2} = R_{AdS}^{2} \frac{dz^{2} + dt_{E}^{2} + dx^{2}}{z^{2}} , \quad S_{A} = \frac{c_{AdS}}{3} \log\left(\frac{T_{0}}{\epsilon_{AdS}}\right)$$

The following "double Wick rotation" gives the same result

Starting from Euclidean AdS

$$ds^2 = R_{\text{AdS}}^2 \frac{dz^2 + dt_{\text{E}}^2 + dx^2}{z^2} \quad , \quad S_A = \frac{c_{\text{AdS}}}{3} \log\left(\frac{T_0}{\epsilon_{\text{AdS}}}\right)$$

The following "double Wick rotation" gives the same result 1. $z = -i\eta$, $R_{AdS} = -iR_{dS}$ ($\epsilon_{AdS} = -i\epsilon_{dS}$, $c_{AdS} = -ic_{dS}$)

Starting from Euclidean AdS

$$ds^{2} = R_{AdS}^{2} \frac{dz^{2} + dt_{E}^{2} + dx^{2}}{z^{2}} , \quad S_{A} = \frac{c_{AdS}}{3} \log\left(\frac{T_{0}}{\epsilon_{AdS}}\right)$$

The following "double Wick rotation" gives the same result

1.
$$z = -i\eta$$
, $R_{AdS} = -iR_{dS}$ ($\epsilon_{AdS} = -i\epsilon_{dS}$, $c_{AdS} = -ic_{dS}$)
2. $t_E = it$ ($T_0 \rightarrow iT_0$)

Starting from Euclidean AdS

$$ds^{2} = R_{\text{AdS}}^{2} \frac{dz^{2} + dt_{\text{E}}^{2} + dx^{2}}{z^{2}} \quad , \quad S_{A} = \frac{c_{\text{AdS}}}{3} \log\left(\frac{T_{0}}{\epsilon_{\text{AdS}}}\right)$$

The following "double Wick rotation" gives the same result 1. $z = -i\eta$, $R_{AdS} = -iR_{dS}$ ($\epsilon_{AdS} = -i\epsilon_{dS}$, $c_{AdS} = -ic_{dS}$) 2. $t_E = it (T_0 \rightarrow iT_0)$





- ▶ TEE defined by the Wick rotation of
 - 1. the replica trick results
 - 2. the coordinates

lead to the same results in CFT_2

- ▶ TEE defined by the Wick rotation of
 - 1. the replica trick results
 - 2. the coordinates

lead to the same results in CFT_2

 Using def. 2, the analytic continuation and the direct (numerical) calculation using the *real-time* approach perfectly agree

- ▶ TEE defined by the Wick rotation of
 - 1. the replica trick results
 - 2. the coordinates

lead to the same results in CFT_2

- Using def. 2, the analytic continuation and the direct (numerical) calculation using the *real-time* approach perfectly agree
- For holographic theories, there is an *extermination* procedure justified by extermination the geodesics with real and imaginary values

- ▶ TEE defined by the Wick rotation of
 - 1. the replica trick results
 - 2. the coordinates

lead to the same results in CFT_2

- Using def. 2, the analytic continuation and the direct (numerical) calculation using the *real-time* approach perfectly agree
- For holographic theories, there is an *extermination* procedure justified by extermination the geodesics with real and imaginary values
- Pseudo Entropy in field theories dual to dS spacetime followed by a "*double Wick rotation*" is the same as TEE in field theories dual to AdS

Holographic TEE: Geodesics in Global AdS₃

• Geodesic length between $(\rho_{\infty}, \frac{T_0}{2}, \phi_0)$ and $(\rho_{\infty}, -\frac{T_0}{2}, \phi_0)$

$$D = \cosh^{-1} \left(\cosh^2 \rho_{\infty} \cos T_0 - \sinh^2 \rho_{\infty} \right)$$
$$\simeq \pi i + \log \left(\frac{\sin^2 \frac{T_0}{2}}{\epsilon^2} \right),$$

(~: leading order in cut-off $\epsilon = e^{-\rho_{\infty}}$)

Holographic TEE: Geodesics in Global AdS₃

• Geodesic length between $(\rho_{\infty}, \frac{T_0}{2}, \phi_0)$ and $(\rho_{\infty}, -\frac{T_0}{2}, \phi_0)$

$$D = \cosh^{-1} \left(\cosh^2 \rho_{\infty} \cos T_0 - \sinh^2 \rho_{\infty} \right)$$
$$\simeq \pi i + \log \left(\frac{\sin^2 \frac{T_0}{2}}{\epsilon^2} \right),$$

(~: leading order in cut-off $\epsilon = e^{-\rho_\infty})$

▶ Timelike part comes from the homology condition

$$S_A^{(\mathrm{T})} = \frac{c}{3} \log \left(\frac{2 \sin \frac{T_0}{2}}{\epsilon} \right) + \frac{c}{6} \pi i$$

same as CFT result with $R = 2\pi$

Ali Mollabashi

Numerical Results (Imaginary Part): Finite Size R=100, $\epsilon=1$, $m=10^{-8}$



Ali Mollabashi

Timelike EE