

# On Complexity Growth for $F(R)$ and Critical Gravity

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Based on:

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# Outline

- 1 Computational Complexity
- 2 Complexity growth in higher curvatures gravities
- 3 Sum up and Outlook

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Black holes are nature's extreme memory!

Holographic bound implies amount of information can be stored in a region is limited by area of the region  $S \leq \frac{A}{4G}$ .

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For a system with  $N$  dof., the scrambling time  $t_*$  is a measure of how long it takes for information about a small  $\mathcal{O}(1)$  perturbation to spread over  $\mathcal{O}(N)$  dof.

Typical system:  $t_* \sim N^{\frac{1}{d}}$       Black holes:  $t_* = \frac{1}{2\pi T} \log S \sim \log N$

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Are black holes also the fastest computer?

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- Is there any bound on the computational speed?

## Lloyd's Conjecture

$$\text{max. rate of computation} \leq E$$

# Computation and Complexity

- How difficult is it to do a task?
- How difficult is it to prepare a particular state?

## Computational Complexity

Minimum number of quantum gates (operations) required to prepare the desired state.

- It grows linearly with time even after thermalization!
- It is bounded by Lloyd,s bound:  $\frac{dC(|\psi(t)\rangle)}{dt} \leq \frac{2\text{Energy}}{\pi\hbar}$   
[Brown, Roberts,Susskind, Swingle,Zhao'15]

# Holographic Complexity

- There are (at least) two proposals for holographic complexity:

Complexity=Volume (CV) [Stanford, Susskind'14]

$$\text{Complexity} \sim \frac{\text{Volume}}{G\ell}$$

Complexity=Action (CA) [Brown, Roberts, Susskind, Swingle, Zhao'15]

$$\text{Complexity} \sim \frac{\text{Action}}{\pi\hbar}$$

See also complexity of reduced state [Alishahiha'15] and complexity=volume version 2.0 [Couch, Fischler, Nguyen'17]

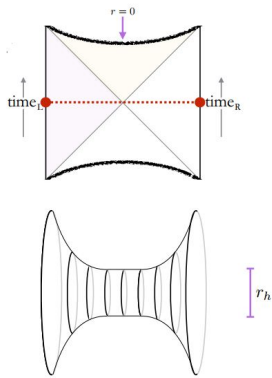
# Complexity=Volume (CV) [Susskind, Stanford'14]

## CV- duality

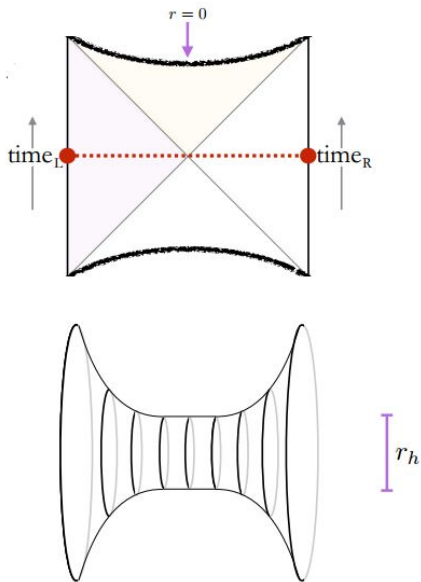
The complexity of the boundary state is proportional to the spatial volume  $V$  of a maximal slice behind the horizon  $\mathcal{C} \sim \frac{\text{Volume}}{G\ell}$ .

$\ell$  is unfixed length scale that has to be chosen appropriately. It is typically AdS radius of horizon radius.

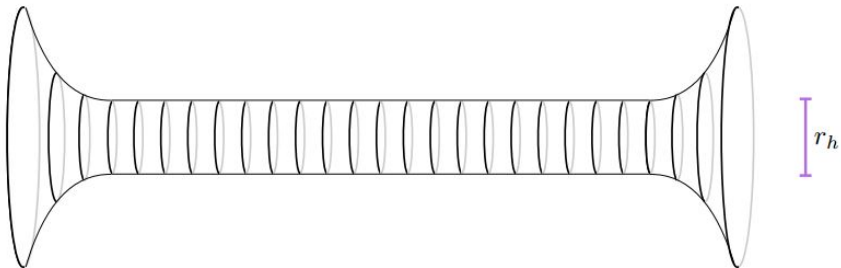
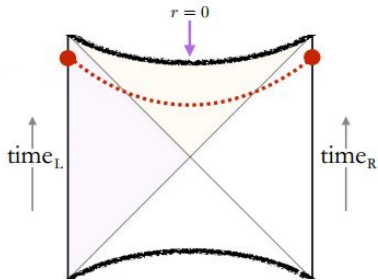
- It is a criterion of wormhole (Einstein-Rosen bridge) size.
- It grows linearly with time after thermalization:  $\mathcal{C}_V \sim TS(t_L + t_R)$
- It passes several tests.
- It has two unpleasant features:
  - Arbitrary scale  $\ell$
  - Why should the maximal slice play a preferred role?







[I've borrowed this pic from Brown's slides!]



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# Complexity=Action (CA) [Brown, Roberts,Susskind, Swingle,Zhao'15]

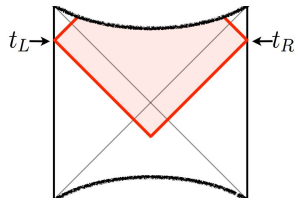
## CA- duality

The complexity of the boundary state is equal to the classical action of Wheeler-DeWitt patch.

Wheeler-DeWitt (WDW) patch is defined as the bulk domain of dependence of a Cauchy slice anchored at the boundary state.

$$C(|\psi(t_R, t, L)\rangle\rangle) = \frac{\text{Action}}{\pi\hbar}$$

$$\frac{d}{dt}C(|\psi(t_R, t_L)\rangle\rangle) \geq \frac{2E}{\pi\hbar}$$



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# $F(R)$ gravity

$$I = \frac{1}{16\pi G} \int_{\mathcal{M}} d^{d+2}x \sqrt{-g} F(R) + \frac{1}{8\pi G} \int_{\partial\mathcal{M}} d^{d+1}x \sqrt{-\gamma} F'(R) K,$$

where  $F'(R) = \frac{\partial F(R)}{\partial R}$ .

$$F'(R)R_{\mu\nu} - \frac{1}{2}F(R)g_{\mu\nu} - (\nabla_{\mu}\nabla_{\nu} - g_{\mu\nu}\square)F'(R) = 0.$$

- AdS-Schwarzschild is a solution to this equation.

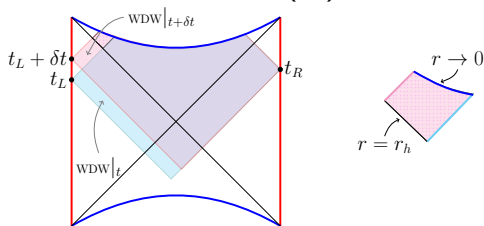
$$ds^2 = -f dt^2 + f^{-1} dr^2 + r^2 d\Omega_d^2, \quad f(r) = 1 + \frac{r^2}{\ell^2} + \frac{M_0}{r^{d-1}}$$

ADM mass of this solution is given as

$$\mathcal{M} = \frac{\Omega_d d}{16\pi G \ell^2} (\ell^2 + r_h^2) F'(R) r_h^{d-1}$$

# Complexity/Action Growth in $F(R)$

- To evaluate the on-shell action on the WDW patch, besides the generalized Gibbons-Hawking terms, one also needs boundary terms for null boundaries and intersections.
- In general it is very hard to find them. However, they are known for Einstein gravity [Parattu,Chakraborty, Majhi Padmanabhan ], [Lehner, Myers ,Poisson ,Sorkin'16].
- As far as the late time behavior of the complexity growth is concerned one may circumvent this challenge by following the approach considered in [Complexity=Action (CA)[Brown, Roberts,Susskind, Swingle,Zhao'15]].
- It matches with more rigorous calculation based on all boundary terms (it has been shown for Einstein theory [Lehner, Myers ,Poisson ,Sorkin'16] )
- I will calculate the rate of increase of action of a WDW patch of the two-sided black hole (which could be dual to the rate of growth of complexity of the boundary state).

Complexity/Action Growth in  $F(R)$  cont.

- Bulk term**

$$\begin{aligned} \delta I_{\mathcal{M}} &= I_{\mathcal{M}} [\text{WDW}|_{t+\delta t}] - I_{\mathcal{M}} [\text{WDW}|_t] \\ &= \frac{1}{16\pi G} \int_{r \rightarrow 0}^{r_h} \int_t^{t+\delta t} \int_{S^d} \mathcal{L}_{F(R)} d\Omega_d dt dr \quad (\text{at late time}) \end{aligned}$$

- Boundary term**

$$\delta I_{\partial\mathcal{M}} = \frac{1}{8\pi G} \int_t^{t+\delta t} \int_{\Omega^d} dt d\Omega_d K F'(R) \Big|_{r \rightarrow 0}^{r_h}$$

## Complexity/Action Growth in $F(R)$ cont.

- Bulk Term

$$\delta I_{\mathcal{M}} = \frac{\Omega_d F(R)}{16\pi G(d+1)} r_h^{d+1} \delta t$$

- Boundary Term

$$\delta I_{\partial\mathcal{M}} = -\frac{\Omega_d F(R)}{16\pi G(d+1)} ((d+1)r_h^2 + d\ell^2) r_h^{d-1} \delta t.$$

- ADM mass

$$\mathcal{M} = \frac{\Omega_d d}{16\pi G \ell^2} (\ell^2 + r_h^2) F'(R) r_h^{d-1}$$

- Bulk + Boundary

$$\delta(I_{\mathcal{M}} + I_{\partial\mathcal{M}}) = 2\mathcal{M}\delta t$$

### Complexity growth

It saturates complexity growth bound  $\dot{\mathcal{C}} = \frac{\delta I}{\delta t} = 2\mathcal{M}$ .



# Critical Gravity: Bulk Action

$$I_{\mathcal{M}} = \frac{1}{16\pi G} \int_{\mathcal{M}} d^{d+2}x \sqrt{-g} \left[ R - 2\Lambda - \frac{1}{m^2} \left( R^{\mu\nu} R_{\mu\nu} - \frac{d+2}{4(d+1)} R^2 \right) \right],$$

- $m$  is a dimensionful parameter.

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- $m$  is a dimensionful parameter.
- It admits AdS and AdS black holes with radius  $\ell$ .
- It is known that at the critical point  $m^2 = \frac{d^2}{2\ell^2}$  the model degenerates yielding to a log-gravity [Alishahiha, Fareghbal'11]

## Critical Gravity: Boundary Term

bulk term

$$I_{\mathcal{M}} = \frac{1}{16\pi G} \int_{\mathcal{M}} d^{d+2}x \sqrt{-g} \left[ R - 2\Lambda - \frac{1}{m^2} \left( R^{\mu\nu} R_{\mu\nu} - \frac{d+2}{4(d+1)} R^2 \right) \right],$$

boundary term [Hohm, Tonni'10]

$$I_{\partial\mathcal{M}} = \frac{1}{16\pi G} \int_{\partial\mathcal{M}} d^{d+1}x \sqrt{-\gamma} \left( -2K - \hat{f}^{ij} K_{ij} + \hat{f} K \right).$$

where

$$\hat{f}^{ij} = f^{ij} + 2h^{(i} N^{j)} + s N^i N^j,$$

$$f^{\mu\nu} = \begin{pmatrix} s & h^i \\ h^i & f^{ij} \end{pmatrix}. \quad f_{\mu\nu} = -\frac{2}{m^2} \left( R_{\mu\nu} - \frac{1}{2(d+1)} R g_{\mu\nu} \right).$$

and  $N^i$  is shift function in ADM decomposition.

# Critical Gravity: Static Solution

- AdS-Schwarzschild is a solution to this theory.

$$ds^2 = -f dt^2 + f^{-1} dr^2 + r^2 d\Omega_d^2, \quad f(r) = 1 + \frac{r^2}{\ell^2} + \frac{16\pi G}{\Omega_d d} \frac{M_0}{r^{d-1}}$$
$$\frac{2\Lambda}{d(d+1)} \ell^4 + \ell^2 - \frac{d(d-2)}{4m^2} = 0.$$

- The ADM mass/energy of this solution is

$$\mathcal{M} = M_0 \left( 1 - \frac{d^2}{2\ell^2 m^2} \right)$$

# Complexity growth in critical gravity: Static

- bulk term

$$\delta I_{\mathcal{M}} = -\frac{\Omega_d r_h^{d+1}}{8\pi G \ell^2} \left(1 - \frac{d^2}{2\ell^2 m^2}\right) \delta t$$

- boundary term term

$$\delta I_{\partial\mathcal{M}} = \frac{\Omega_d}{8\pi G} \left( d r_h^{d-1} + \frac{(d+1)r_h^{d+1}}{\ell^2} \right) \left(1 - \frac{d^2}{2\ell^2 m^2}\right) \delta t$$

- bulk + boundary  $\delta(I_{\mathcal{M}} + I_{\partial\mathcal{M}}) = 2\mathcal{M}\delta t$

## Complexity growth

It saturates complexity growth bound  $\dot{C} = 2\mathcal{M} = 2M_0(1 - \frac{d^2}{2\ell^2 m^2})$ .

at the critical point  $m^2 = \frac{d^2}{2\ell^2}$  where the model develops a log gravity, the rate of growth vanishes!

## 3D Critical Gravity (New Massive Gravity): Stationary

- 3d critical gravity is known as New Massive Gravity (NMG).
- BTZ black hole is a solution to this theory

$$ds^2 = \frac{dr^2}{f(r)} - f(r)dt^2 + r^2 \left( d\phi - \frac{8GJ_0}{2r^2} dt \right)^2,$$
$$f(r) = \frac{r^2}{\ell^2} - 8GM_0 + \frac{(8GJ_0)^2}{4r^2},$$

- The ADM mass and angular momentum of this solution is

$$\mathcal{M} = M_0 \left( 1 - \frac{1}{2\ell^2 m^2} \right), \quad \mathcal{J} = J_0 \left( 1 - \frac{1}{2\ell^2 m^2} \right).$$

- Note that at the critical point,  $m^2 = \frac{1}{2\ell^2}$  where the model develops a log gravity, the rate of growth vanishes.

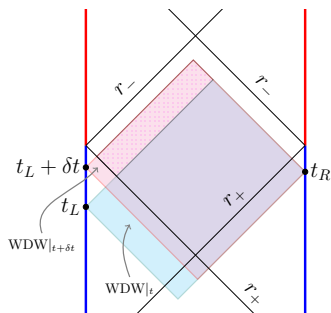
# Complexity growth in NMG: Stationary (BTZ)

- bulk term

$$\begin{aligned} \delta I_{\mathcal{M}} &= I_{\mathcal{M}}[\text{WDW}|_{t+\delta t}] - I_{\mathcal{M}}[\text{WDW}|_t] \\ &= \int_{r_-}^{r_+} \int_t^{t+\delta t} \int_0^{2\pi} \mathcal{L}_{\text{NMG}} \quad (\text{at late time}) \end{aligned}$$

- boundary term

$$\delta I_{\partial\mathcal{M}} = \int_t^{t+\delta t} \int_0^{2\pi} (\text{Boundary Term}) \Big|_{r \rightarrow 0}^{r_h}$$



## Complexity growth

$$\dot{C} = 2\sqrt{\mathcal{M}^2 - \frac{\mathcal{J}^2}{\ell^2}} < 2\mathcal{M}. \text{ It respects to Lloyd bound.}$$

Intuitively, the conservation of angular momentum set a barrier to rapid complexification.



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# Sum up

- Using “complexity=action” proposal we study complexity growth of  $f(R)$  and critical gravity.
- We show action that the action growth for neutral black hole saturate the complexity growth bound.
- We also study effect of shock wave on black hole in critical gravity.
- The presence of **massive spin-2 slows down** the rate of growth.

# Outlook

- Higher derivative terms in critical gravity may slow down complexity growth. **What does it mean in dual CFT?**
- Finding the **boundary terms for null boundary and intersections.**
- Using them one may study complexity **Not complexity growth.**
- What about **Gauss-Bonnet and Lovelock** or any theory with **Riemann**? They suffers from **singularity** at  $r = 0$  where  $Riem^2$  blows up!!!
- What is the role of quantum gravity near the **singularity!**
- Promotion CV duality for higher curvature theory. We suggest

$$C_V = -\frac{1}{\ell} \int_{\mathcal{B}} \frac{\partial \mathcal{L}}{\partial R_{\mu\nu\alpha\beta}} \epsilon_{\mu\nu} \epsilon_{\alpha\beta} \quad (1)$$

$$C_V = -\frac{1}{\ell} \int_{\mathcal{B}} \left[ \frac{\partial \mathcal{L}}{\partial R_{\mu\nu\alpha\beta}} \left( a n_\mu n_\beta h_{\nu\alpha} + \frac{d+5}{d(d+1)} h_{\mu\beta} h_{\nu\alpha} \right) + c \right] \quad (2)$$

- **Lots to explore!**

# Thank you!