Quantum Null Energy Condition In two dimensions

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Based on 1901.04499

Interesting physical consequences from mathematical inequalities

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Example: unitarity constraints on physical parameters in quark mixing matrix if Standard Model correct then measurements must reproduce unitarity triangle

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green: localized in coordinate space (x), delocalized in momentum space (p) blue: mildly (de-)localized in coordinate and momentum space orange: delocalized in coordinate space (x), localized in momentum space (p)

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 - Definition: (local) inequalities on the stress tensor T_{μν} e.g. Null Energy Condition (NEC)

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For instance: Penrose singularity theorem from Raychaudhuri eq.

$$\frac{\mathrm{d}^2 \mathrm{area}}{\mathrm{d}k^2} = -\left(\frac{\mathrm{d}\operatorname{area}}{\mathrm{d}k}\right)^2 - \mathrm{shear}^2 - 8\pi G T_{kk} \leq -8\pi G T_{kk} \stackrel{\mathrm{NEC}}{\leq} 0$$

If $T_{kk} \ge 0$ (NEC) \Rightarrow focussing! (negative acceleration of area)

For experts: $\frac{d \operatorname{area}}{dk} = \theta$ is null expansion



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NEC violated by Casimir energy, accelerated mirrors, Hawking radiation, ...

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Are there quantum energy conditions? [How is 2^{nd} law saved?]

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ANEC proved under rather generic assumptions

Faulkner, Leigh, Parrikar and Wang 1605.08072 Hartman, Kundu and Tajdini 1610.05308

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Is there a local quantum energy condition?

Proposed by Bousso, Fisher, Leichenauer and Wall in 1506.02669

QNEC (in
$$D>2)$$
 is the following inequality
$$\langle T_{kk}\rangle \geq \frac{\hbar}{2\pi\sqrt{\gamma}}\,S''$$

Physical motivation from focussing properties and second law: Replace area by area + 4G (entanglement entropy) Modified Raychaudhuri eq., schematically:

$$\frac{\mathrm{d}^2 \mathrm{area}}{\mathrm{d}k^2} + 4G\,S'' = -8\pi G\,T_{kk} + 4G\,S'' \stackrel{\mathrm{QNEC}}{\leq} 0$$

requires for focussing property (=2nd law) QNEC

fineprint: above we set expansion to zero, $\frac{d \operatorname{area}}{dk} = 0$, and shear to zero; we also set the area to unity, $\sqrt{\gamma} = 1$ thus, QNEC is implied from quantum focussing for special congruences

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Obvious observations:

- ▶ if r.h.s. vanishes: semi-classical version of NEC
- if r.h.s. negative: weaker condition than NEC (NEC can be violated while QNEC holds)
- if r.h.s. positive: stronger condition than NEC (if QNEC holds also NEC holds)

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T_{kk} = T_{µν}k^µk^ν with k_µk^µ = 0 and ⟨⟩ denotes expectation value
S'': 2nd variation of EE for entangling surface deformations along k_µ

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• $T_{kk} = T_{\mu\nu}k^{\mu}k^{\nu}$ with $k_{\mu}k^{\mu} = 0$ and $\langle \rangle$ denotes expectation value • S'': 2nd variation of EE for entangling surface deformations along k_{μ} • $\sqrt{\gamma}$: induced volume form of entangling region (black boundary curve) Quantum null energy condition (QNEC) Proposed by Bousso, Fisher, Leichenauer and Wall in 1506.02669

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Proofs of QNEC in D > 2:

- For free QFTs: Bousso, Fisher, Koeller, Leichenauer and Wall, 1509.02542
- For holographic CFTs: Koeller and Leichenauer, 1512.06109
- ▶ For general CFTs: Balakrishnan, Faulkner, Khandker and Wang, 1706.09432
- QNEC from ANEC in QFTs: Ceyhan, Faulkner 1812.04683

QNEC in CFT_2

QNEC (in CFT₂) is the following inequality $2\pi \left< T_{kk} \right> \geq S'' + \frac{6}{c} \left(S' \right)^2$ $c>0 \text{ is the central charge of the CFT}_2$

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Focus here on QNEC in AdS_3/CFT_2

▶ Solutions to AdS₃ Einstein gravity with Brown–Henneaux bc's:

$$ds^{2} = \frac{dz^{2} - dx^{+} dx^{-}}{z^{2}} + \mathcal{L}^{+}(x^{+})(dx^{+})^{2} + \mathcal{L}^{-}(x^{-})(dx^{-})^{2} + \mathcal{O}(z^{2})$$

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Holographic dictionary for vev of stress tensor:

$$2\pi \left\langle \mathcal{L}^+, \, \mathcal{L}^- | T_{\pm\pm}(x^{\pm}) | \mathcal{L}^+, \, \mathcal{L}^- \right\rangle = \frac{c}{6} \, \mathcal{L}^{\pm}(x^{\pm})$$

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Uniform result for HEE from (H)RT

$$S = S^+ + S^ S^{\pm} = \frac{c}{6} \ln \left(\ell^{\pm}(x_1^{\pm}, x_2^{\pm}) / \epsilon_{\text{UV}} \right)$$

from local diffeo to Poincaré patch AdS_3 , $\{z, x^{\pm}\} \rightarrow \{z_P, x_P^{\pm}\}$

 x_1^{\pm}, x_2^{\pm} : boundary points of entangling region ϵ_{UV} : cutoff

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from local diffeo to Poincaré patch AdS₃, $\{z, x^{\pm}\} \rightarrow \{z_P, x_P^{\pm}\}$ $\ell^{\pm}(x_1^{\pm}, x_2^{\pm}) = \psi_1^{\pm}(x_1^{\pm})\psi_2^{\pm}(x_2^{\pm}) - \psi_2^{\pm}(x_1^{\pm})\psi_1^{\pm}(x_2^{\pm})$ contain solutions to Hill's equation

$$\psi^{\pm \prime \prime} - \mathcal{L}^{\pm} \, \psi^{\pm} = 0$$

with unit Wronskian $\psi_1^\pm\psi_2^{\pm\prime}-\psi_2^\pm\psi_1^{\pm\prime}=\pm 1$

▶ HEE transforms like anomalous scalar under diffeos (Wall '11)

$$\delta_{\xi}S = -\xi^{\mu}\,\partial_{\mu}S + \frac{c}{12}\,\partial_{\mu}\xi^{\mu}$$

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QNEC saturation for all states dual to Bañados geometries

QNEC saturation for vacuum-like states obvious from symmetries

see arguments in Khandker, Kundu, Li '18 for QNEC saturation in absence of bulk matter

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- Nevertheless, physically interesting examples such as far from equilibrium flow Bhaseen, Dyon, Luca, Schalm '13, Erdmenger, Fernandez, Flory, Megias, Straub, Witkowski '17

$$\mathcal{L}^{+}(x) = \mathcal{L}^{-}(-x) = \pi^{2} \left(T_{L}^{2} + \theta(x) \left(T_{R}^{2} - T_{L}^{2} \right) \right)$$

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- Far from equilibrium transport in strongly coupled CFT
- Long-time energy transport universally via steady-state
- In AdS₃/CFT₂: specific Bañados geometry with step function
- Our results imply QNEC saturation at all times

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Daniel Grumiller — Quantum Null Energy Condition

QNEC in dual of AdS_3 -Vaidya

• Vaidya = simple model for bulk matter; mass function M(t)

$$ds^{2} = \frac{1}{z^{2}} \left(-(1 - M(t)z^{2}) dt^{2} - 2 dt dz + dx^{2} \right)$$

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If time is much larger than entangling region we find QNEC saturation

$$\lim_{t_0 \gg l} \frac{S'' + \frac{6}{c} (S')^2}{2\pi \langle T_{kk} \rangle} = 1 + \mathcal{O}(\epsilon)$$

▶ Finite-*c* correction to EE (Faulkner, Lewkowycz, Maldacena '13)

 $S = S_{\rm HRT} + S_{\rm bulk}$

Finite- $c\ {\rm corrections}\ {\rm to}\ {\rm EE}\ {\rm and}\ {\rm QNEC}$

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$$\phi = \frac{a}{\sqrt{2\pi}} \frac{e^{-2iht}}{(1+r^2)^h} + \frac{a^{\dagger}}{\sqrt{2\pi}} \frac{e^{2iht}}{(1+r^2)^h} \qquad [a, a^{\dagger}] = 1$$

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$$G_i(r) = 1 - \frac{12h}{c} \left(1 - (i-1)/(1+r^2)^{2h-1}\right) + \mathcal{O}(h^2/c^2)$$

NOT a Bañados geometry

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 $S_{\rm bulk}:$ bulk entanglement of low energy dof's

▶ Consider global AdS₃ backreacted by bulk scalar field ϕ with mass $m^2 = 4h(h-1) \ge -1$ (see Belin, Iqbal, Lokhande '18)

$$\phi = \frac{a}{\sqrt{2\pi}} \frac{e^{-2iht}}{(1+r^2)^h} + \frac{a^{\dagger}}{\sqrt{2\pi}} \frac{e^{2iht}}{(1+r^2)^h}$$
$$ds^2 = -\left(r^2 + G_1(r)^2\right) dt^2 + \frac{dr^2}{r^2 + G_2(r)^2} + r^2 d\varphi^2$$
$$G_i(r) = 1 - \frac{12h}{c} \left(1 - (i-1)/(1+r^2)^{2h-1}\right) + \mathcal{O}(h^2/c^2)$$

Asymptotically conical defect, $2\pi \langle T_{\pm\pm} \rangle = -\frac{c}{24} + h + O(h^2/c)$

▶ Calculate S_{HRT} for interval bounded by origin and $t = \lambda$, $\varphi = \Delta \varphi - \lambda$

Calculate S_{HRT} for interval bounded by origin and t = λ, φ = Δφ − λ
 Take λ-derivatives to get QNEC

$$S_{\rm HRT}'' + \frac{6}{c} \left(S_{\rm HRT}' \right)^2 = -\frac{c}{24} + h - h \, \frac{\sqrt{\pi} \, \Gamma[2h+2]}{4\Gamma[2h+\frac{3}{2}]} \, \sin^{4h-2} \frac{\Delta \varphi}{2} + \mathcal{O}(1/c)$$

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$$\delta S_{\text{bulk}} = 2\pi \left\langle \delta H_{\text{bulk}} \right\rangle + \mathcal{O}(\delta^2)$$

$\delta:$ small perturbation of the vacuum

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$$\langle \delta H_{\text{bulk}} \rangle = \int_{\omega, k} \omega \left(|\alpha_{\omega, k}|^2 + |\beta_{\omega, k}|^2 \right)$$

 $\begin{aligned} &\alpha_{\omega,\,k}, \beta_{\omega,\,k} \text{: Bogoliubov coefficients} \\ & \text{calculated in appendix of Belin, Iqbal, Lokhande '18} \\ & \text{map entanglement wedge to Rindler space (see Casini, Huerta, Myers '11) and} \\ & \text{expand scalar field } \phi \text{ in Rindler modes} \end{aligned}$

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QNEC holds. QNEC saturation up to $\mathcal{O}(\Delta \varphi^{4h})$

Pragmatically, need small parameter for S_{bulk}

reason: for Renyi entropies need reduced density matrix

- ρ : bulk density matrix
- ρ_1 : reduced density matrix
- ho_0 : reduced density matrix for vacuum
- $\delta \rho = \rho_1 \rho_0$: small for small interval
- e.g. small interval limit: $\Delta \varphi$ is small parameter

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Recovered from CFT_2 calculation

$$\delta S_n = \frac{1}{1-n} \ln \left[\left(\frac{1}{n} \sin(\Delta \varphi/2) \right)^{4nh} \text{Hf} \right]$$

by taking suitable limit of Hafnian for $\Delta \varphi = \pi - \epsilon$

$$\mathrm{Hf} = \mathrm{Hf}_0 + (n-1)\,\mathrm{Hf}_1 + \dots$$

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QNEC non-saturation at half-interval:

$$2\pi \langle T_{\pm\pm} \rangle - S'' - \frac{6}{c} (S')^2 |_{\Delta \varphi = \pi, c \gg h \gg 1} = \frac{1}{4} h$$

We scratched the surface of QNEC in AdS_3/CFT_2

- QNEC saturates for vacuum-like states
- Curious half-saturation for states dual to AdS₃-Vaidya
- Leading 1/c correction
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 - QNEC-gap at half-interval
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- QNEC as constraint for semi-classical model building (black holes)
- QNEC saturation as quantum equilibration?
- QNEC-like inequalities for QFTs that are not CFTs?
 - Galilean/ultrarelativistic CFTs and flat space holography? Bagchi et al '09-'18
 - warped CFTs (and their holographic duals)? Detournay, Hartman, Hofman '12
 - other non-standard QFTs?

Much to be learned about QNEC and its potential applications



Families of geodesics used in AdS3-Vaidya for numerical QNEC determination

Thanks for your attention!